Short Range Order in High-Entropy Alloys: First Principles Theory and Atomistic Modelling

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Our Description

- On lattice.
- Specify configuration by $\{\xi_{i\alpha}\}$.
- Interested in the average value of these:

$$c_{i\alpha} \equiv \langle \xi_{i\alpha} \rangle.$$

- High temperature, homogeneous: $c_{i\alpha} = c_{\alpha}$.
- ▶ Perturb homogeneous state $c_{i\alpha} = c_{\alpha} + \delta c_{i\alpha}$ and see what favourable correlations are.



(1)

Cantor-Wu Alloys

Summary O

Concentration Waves

$$c_{i\alpha} = c_{\alpha} + \eta_{\alpha} \frac{1}{2} \left(e^{i\mathbf{q}\cdot\mathbf{R}_{i}} + e^{-i\mathbf{q}\cdot\mathbf{R}_{i}} \right), \quad \mathbf{q} = \left(\frac{1}{2}, \frac{1}{2} \right).$$

$$(2)$$

$$\eta = (0,0) \qquad \eta = (0.25, -0.25) \qquad \eta = (0.5, -0.5)$$

Our Approach oo●o	Cantor-Wu Alloys 0000	

Energetics

- Evaluate cost of fluctuations ab initio via DFT, using KKR-CPA and a linear response theory^{ab}.
- Mean-field free energy:

$$G = \beta^{-1} \sum_{i\alpha} c_{i\alpha} \log c_{i\alpha} - \sum_{i\alpha}' \nu_{i\alpha} c_{i\alpha} + \langle \Omega_{el} \rangle_0 [\{c_{i\alpha}\}] \quad (3)$$

Important quantity:

$$S_{i\alpha;j\alpha'}^{(2)} \equiv \frac{\partial^2 \langle \Omega_{\mathsf{el}} \rangle_0}{\partial c_{i\alpha} \partial c_{j\alpha'}} \rightsquigarrow \Psi_{\alpha\alpha'}^{-1}(\mathbf{q}) \tag{4}$$

^aS. N. Khan, J. B. Staunton, G. M. Stocks, Phys. Rev. B **93** 054206 (2016) ^bB. L. Gyorffy, G. M. Stocks, Phys. Rev. Lett. **50**, 374 (1983)

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Energetics

Bragg-Williams Hamiltonian for atomistic modelling:

$$H(\{\xi_{i\alpha}\}) = \frac{1}{2} \sum_{\substack{i\alpha\\j\alpha'}} V_{i\alpha;j\alpha'} \xi_{i\alpha} \xi_{j\beta} + \sum_{i\alpha} \nu_{i\alpha} \xi_{i\alpha}.$$
 (5)

► If *H* as above,
$$V_{i\alpha;j\alpha'}$$
 are exactly $-S^{(2)}_{i\alpha;j\alpha'}$.



Our Approach	Cantor-Wu Alloys	Summa
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Cantor-Wu Alloys

Prototypical FCC high entropy alloy is the "Cantor alloy" - NiCoFeMnCr.



> At what temperature will order emerge? What is the nature of order?

We looked at a series of three Cantor-Wu alloys: NiCoCr, NiCoFeCr, NiCoFeMnCr, along with some binary subsystems¹.

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¹C. D. Woodgate, J. B. Staunton, Phys. Rev. B **105** 115124 (2022)

Linear Response

Eigenvalues of $\Psi_{\alpha\alpha'}^{-1}(\mathbf{q})$ around IBZ at 1200K.



C. D. Woodgate, J. B. Staunton, Phys. Rev. B 105 115124 (2022)

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Atomistic Simulations

NiCoCr



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Cantor-Wu Alloys

Atomistic Simulations



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Summary

- SRO dominated by Co-Cr and Cr-Cr correlations.
- Fe, Mn, serve a diluting effect and stabilise solid solution.
- Predicted ordering temperatures are low, consistent with experiment.
- Method obtains results comparable with conventional techniques for a fraction of the computational cost.

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Density of States



Linear Response

System	<i>k</i> -vector	δc_1	δc_2	δc_3	δc_4	δc_5	$T_{ m order}$ (K)
NiCr	(0, 0, 0.6)	0.70711	-0.70711				200
CoCr	(0, 0, 1)	0.70711	-0.70711				793
NiCo	(0, 0.6, 0.6)	0.70711	-0.70711				83
NiCoCr	(0, 0, 1)	-0.034613	-0.68916	0.72378			606
NiCoFeCr	(0, 0, 1)	0.013671	-0.68858	-0.048489	0.72340		404
NiCoFeMnCr	(0, 0, 1)	0.033024	-0.68516	-0.081509	0.010666	0.72298	281

Fitting to Pairwise Interaction

We approximate the true interaction by fitting to a Bragg-Williams-like Hamiltonian. That is, we approximate the *ab initio* data by $V_{\alpha\beta}(\mathbf{q})$, where

$$V_{\alpha\beta}(\mathbf{q}) = V_{\alpha\beta}^{0} + V_{\alpha\beta}^{1} \left(\sum_{\{\mathbf{R}_i\}_1} e^{i\mathbf{q}\cdot\mathbf{R}_i} \right) + \dots + V_{\alpha\beta}^{N} \left(\sum_{\{\mathbf{R}_i\}_N} e^{i\mathbf{q}\cdot\mathbf{R}_i} \right), \quad (6)$$

a constant plus a nearest-neighbour term, next-nearest neighbour term, etc.

