

Short Range Order in Refractory High-Entropy Alloys: First Principles Theory and Atomistic Modelling

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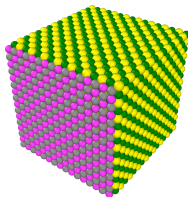
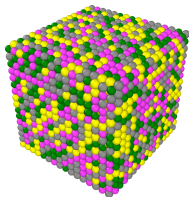
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High Entropy Alloys

- ▶ High Entropy Alloy (HEA): *multiple* principle components.
- ▶ Solid solution stabilised by large “entropy of mixing”

$$TS = -k_B TN \sum_{\alpha} c_{\alpha} \log c_{\alpha}. \quad (1)$$



- ▶ At what temperature will order emerge? What is the nature of order? Is order beneficial or detrimental?

Refractory HEAs

- ▶ Prototypical BCC HEAs.
- ▶ Typically V, Nb, Mo, Ta, W.

Periodic Table of the Elements

The image shows a standard periodic table of elements. A red rectangular box highlights a vertical column of five elements: Vanadium (V), Niobium (Nb), Molybdenum (Mo), Tantalum (Ta), and Tungsten (W). These elements are located in the transition metal block, specifically in the d-block. The table includes element symbols, atomic numbers, and names. Below the main table, there are two rows of lanthanide and actinide elements. At the bottom, there is a legend for element categories: Alkali, Alkaline Earth, Transition Metal, Lanthanide, Actinide, Metalloid, Nonmetal, Halogen, Noble Gas, and Metal.

- ▶ Apply the modelling techniques used successfully on Ni-based HEAs¹ to refractory systems.

¹C. D. Woodgate, J. B. Staunton, Phys. Rev. B **105** 115124 (2022)

Our Description

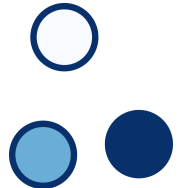
- ▶ On lattice
- ▶ Specify configuration by $\{\xi_{i\alpha}\}$
- ▶ Interested in the *average* value of these:

$$c_{i\alpha} \equiv \langle \xi_{i\alpha} \rangle$$

- ▶ Perturb high temperature, homogeneous state

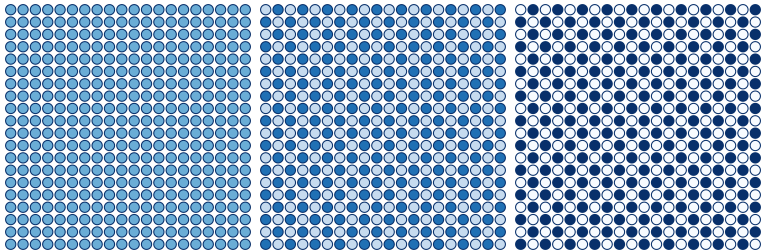
$$c_{i\alpha} = c_{\alpha} + \delta c_{i\alpha}$$

and see what favourable correlations are.



Concentration Waves

$$c_{i\alpha} = c_\alpha + \eta_\alpha \frac{1}{2} \left(e^{i\mathbf{q}\cdot\mathbf{R}_i} + e^{-i\mathbf{q}\cdot\mathbf{R}_i} \right), \quad \mathbf{q} = \left(\frac{1}{2}, \frac{1}{2} \right)$$



$$\eta = (0, 0)$$

$$\eta = (0.25, -0.25)$$

$$\eta = (0.5, -0.5)$$

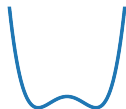
Energetics

- ▶ Evaluate cost of fluctuations *ab initio* via DFT, using KKR-CPA and a Landau-type theory^{ab}.
- ▶ Mean-field free energy:

$$G = \beta^{-1} \sum_{i\alpha} c_{i\alpha} \log c_{i\alpha} - \sum_{i\alpha}' \nu_{i\alpha} c_{i\alpha} + \langle \Omega_{\text{el}} \rangle_0 [\{c_{i\alpha}\}]$$

- ▶ Important quantities:

$$S_{i\alpha; j\alpha'}^{(2)} \equiv \frac{\partial^2 \langle \Omega_{\text{el}} \rangle_0}{\partial c_{i\alpha} \partial c_{j\alpha'}}, \quad \frac{\partial^2 G}{\partial c_{i\alpha} \partial c_{j\alpha'}} \rightsquigarrow \Psi_{\alpha\alpha'}^{-1}(\mathbf{q})$$



^aKhan, Staunton, Stocks, Phys. Rev. B **93** 054206 (2016)

