A Quantitative Central Limit Theorem for Simple Symmetric Exclusion Process

Vitalii Konarovskyi

University of Hamburg and Institute of Mathematics of NAS of Ukraine

Joint Kyiv-Warwick Stochastic Analysis seminar

joint work with Benjamin Gess

National Academy of Sciences of Ukraine **INSTITUTE OF MATHEMATICS**

Vitalii Konarovskyi (University of Hamburg and Institutive CLT for SSEP December 10, 2024 1/22

Simple symmetric exclusion process

On the d-dim discrete torus

$$
\mathbb{T}_n^d := \left\{ \frac{k}{n} : k \in \mathbb{Z}_n^d := \{-m, \ldots, m\}^d \right\} \subset \mathbb{T}^d = \left(\mathbb{R} / \{ \mathbb{Z} - 1/2 \} \right)^d
$$

for $n = 2m + 1$ we consider a **Simple Symmetric Exclusion Process** (SSEP)

∍

State space and generator

Particle configuration $\eta \in \{0,1\}^{\mathbb{T}_n^d}.$

 $\eta(x) = 0 \Leftrightarrow$ side x is empty $\eta(x) = 1 \Leftrightarrow$ side x is occupied

η

$$
\eta^{x \leftrightarrow y}(z) = \begin{cases} \eta(z), & z \neq x, y, \\ \eta(y), & z = x, \\ \eta(x), & z = y, \end{cases}
$$

 298

∍

State space and generator

Particle configuration $\eta \in \{0,1\}^{\mathbb{T}_n^d}.$

 $\eta(x) = 0 \Leftrightarrow$ side x is empty $\eta(x) = 1 \Leftrightarrow$ side x is occupied

$$
\mathcal{G}_n^{EP}F(\eta) := \frac{n^2}{2} \sum_{j=1}^d \sum_{x \in \mathbb{T}_n} \left[F(\eta^{x \leftrightarrow x + e_j}) - F(\eta) \right]
$$
 [Kipnis, Landim '99]

SSEP is already parabolically rescaled: space $\sim \frac{1}{n}$ $\sim \frac{1}{n}$ $\sim \frac{1}{n}$ time $\sim n^2$.

Let $\eta_t^n, t \geq 0$, be a SSEP.

D.

4 D F

造

Let $\eta_t^n, t \geq 0$, be a SSEP.

Equilibrium SSEP:

For fixed $\rho_0\in[0,1]$, the measure $B(\rho_0)^{\otimes \mathbb{T}_n^d}$ is invariant for SSEP η_t^n : If $\eta_0^n(x)\sim B(\rho_0),\,x\in\mathbb{T}_n^d$, are independent then $\eta_t^n(x)\sim B(\rho_0),\,x\in\mathbb{T}_n^d$, are independent. \rightsquigarrow $\mathbb{E}\eta_t^n(x) = \rho_0, \text{ Var } \eta_t^n = \rho_0(1-\rho_0).$

Let $\eta_t^n, t \geq 0$, be a SSEP.

Equilibrium SSEP:

For fixed $\rho_0\in[0,1]$, the measure $B(\rho_0)^{\otimes \mathbb{T}_n^d}$ is invariant for SSEP η_t^n : If $\eta_0^n(x)\sim B(\rho_0),\,x\in\mathbb{T}_n^d$, are independent then $\eta_t^n(x)\sim B(\rho_0),\,x\in\mathbb{T}_n^d$, are independent. \rightsquigarrow $\mathbb{E}\eta_t^n(x) = \rho_0, \text{ Var } \eta_t^n = \rho_0(1-\rho_0).$

Non-Equilibrium SSEP:

We consider $\rho_0: \mathbb{T}^d \to [0,1]$ and let $\eta_0^n(\mathsf{x}) \sim B(\rho_0(\mathsf{x})), \, \mathsf{x} \in \mathbb{T}_n^d,$ are independent.

Let $\eta_t^n, t \geq 0$, be a SSEP.

Equilibrium SSEP:

For fixed $\rho_0\in[0,1]$, the measure $B(\rho_0)^{\otimes \mathbb{T}_n^d}$ is invariant for SSEP η_t^n : If $\eta_0^n(x)\sim B(\rho_0),\,x\in\mathbb{T}_n^d$, are independent then $\eta_t^n(x)\sim B(\rho_0),\,x\in\mathbb{T}_n^d$, are independent. \rightsquigarrow $\mathbb{E}\eta_t^n(x) = \rho_0, \text{ Var } \eta_t^n = \rho_0(1-\rho_0).$

Non-Equilibrium SSEP:

We consider $\rho_0: \mathbb{T}^d \to [0,1]$ and let $\eta_0^n(\mathsf{x}) \sim B(\rho_0(\mathsf{x})), \, \mathsf{x} \in \mathbb{T}_n^d,$ are independent. What is $\mathbb{E} \eta_t^n(x)$?

