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# Structure and Phase Transitions in Multiferroics: The Case of $\text{BiFeO}_3$



Phil Lightfoot, EastCHEM, University of St Andrews

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# Structure and Phase Transitions in Multiferroics: The Case of $\text{BiFeO}_3$

With thanks to:

**Dr Donna Arnold**

Dr Finlay Morrison  
Dr Richard Goff  
Ms Alex Gibbs  
Dr Kevin Knight (ISIS)

Prof. Jim Scott (Cambridge)  
Dr Simon Redfern (Cambridge)  
Dr Gustau Catalan (Cambridge)

EPSRC, STFC



Perhaps the most studied multiferroic?\*

Ferroelectric  $T_C \sim 820 \text{ }^\circ\text{C}$ , magnetic  $T_N \sim 360 \text{ }^\circ\text{C}$

Three suggested (structural) phase transitions above RT:

1. At  $T_C$ :  $\alpha$  (ferroelectric) to  $\beta$  (paraelectric)
2. Region 925 – 933  $^\circ\text{C}$ :  $\beta$  (insulator) to  $\gamma$  (metallic?)
3. Near 185  $^\circ\text{C}$ : 'Polomska' transition

\*Catalan and Scott, *Adv. Mater.* *In press.*

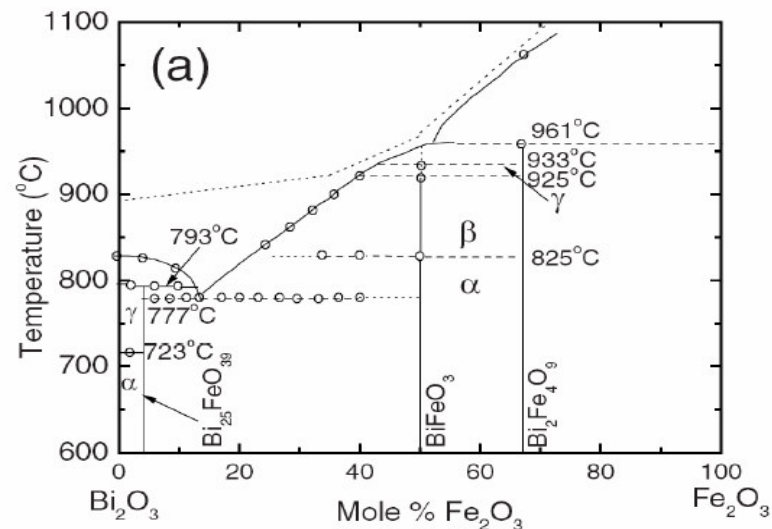
Structural nature of these transitions??

## Challenges:

1.  $\text{BiFeO}_3$  is metastable above 700 °C
2. X-rays can barely 'see' O atoms in the presence of Bi
3. Several previous studies have therefore 'failed'

## Solution:

Powder neutron diffraction (HRPD-ISIS)



Palai et al., PRB, 2008, 77,014110

# BiFeO<sub>3</sub> ferroelectric phase transition

COMMUNICATION

ADVANCED  
MATERIALS

DOI: 10.1002/adma.200800218

## The Ferroic Phase Transitions of BiFeO<sub>3</sub>\*\*

By Sverre M. Selbach, Thomas Tybell, Mari-Ann Einarsrud, and Tor Grande\*

Multiferroics<sup>[1]</sup> have rapidly gained attention due to a wide range of potential applications<sup>[2]</sup> in microelectronic and spintronic devices, and the possibility of controlling magnetic order by electric fields,<sup>[3]</sup> or vice versa,<sup>[4]</sup> through magnetolectric coupling.<sup>[5]</sup> BiFeO<sub>3</sub> is termed a multiferroic material due to coexistence of ferroelasticity, antiferromagnetism, and ferroelectricity.<sup>[6–9]</sup> A large spontaneous polarization<sup>[9–11]</sup>  $P_S$  of 88–100  $\mu\text{C cm}^{-2}$  and a high  $T_C$  of 830 °C makes

structure of the high temperature paraelectric phase is centrosymmetric  $R\bar{3}c$ . The ferroelectric transition is first order and accompanied by discontinuous volume and abrupt changes in atomic positions. A continuous lattice anomaly is associated with the second order magnetic phase transition at  $T_N$ . Electrical conductivity anomalies associated with both phase transitions are observed. The present findings provide a basis for fundamental insight to the multiferroic properties.

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## Atomic displacements in BiFeO<sub>3</sub> as a function of temperature: neutron diffraction study

A. Palewicz,<sup>1\*</sup> R. Przeniosło,<sup>2\*</sup> I. Sosnowska<sup>3</sup> and A. W. Hewat<sup>4</sup>

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The parameters of the crystal structure of BiFeO<sub>3</sub>, described within the space group  $R\bar{3}c$ , have been determined by high-resolution neutron powder diffraction for temperatures from 293 to 923 K. It was found that there is a local minimum for the rhombohedral angle  $\alpha_{\text{rh}}$ , near the Néel temperature  $T_N \approx 640$  K, a gradual rotation of the FeO<sub>6</sub> octahedra and an increase of the Fe—O—Fe angle. The displacement of the Bi<sup>3+</sup> ions from the FeO<sub>6</sub> octahedra which influence the electric polarization decreases with temperature. One of the Bi—Fe distances also has a local maximum near  $T_N$ . The atomic vibrations of Bi<sup>3+</sup> and O<sup>2-</sup> ions show a significant anisotropy.