Pairwise Interactions

$V^{(1)}_{lphaeta}$	Ni	Co	
Ni	0.023	-0.023	
Co	-0.023	0.023	
$V^{(1)}_{\alpha\beta}$	Ni	Cr	
Ni	0.199	-0.199	
Cr	-0.199	0.199	
$V^{(1)}_{lphaeta}$	Co	Cr	
$\frac{V^{(1)}_{\alpha\beta}}{Co}$	Co 1.026	Cr -1.026	
$\begin{array}{c} V^{(1)}_{\alpha\beta} \\ \hline \text{Co} \\ \text{Cr} \end{array}$	Co 1.026 -1.026	Cr -1.026 1.026	
$\frac{V^{(1)}_{\alpha\beta}}{Co}$	Co 1.026 -1.026	Cr -1.026 1.026	
$\begin{array}{c} V^{(1)}_{\alpha\beta}\\ \hline Co\\ Cr\\ V^{(1)}_{\alpha\beta} \end{array}$	Co 1.026 -1.026 Ni	Cr -1.026 1.026 Co	Cr
	Co 1.026 -1.026 Ni -0.218	Cr -1.026 1.026 Co 0.682	Cr -0.465
$\begin{array}{c} V^{(1)}_{\alpha\beta}\\ \hline \text{Co}\\ \text{Cr}\\ \hline V^{(1)}_{\alpha\beta}\\ \hline \text{Ni}\\ \text{Co}\\ \end{array}$	Co 1.026 -1.026 Ni -0.218 0.682	Cr -1.026 1.026 Co 0.682 0.672	Cr -0.465 -1.351

Pairwise Interactions

$V^{(1)}_{\alpha\beta}$	Ni	Co	Fe	Cr
Ni	-0.338	0.606	0.097	-0.367
Co	0.606	0.656	-0.049	-1.213
Fe	0.097	-0.049	-0.019	-0.029
Cr	-0.367	-1.213	-0.029	1.609
$V^{(2)}_{lphaeta}$	Ni	Co	Fe	Cr
Ni	0.316	0.058	-0.061	-0.313
Co	0.058	0.005	-0.007	-0.057
Fe	-0.061	-0.007	-0.010	0.058
Cr	-0.313	-0.057	0.058	0.312
$V^{(3)}_{lphaeta}$	Ni	Co	Fe	Cr
Ni	0.002	0.090	0.008	-0.100
Co	0.090	0.053	-0.009	-0.135
Fe	0.008	-0.009	-0.002	0.003
Cr	-0.100	-0.135	0.003	0.232



Atomistic Simulations

NiCoFeCr





Atomistic Simulations

NiCoFeMnCr



Ab Initio Theory

Expand free energy around homogeneous reference state:

$$\Omega(\{\bar{c}_{i\alpha}\}) = \Omega(\{c_{\alpha}\}) + \sum_{i\alpha} \frac{\partial\Omega}{\partial\bar{c}_{i\alpha}} \Big|_{\{c_{\alpha}\}} \Delta\bar{c}_{i\alpha} + \frac{1}{2} \sum_{i\alpha;j\alpha'} \frac{\partial^{2}\Omega}{\partial\bar{c}_{i\alpha}\partial\bar{c}_{j\alpha'}} \Big|_{\{c_{\alpha}\}} \Delta\bar{c}_{i\alpha}\Delta\bar{c}_{j\alpha'} + \dots$$
(7)

Ab Initio Theory

Important bit is second-order term.

$$\delta\Omega = \frac{1}{2} \sum_{i,j} \sum_{\alpha,\alpha'} \Delta \bar{c}_{i\alpha} [\beta^{-1} C_{\alpha,\alpha'}^{-1} - S_{i\alpha,j\alpha'}^{(2)}] \Delta \bar{c}_{j\alpha'}, \tag{8}$$

where
$$-rac{\partial^2 \langle \Omega_{
m el}
angle_0}{\partial ar{c}_{ilpha \partial} ar{c}_{jlpha'}} \equiv S^{(2)}_{ilpha;jlpha'}$$
, and $C^{-1}_{lpha lpha'} = rac{\delta_{lpha lpha'}}{c_{lpha}}$.

Ab Initio Theory

► Or, in reciprocal space,

$$\delta\Omega = \frac{1}{2} \sum_{\mathbf{k}} \sum_{\alpha,\alpha'} \Delta \bar{c}_{\alpha}(\mathbf{k}) [\beta^{-1} C_{\alpha\alpha'}^{-1} - S_{\alpha\alpha'}^{(2)}(\mathbf{k})] \Delta \bar{c}_{\alpha'}(\mathbf{k}).$$
(9)