 QQ

Set $\rho_t^n(x) := \mathbb{E} \eta_t^n(x)$.

造

ヨメ イヨメ

◂**◻▸ ◂◚▸**

Set $\rho_t^n(x) := \mathbb{E} \eta_t^n(x)$.

Then

$$
d\rho_t^n(x) = \mathbb{E}\mathcal{G}_n^{EP}\eta_t^n(x)dt = \mathbb{E}\frac{n^2}{2}\sum_{j=1}^d\sum_{y\in\mathbb{T}_n}\left[\eta^{y\leftrightarrow y+e_j}(x) - \eta(x)\right]dt
$$

化重复化重复

← ロ → → ← 何 →

Set $\rho_t^n(x) := \mathbb{E} \eta_t^n(x)$.

Then

$$
d\rho_t^n(x) = \mathbb{E}\mathcal{G}_n^{EP}\eta_t^n(x)dt = \mathbb{E}\frac{n^2}{2}\sum_{j=1}^d\sum_{y\in\mathbb{T}_n}\left[\eta^{y\leftrightarrow y+e_j}(x) - \eta(x)\right]dt
$$

$$
= \frac{n^2}{2}\sum_{j=1}^d\mathbb{E}\left(\eta_t^n(x+e_j) + \eta_t^n(x-e_j) - 2\eta(x)\right)dt
$$

4 0 8

- ← 冊 →

 \rightarrow \rightarrow \rightarrow

э

Set $\rho_t^n(x) := \mathbb{E} \eta_t^n(x)$.

Then

$$
d\rho_t^n(x) = \mathbb{E}\mathcal{G}_n^{EP}\eta_t^n(x)dt = \mathbb{E}\frac{n^2}{2}\sum_{j=1}^d\sum_{y\in\mathbb{T}_n}\left[\eta^{y\leftrightarrow y+e_j}(x) - \eta(x)\right]dt
$$

$$
= \frac{n^2}{2}\sum_{j=1}^d\mathbb{E}\left(\eta_t^n(x+e_j) + \eta_t^n(x-e_j) - 2\eta(x)\right)dt
$$

$$
= \frac{1}{2}\Delta_n\rho_t^n(x)dt.
$$

 \rightarrow \rightarrow \rightarrow

э.

4 0 8

- ← 冊 →

重

Set $\rho_t^n(x) := \mathbb{E} \eta_t^n(x)$.

Then

$$
d\rho_t^n(x) = \mathbb{E}\mathcal{G}_n^{EP}\eta_t^n(x)dt = \mathbb{E}\frac{n^2}{2}\sum_{j=1}^d\sum_{y\in\mathbb{T}_n}\left[\eta^{y\leftrightarrow y+e_j}(x) - \eta(x)\right]dt
$$

$$
= \frac{n^2}{2}\sum_{j=1}^d\mathbb{E}\left(\eta_t^n(x+e_j) + \eta_t^n(x-e_j) - 2\eta(x)\right)dt
$$

$$
= \frac{1}{2}\Delta_n\rho_t^n(x)dt.
$$

In particular, the empirical distribution

$$
\tilde{\rho}_t^n := \frac{1}{n^d} \sum_{x \in \mathbb{T}_n^d} \rho_t(x) \delta_x,
$$

solves

$$
d\langle \varphi,\tilde{\rho}_t^n\rangle=\frac{1}{2}\langle \Delta_n\varphi,\tilde{\rho}_t^n\rangle dt
$$

Vitalii Konarovskyi (University of Hamburg and Institutive CLT for SSEP December 10, 2024 5/22

造

 \triangleright \rightarrow \exists \rightarrow

◂**◻▸ ◂⊓▸**

Set $\rho_t^n(x) := \mathbb{E} \eta_t^n(x)$.

Then

$$
d\rho_t^n(x) = \mathbb{E}\mathcal{G}_n^{EP}\eta_t^n(x)dt = \mathbb{E}\frac{n^2}{2}\sum_{j=1}^d\sum_{y\in\mathbb{T}_n}\left[\eta^{y\leftrightarrow y+e_j}(x) - \eta(x)\right]dt
$$

$$
= \frac{n^2}{2}\sum_{j=1}^d\mathbb{E}\left(\eta_t^n(x+e_j) + \eta_t^n(x-e_j) - 2\eta(x)\right)dt
$$

$$
= \frac{1}{2}\Delta_n\rho_t^n(x)dt.
$$

In particular, the empirical distribution

$$
\tilde{\rho}_t^n := \frac{1}{n^d} \sum_{x \in \mathbb{T}_n^d} \rho_t(x) \delta_x,
$$

solves

$$
d\langle \varphi, \tilde{\rho}_t^n \rangle = \frac{1}{2} \langle \Delta_n \varphi, \tilde{\rho}_t^n \rangle dt \stackrel{n \to \infty}{\to} d\rho_t^{\infty} = \frac{1}{2} \Delta \rho_t^{\infty} dt, \quad \rho_0^{\infty} = \rho_0
$$