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Accepted 15 May 2007

PHYSICAL REVIEW B 77, 014110 (2008)

### $\beta$ phase and $\gamma$ - $\beta$ metal-insulator transition in multiferroic BiFeO<sub>3</sub>

R. Palai,<sup>1</sup> R. S. Katyar,<sup>1</sup> H. Schmid,<sup>2</sup> P. Tissot,<sup>2</sup> S. J. Clark,<sup>3</sup> J. Robertson,<sup>4</sup> S. A. T. Redfern,<sup>5</sup> G. Catalan,<sup>5</sup> and J. F. Scott<sup>5</sup>

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(Received 17 August 2007; revised manuscript received 17 October 2007; published 28 January 2008)

We report on extensive experimental studies on this film, single crystal, and ceramics of multiferroic bismuth ferrite BiFeO<sub>3</sub> using differential thermal analysis, high-temperature polarized light microscopy, high-temperature and polarized Raman spectroscopy, high-temperature x-ray diffraction, dc conductivity, optical absorption and reflectivity, and domain imaging, and show that epitaxial (001) thin films of BiFeO<sub>3</sub> are clearly monoclinic at room temperature, in agreement with recent synchrotron studies but in disagreement with all other earlier reported results. We report an orthorhombic order-disorder  $R$  phase between 830 and 925

PHYSICAL REVIEW B 78, 134108 (2008)

### Phase stability and structural temperature dependence in powdered multiferroic BiFeO<sub>3</sub>

R. Haumont,<sup>1</sup> Igor A. Kornev,<sup>2\*</sup> S. Lisenkov,<sup>2</sup> L. Bellaiche,<sup>2</sup> J. Kreisel,<sup>3</sup> and B. Dkhil<sup>4</sup>

<sup>1</sup>Laboratoire de Physico-Chimie de l'Etat Solide, ICMMO, CNRS-UMR8182, Université Paris XI, 91405 Orsay, France

<sup>2</sup>Physics Department, University of Arkansas, Fayetteville, Arkansas 72701, USA

<sup>3</sup>Laboratoire Matériaux et Génie Physique (CNRS), Grenoble Institute of Technology, Minatex 38016 Grenoble, France

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We report a temperature-dependent investigation of the multiferroic perovskite bismuth ferrite BiFeO<sub>3</sub> (BFO) by using x-ray powder diffraction together with differential scanning calorimetry measurements. Our results provide evidence that the paraelectric phase above  $T_C \approx 820$  °C is not cubic but distorted and can be well refined in a monoclinic  $P2_1/m$  space group. An equivalent structure can be reconstructed based on the  $C2/m$  monoclinic space group and by assuming two types of bismuth sites. The marked change of the cell volume at  $T_C$  provides evidence for the first-order nature of the  $R\bar{3}c$ -to- $P2_1/m$  transition. The high-temperature  $P2_1/m$  phase is centrosymmetric and characterized by (i) strong oxygen octahedra tilting along the  $b$  axis; (ii) the occurrence of antiferroelectric displacements of the Fe cations; and (iii) an interesting lamellar structure characterized by two different types of BiO<sub>12</sub> cages. The temperature-induced lamellar structure suggests a significant electronic rearrangement in terms of chemical bonding, which in turn might condition anisotropic electronic properties. The occurrence of a lamellar structure provides also an understanding of why BFO decomposes suddenly at higher temperatures. Finally, an anomaly in the evolution of the cell parameters at  $T_N$  underlines the spin-lattice coupling in proximity of the magnetic transition.

PRL 99, 227602 (2007)

PHYSICAL REVIEW LETTERS

week ending  
30 NOVEMBER 2007

### Finite-Temperature Properties of Multiferroic BiFeO<sub>3</sub>

Igor A. Kornev,<sup>1,2</sup> S. Lisenkov,<sup>1</sup> R. Haumont,<sup>3</sup> B. Dkhil,<sup>4</sup> and L. Bellaiche<sup>1</sup>

<sup>1</sup>Physics Department, University of Arkansas, Fayetteville, Arkansas 72701, USA

<sup>2</sup>Mads Clausen Institute for Product Innovation, University of Southern Denmark, Alsion 2, DK-6400 Sønderborg, Denmark

<sup>3</sup>Laboratoire de Physico-Chimie de l'Etat Solide (CNRS), ICMMO, Université Paris XI, 91405 Orsay, France

<sup>4</sup>Laboratoire Structures, Propriétés et Modélisation des Solides, Ecole Centrale Paris, CNRS-UMR8580,

Grande Voie des Vignes, 92295 Châtenay-Malabry Cedex, France

(Received 5 July 2007; published 30 November 2007)

An effective Hamiltonian scheme is developed to study finite-temperature properties of multiferroic BiFeO<sub>3</sub>. This approach reproduces very well (i) the symmetry of the ground state, (ii) the Néel and Curie temperatures, and (iii) the intrinsic magnetolectric coefficients (that are very weak). This scheme also predicts (a) an overlooked phase above  $T_C \approx 1100$  K that is associated with antiferrodistortive motions, as consistent with our additional x-ray diffractions, (b) improperlike dielectric features above  $T_C$ , and (c) that the ferroelectric transition is of first order with no group-subgroup relation between the paraelectric and polar phases.