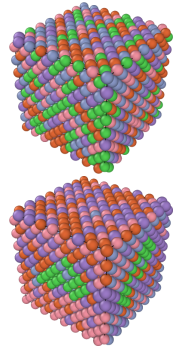
^bGyorffy, Stocks, Phys. Rev. Lett. **50**, 374 (1983)

Energetics

- ▶ Bragg-Williams Hamiltonian for atomistic modelling:

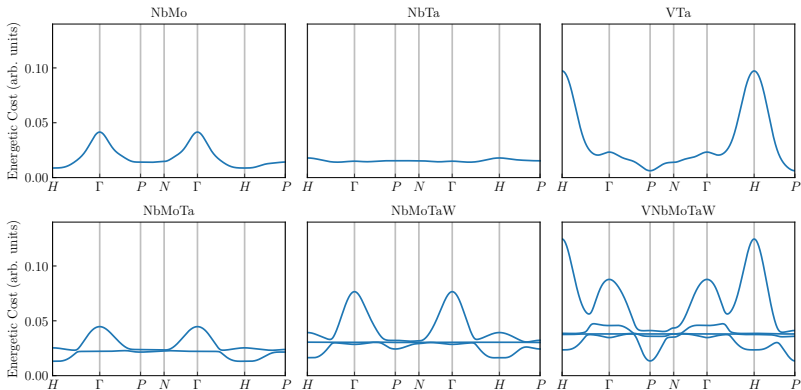
$$H(\{\xi_{i\alpha}\}) = \frac{1}{2} \sum_{\substack{i\alpha \\ j\alpha'}} V_{i\alpha;j\alpha'} \xi_{i\alpha} \xi_{j\beta} + \sum_{i\alpha} \nu_{i\alpha} \xi_{i\alpha}$$

- ▶ If H as above, $V_{i\alpha;j\alpha'}$ are *exactly* $-S_{i\alpha;j\alpha'}^{(2)}$.



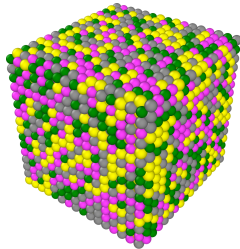
Linear Response

Eigenvalues of $\Psi_{\alpha\alpha'}^{-1}(\mathbf{q})$ around IBZ at 1200K.

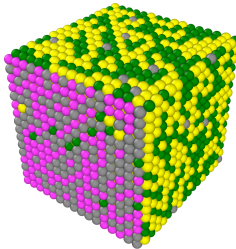


Atomistic Simulations

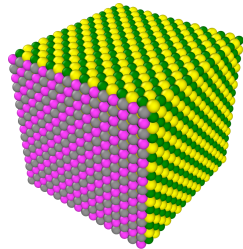
NbMoTaW



1200K

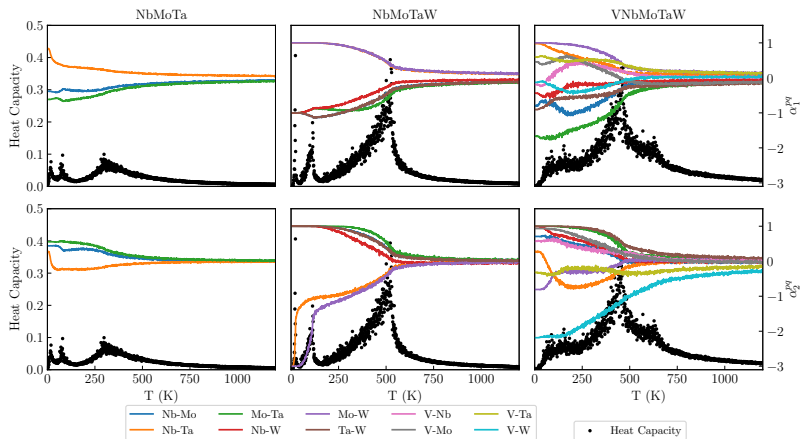


300K



10K

Atomistic Simulations



Summary

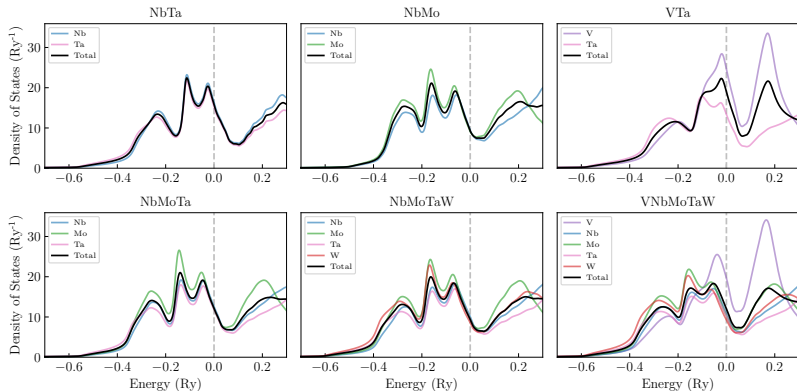
- ▶ B2-like ordering successfully predicted for NbMoTaW.
- ▶ SRO in VNbMoTaW dominated by V.
- ▶ Insight into physical origins of order.
- ▶ **Method obtains results comparable with conventional techniques for a fraction of the computational cost.**