The Edinburgh and St Andrews  
Research School of Chemistry

School of Chemistry, University of St Andrews



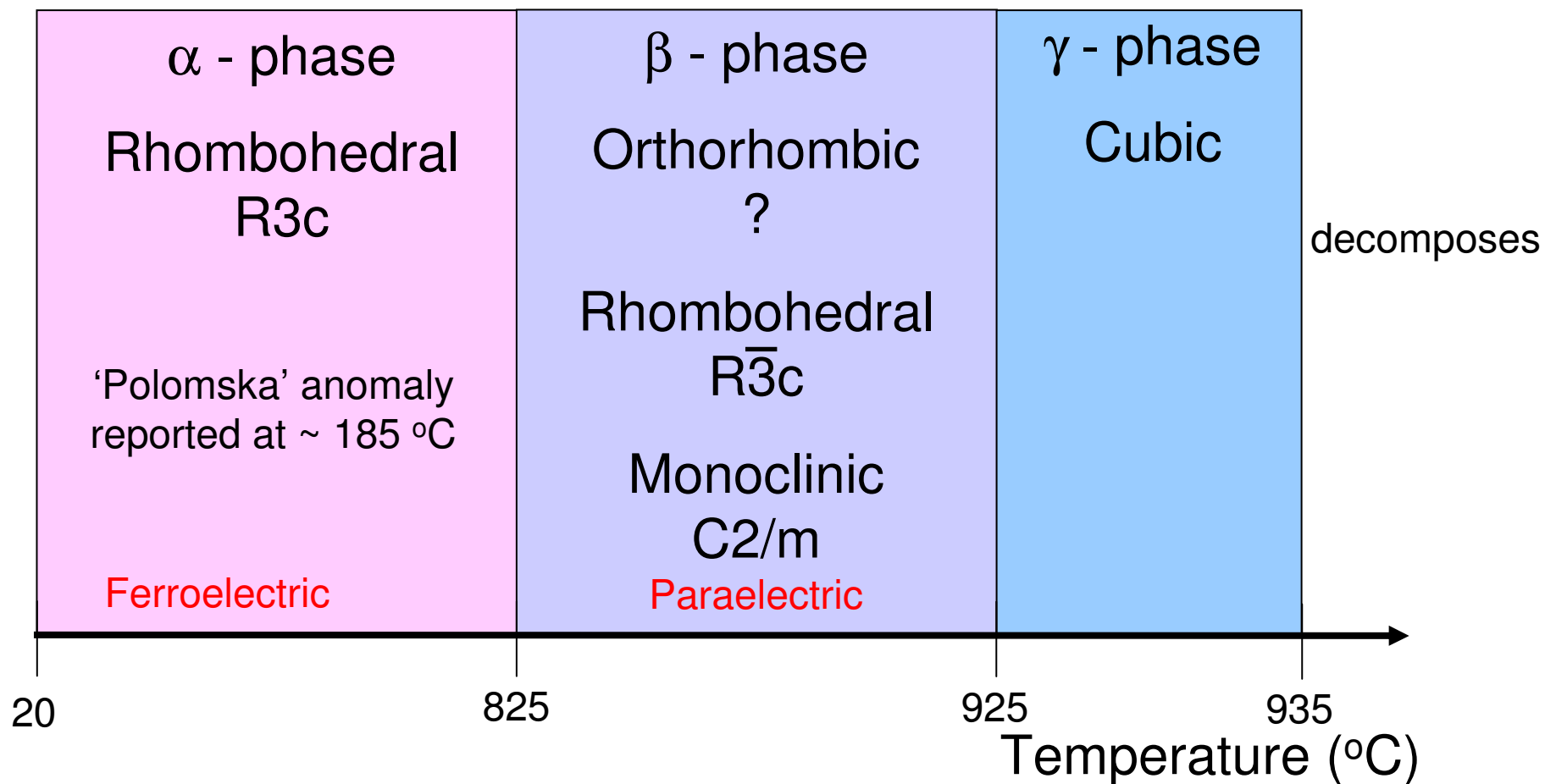
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## The $\alpha - \beta$ and $\beta - \gamma$ transitions

Recent studies (2007 – 08):

1. PND: RT – 700 °C – retains RT structure (R3c – paraelectric).  
Sample degrades at high T
2. PXR: RT - 900 °C – suggests R3c – R-3c transition at  $T_C$
3. Theoretical: suggests R3c – I4/mcm- Pm-3m  
(PXR suggests C2/m not I4/mcm!)
4. PXR – P2<sub>1</sub>/m or C2/m above  $T_C$
5. Raman, XRD, conductivity etc: R3c – orth(?) – cubic  
(M-I transition  $\gamma - \beta$ )

## The $\alpha - \beta$ and $\beta - \gamma$ transitions



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Our PND experiments (HRPD at ISIS):

(Sample preparation): elimination of  $\text{Bi}_2\text{Fe}_4\text{O}_9$

1. Sept 08 – to study the  $\alpha - \beta$  transition (what happens at  $T_C$ ?)

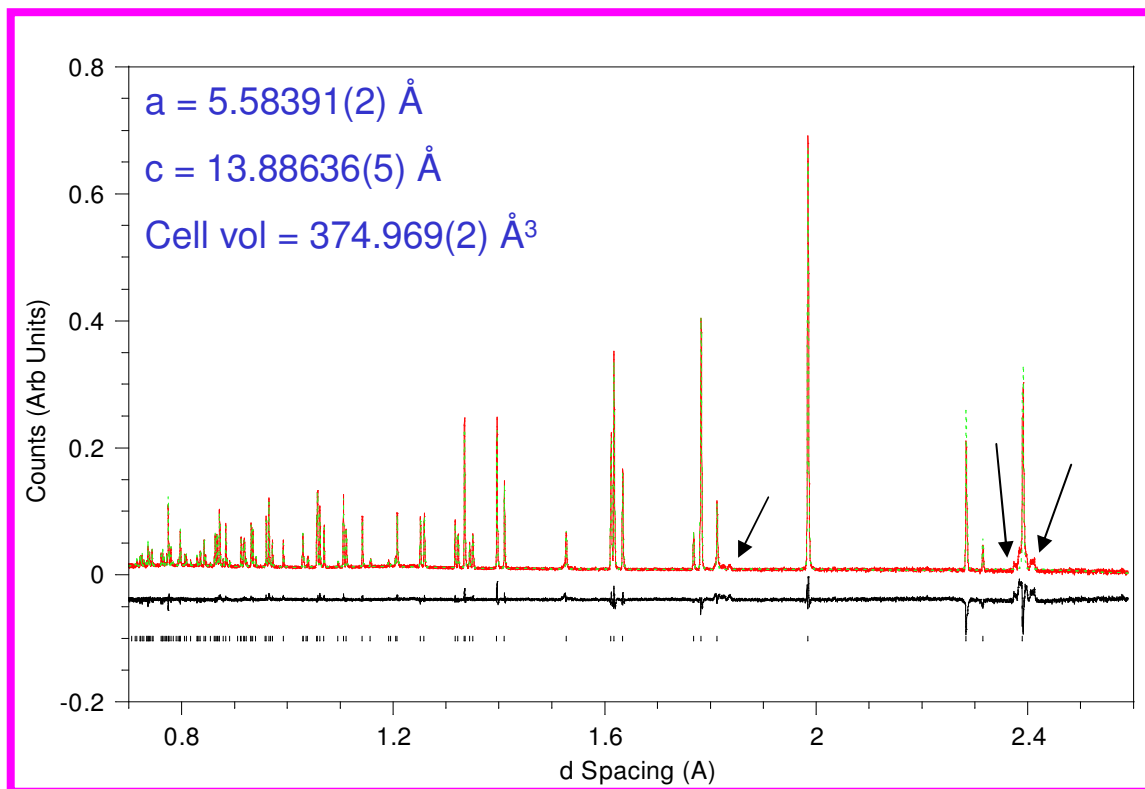
V can / Sealed quartz tube RT – 900 °C; 20 min – 1hr runs

2. March 09 – to revisit the  $\beta - \gamma$  and ‘Polomska’ transitions

Open vessel – flowing  $\text{N}_2$ ; 900 – 960 °C; 5 min runs



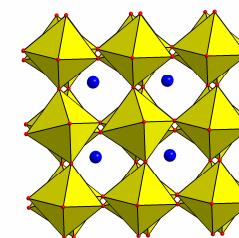
# BiFeO<sub>3</sub> α – phase



R3c

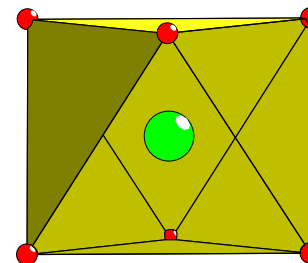
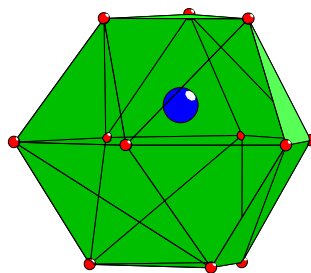
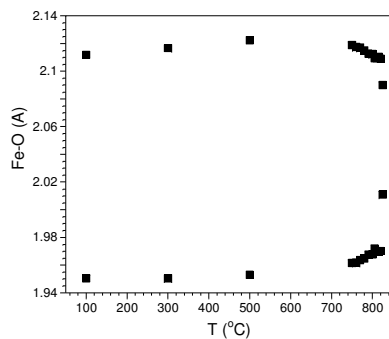
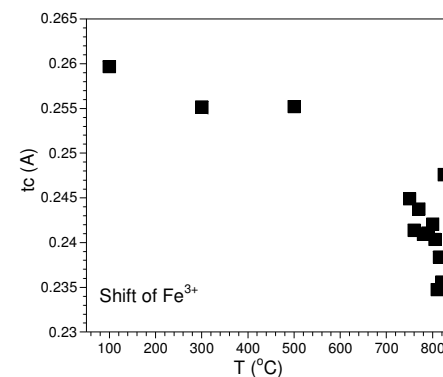
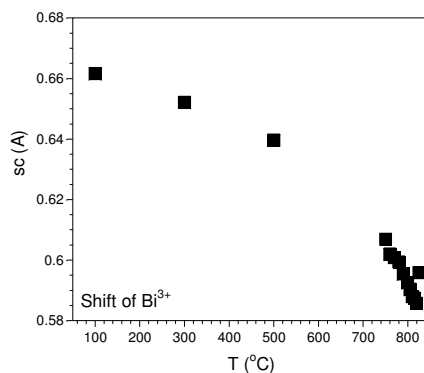
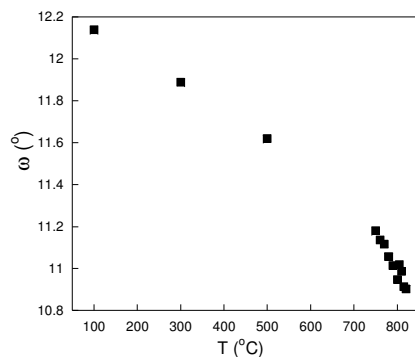
Ferroelectric