Acknowledgements

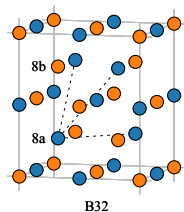
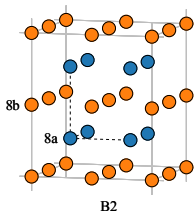
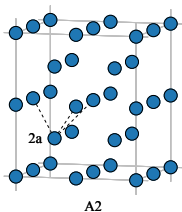
- ▶ I am funded by a studentship in the UK EPSRC-supported CDT in Modelling of Heterogeneous Systems at the University of Warwick, UK.
warwick.ac.uk/hetsys



Density of States



BCC Ordered Structures



Linear Response

Material	$T_{us}(K)$	k_{us}	δc_1	δc_2	δc_3	δc_4	δc_5
NbMoTa	511	$\{0, 0, 1\}$	-0.406	0.816	-0.410		
NbMoTaW	559	$\{0, 0, 1\}$	-0.383	0.594	-0.595	0.383	
VNbMoTaW	742	$\{\frac{1}{2}, \frac{1}{2}, \frac{1}{2}\}$	-0.824	0.012	0.085	0.252	0.500

Fitting to Pairwise Interaction

We approximate the true interaction by fitting to a Bragg-Williams-like Hamiltonian. That is, we approximate the *ab initio* data by $V_{\alpha\beta}(\mathbf{q})$, where

$$V_{\alpha\beta}(\mathbf{q}) = V_{\alpha\beta}^0 + V_{\alpha\beta}^1 \left(\sum_{\{\mathbf{R}_i\}_1} e^{i\mathbf{q}\cdot\mathbf{R}_i} \right) + \dots + V_{\alpha\beta}^N \left(\sum_{\{\mathbf{R}_i\}_N} e^{i\mathbf{q}\cdot\mathbf{R}_i} \right), \quad (2)$$

a constant plus a nearest-neighbour term, next-nearest neighbour term, etc.

Ab Initio Theory

- ▶ Expand free energy around homogeneous reference state:

$$\begin{aligned}\Omega(\{\bar{c}_{i\alpha}\}) &= \Omega(\{c_\alpha\}) + \sum_{i\alpha} \left. \frac{\partial \Omega}{\partial \bar{c}_{i\alpha}} \right|_{\{c_\alpha\}} \Delta \bar{c}_{i\alpha} \\ &+ \frac{1}{2} \sum_{i\alpha; j\alpha'} \left. \frac{\partial^2 \Omega}{\partial \bar{c}_{i\alpha} \partial \bar{c}_{j\alpha'}} \right|_{\{c_\alpha\}} \Delta \bar{c}_{i\alpha} \Delta \bar{c}_{j\alpha'} + \dots\end{aligned}\quad (3)$$

Ab Initio Theory

- ▶ Important bit is second-order term.

$$\delta\Omega = \frac{1}{2} \sum_{i,j} \sum_{\alpha,\alpha'} \Delta\bar{c}_{i\alpha} [\beta^{-1} C_{\alpha,\alpha'}^{-1} - S_{i\alpha;j\alpha'}^{(2)}] \Delta\bar{c}_{j\alpha'}, \quad (4)$$

where $-\frac{\partial^2 \langle \Omega_{el} \rangle_0}{\partial \bar{c}_{i\alpha} \partial \bar{c}_{j\alpha'}} \equiv S_{i\alpha;j\alpha'}^{(2)}$, and $C_{\alpha\alpha'}^{-1} = \frac{\delta_{\alpha\alpha'}}{c_\alpha}$.

Ab Initio Theory

- ▶ Or, in reciprocal space,

$$\delta\Omega = \frac{1}{2} \sum_{\mathbf{k}} \sum_{\alpha, \alpha'} \Delta\bar{c}_{\alpha}(\mathbf{k}) [\beta^{-1} C_{\alpha\alpha'}^{-1} - S_{\alpha\alpha'}^{(2)}(\mathbf{k})] \Delta\bar{c}_{\alpha'}(\mathbf{k}). \quad (5)$$