$a^-a^-a^-$  tilt



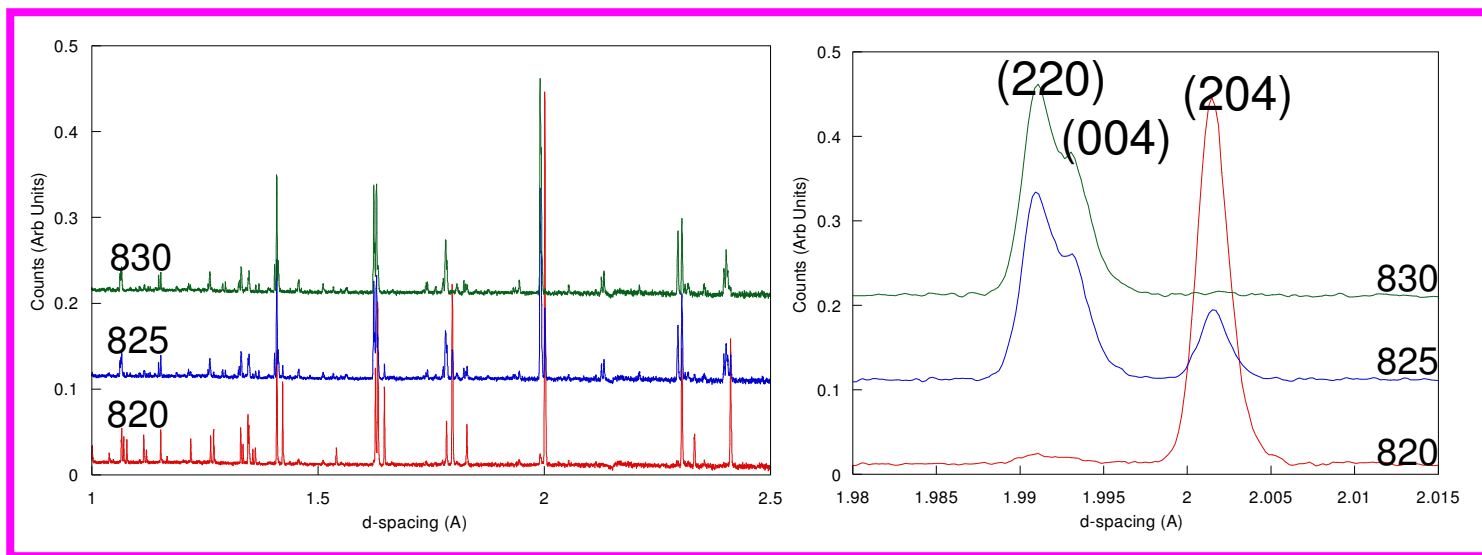
Phys. Rev. Lett., 2009, 102, 027602

# BiFeO<sub>3</sub> evolution of the $\alpha$ phase



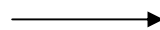
Phys. Rev. Lett., 2009, 102, 027602

## BiFeO<sub>3</sub> ferroelectric phase transition ( $\alpha - \beta$ )



1<sup>st</sup> order transition: phases co-exist at 825 °C

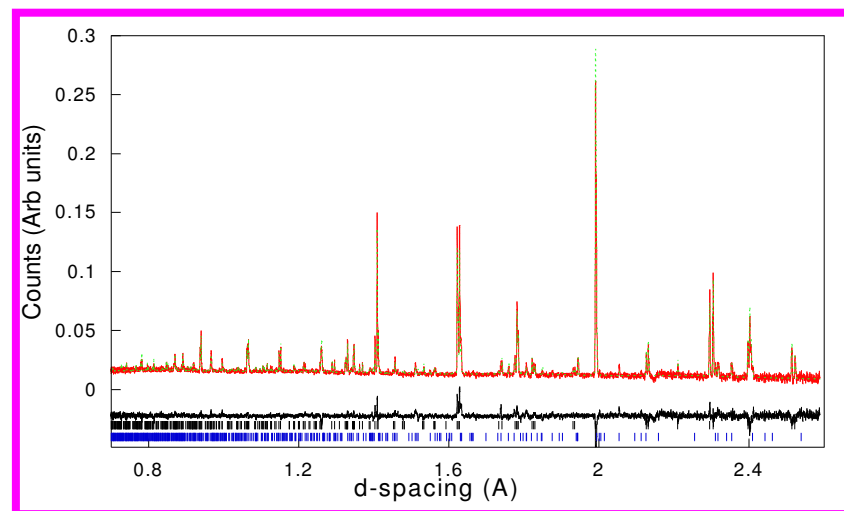
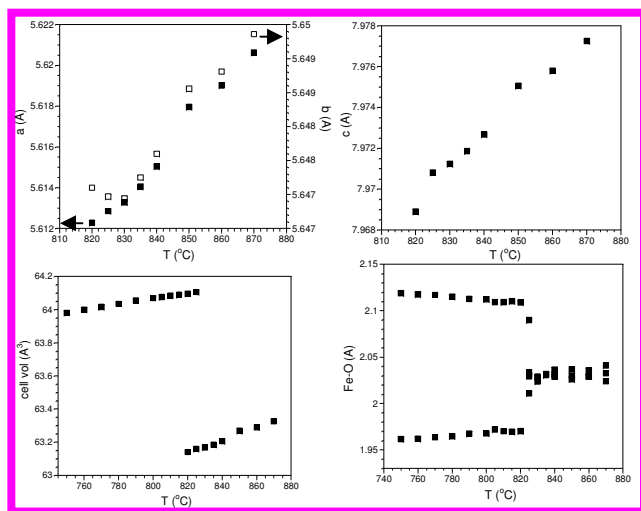
Trigonal R3c,  $a^-a^-a^-$   
Ferroelectric



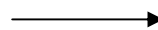
Orthorhombic Pbnm,  $a^-a^+b^+$   
Paraelectric

Phys. Rev. Lett., 2009, 102, 027602

# BiFeO<sub>3</sub> ferroelectric phase transition ( $\alpha - \beta$ )

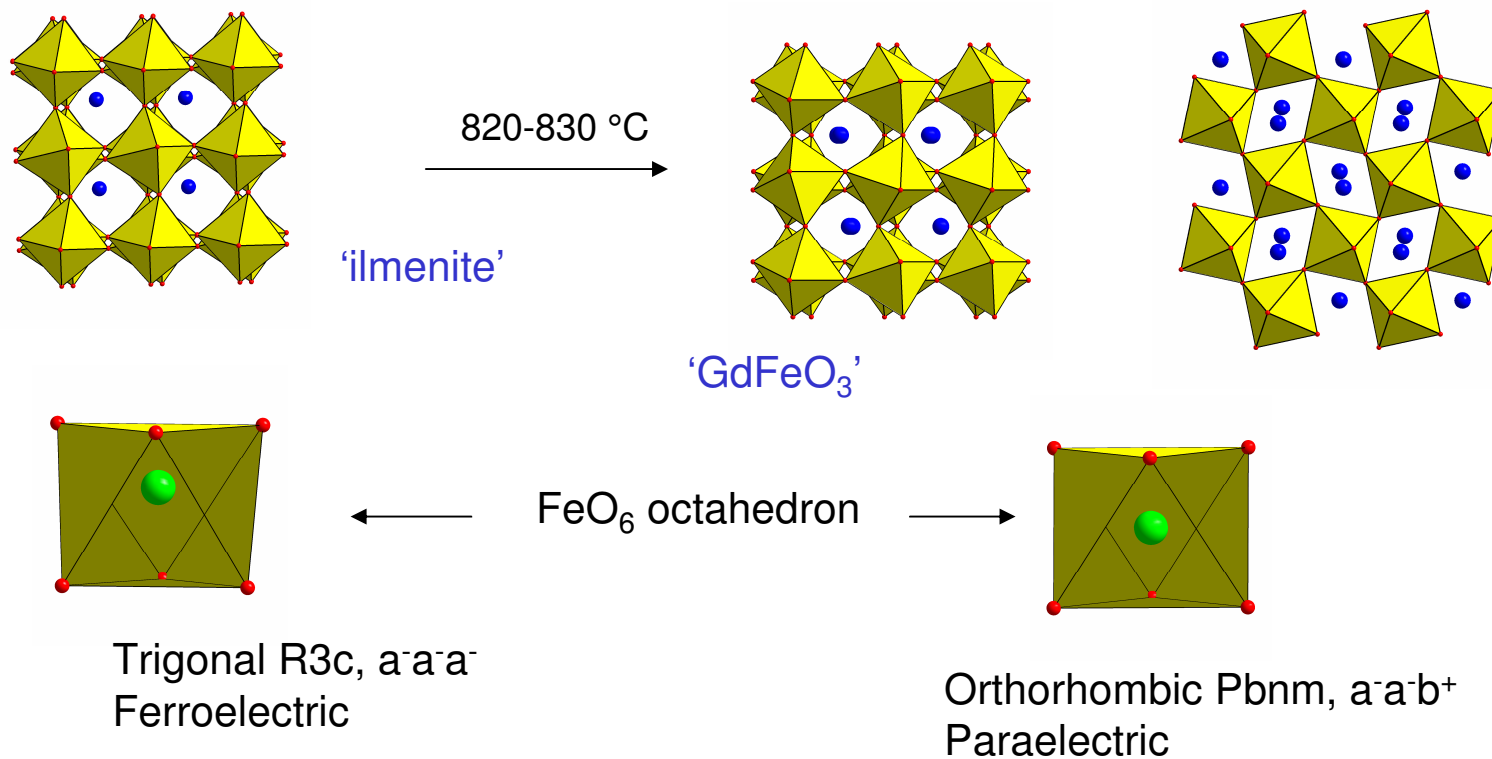


Trigonal R3c,  $a^-a^-a^-$   
Ferroelectric



Orthorhombic Pbnm,  $a^-a^+b^+$   
Paraelectric

# BiFeO<sub>3</sub> ferroelectric phase transition ( $\alpha - \beta$ )

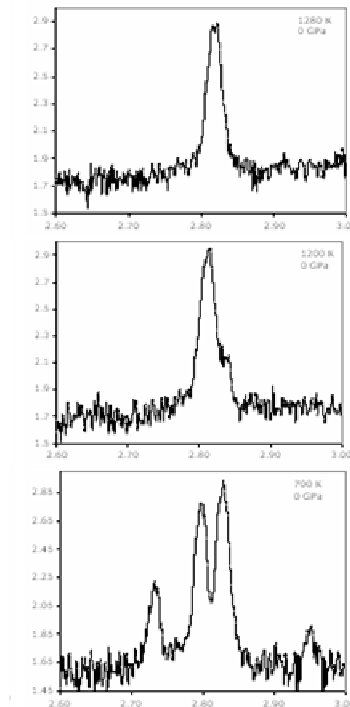
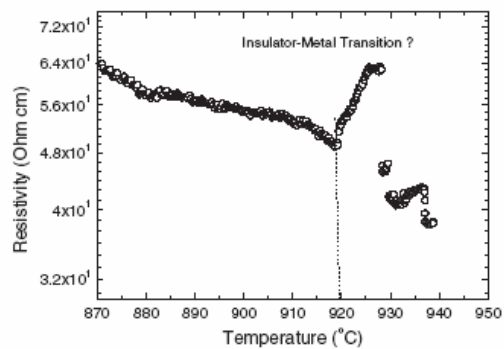


Phys. Rev. Lett., 2009, 102, 027602

## BiFeO<sub>3</sub> β – γ phase transition

Palai et al. suggest cubic γ-phase:

Appears 'cubic' from PXRD



'Cubic'

Orth

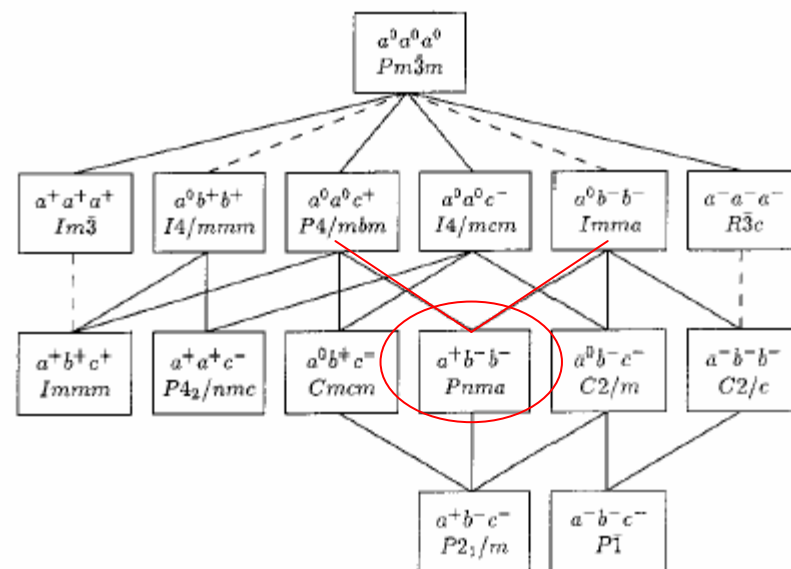
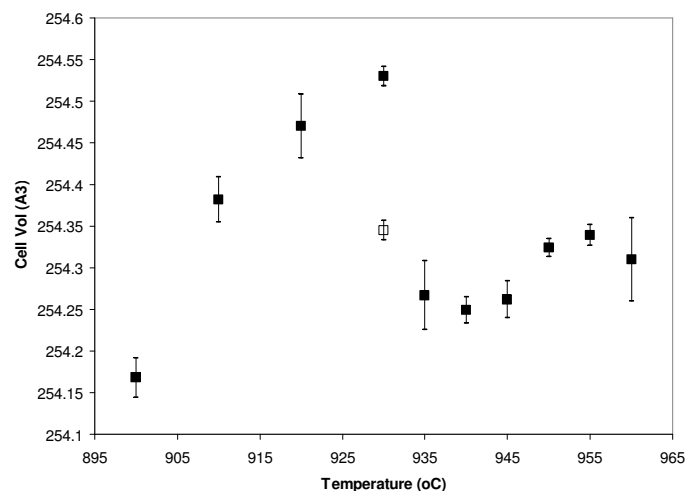
R3c

Our sample decomposes 'quickly' to Bi<sub>2</sub>Fe<sub>4</sub>O<sub>9</sub> under conditions used in Experiment 1 (7% at 835 C, 64% at 870 C)

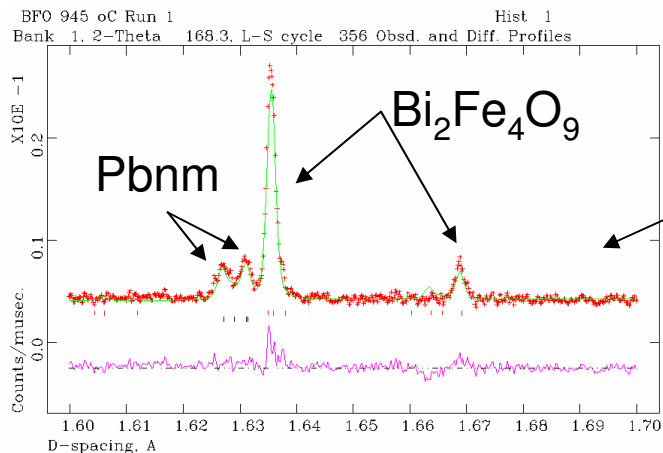
∴ A 'quicker' experiment was carried out to isolate the γ phase (metallic, cubic?)

# BiFeO<sub>3</sub> β – γ phase transition

Unit cell volume suggests abrupt transition.  
 BUT Rietveld still suggests orthorhombic



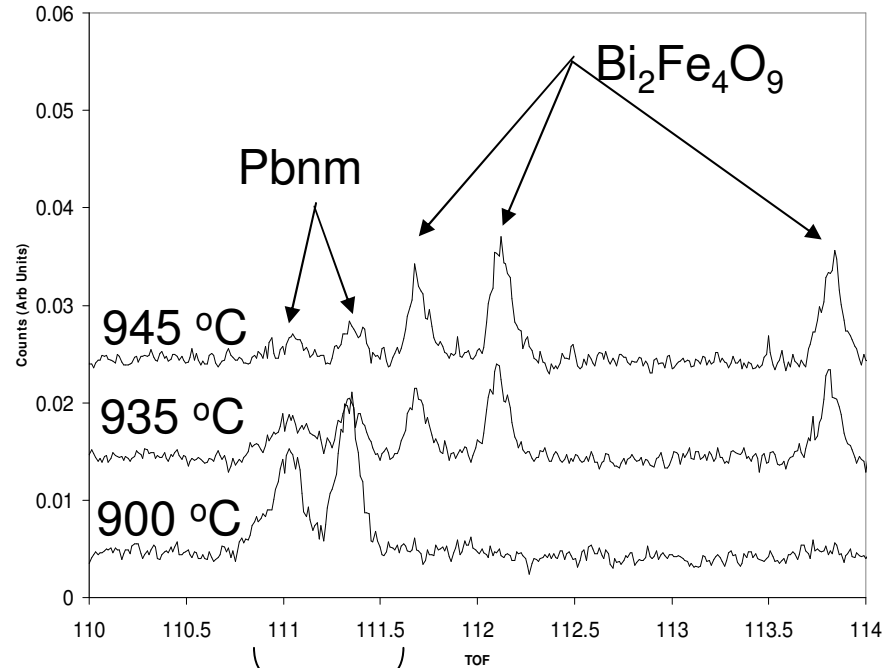
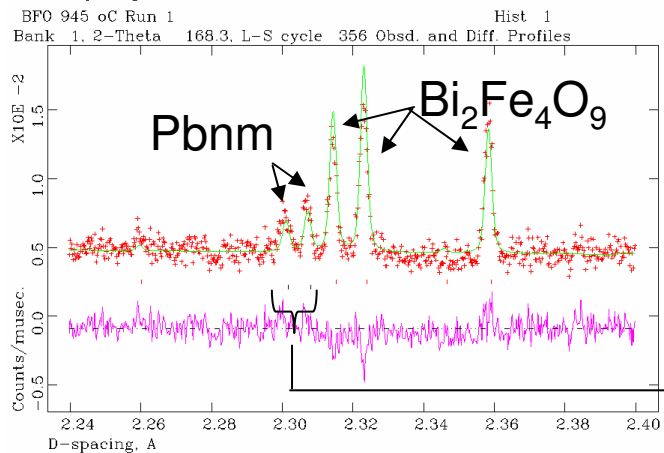
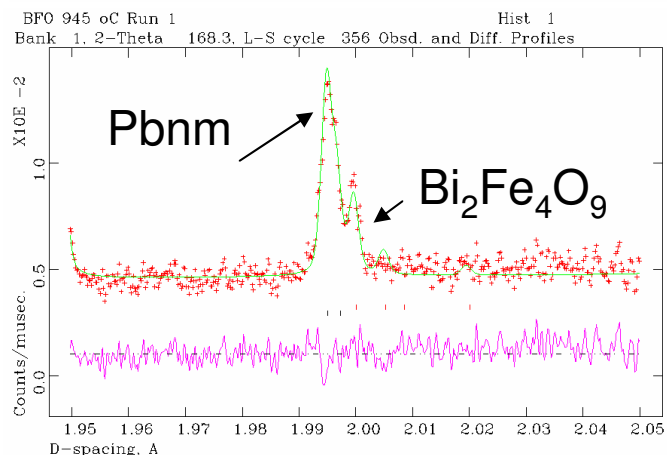
Howard et al., Acta Cryst., 1998, B54, 782



## Persistence of ortho phase to HT

Refinement of 945 °C data

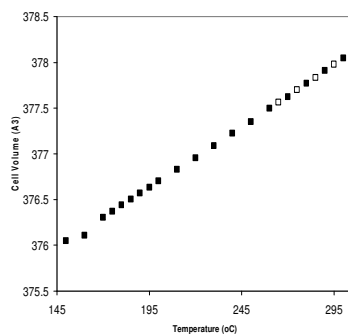
Close inspection of the PND data shows no evidence of the cubic phase - characteristic orthorhombic peaks persist at 945 °C



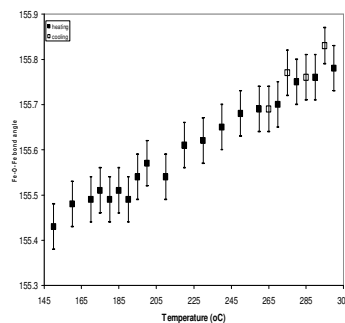
In particular this should be a single peak for cubic or tetragonal symmetry and it clearly remains split



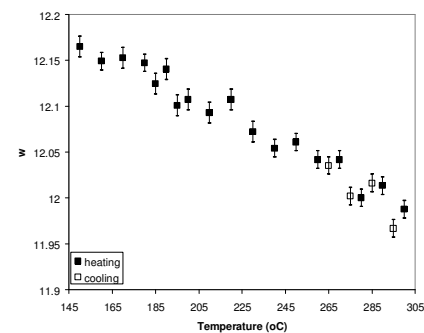
## 'Polomska' transition @ 185 °C



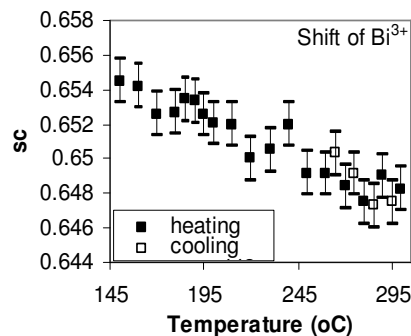
Cell volume



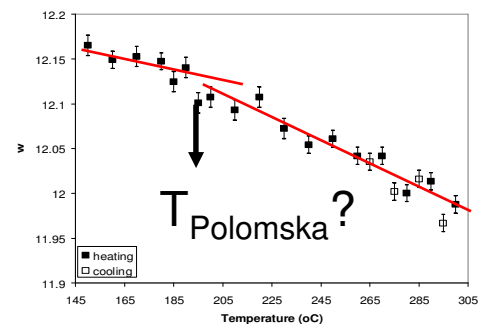
Fe-O-Fe bond angle



Tilt angle



No obvious anomalies,  
BUT....



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## Summary

1.  $\alpha - \beta$  phase transition : R3c – Pbnm
2.  $\beta - \gamma$  phase transition: *probably* Pbnm – Pbnm (Imma?)
3. Cubic phase? *might* depend on kinetics, or is actually *orthorhombic*
4. Polomska transition – no observable (?) structural transition  
(may be magnetic ??)
5. Further phase transitions below RT...
6. WISH diffractometer may clarify (4 and 5).

Thank you for listening!