Vitalii Konarovskyi (University of Hamburg a**nd Institute OLT for SSEP** December 10, 2024 5/22

重

Law of large numbers

Theorem [see e.g. in Kipnis, Landim '99]

Let $\rho_0: \mathbb{T}^d \to [0,1]$ be an initial density profile and $\eta_0^n(x) \sim B(\rho_0(x))$ be independent. Then

$$
\tilde{\eta}_t^n := \frac{1}{n^d} \sum_{\mathbf{x} \in \mathbb{T}_n^d} \eta_t(\mathbf{x}) \delta_{\mathbf{x}}
$$

 $\textsf{converges}$ in probability to $\rho_t^\infty(x)$ dx, where $\rho_t^\infty := P_t^{\textsf{HE}} \rho_0$ solves

$$
d\rho_t^{\infty}=\frac{1}{2}\Delta\rho_t^{\infty}dt, \quad \rho_0^{\infty}=\rho_0.
$$

Law of large numbers

Theorem [see e.g. in Kipnis, Landim '99]

Let $\rho_0: \mathbb{T}^d \to [0,1]$ be an initial density profile and $\eta_0^n(x) \sim B(\rho_0(x))$ be independent. Then

$$
\tilde{\eta}_t^n := \frac{1}{n^d} \sum_{\mathbf{x} \in \mathbb{T}_n^d} \eta_t(\mathbf{x}) \delta_{\mathbf{x}}
$$

 $\textsf{converges}$ in probability to $\rho_t^\infty(x)$ dx, where $\rho_t^\infty := P_t^{\textsf{HE}} \rho_0$ solves

$$
d\rho_t^{\infty}=\frac{1}{2}\Delta\rho_t^{\infty}dt, \quad \rho_0^{\infty}=\rho_0.
$$

$$
G_n^{EP} f(\langle \varphi, \tilde{\eta} \rangle) := \frac{n^2}{2} \sum_{j=1}^d \sum_{x \in \mathbb{T}_n} \left[f(\langle \varphi, \tilde{\eta}^{x \leftrightarrow x + \epsilon_j} \rangle) - f(\langle \varphi, \tilde{\eta} \rangle) \right]
$$

$$
= \frac{1}{2} f'(\langle \varphi, \tilde{\eta} \rangle) \langle \Delta_n \varphi, \tilde{\eta} \rangle + \frac{1}{4n^2} f''(\langle \varphi, \tilde{\eta} \rangle) \left\langle \left| \partial_{n,j} \varphi \right|^2, \tau_j \tilde{\eta} + \tilde{\eta} - 2 \tilde{\eta} \tilde{\tau_j} \tilde{\eta} \right\rangle + \dots,
$$

where $\tau_j \eta(x) := \eta(x + e_j)$.

 QQ

We now consider the fluctuations of the SSEP around its mean:

 $\zeta_t^n(x) := n^{\frac{d}{2}} (\eta_t^n(x) - \rho_t^n(x))$.

∍

We now consider the fluctuations of the SSEP around its mean:

 $\zeta_t^n(x) := n^{\frac{d}{2}} (\eta_t^n(x) - \rho_t^n(x))$.

The generator of

$$
\tilde{\zeta}_t^n := \frac{1}{n^d} \sum_{x \in \mathbb{T}_n^d} \zeta_t(x) \delta_x
$$

can be expanded as follows

$$
\mathcal{G}_{n}^{FF} f\left(\langle\varphi,\tilde{\zeta}\rangle\right) = \frac{1}{2} f'\left(\langle\varphi,\tilde{\zeta}\rangle\right) \langle\Delta_{n}\varphi,\tilde{\zeta}\rangle + \frac{n^{d}}{4n^{d}} f''\left(\langle\varphi,\tilde{\zeta}\rangle\right) \langle\left|\partial_{n,j}\varphi\right|^{2}, \tau_{j}\tilde{\eta} + \tilde{\eta} - 2\tilde{\eta}\tilde{\tau_{j}}\tilde{\eta}\rangle + O\left(1/n^{\frac{d}{2}+1}\right)
$$

∍

 QQ

We now consider the fluctuations of the SSEP around its mean:

 $\zeta_t^n(x) := n^{\frac{d}{2}} (\eta_t^n(x) - \rho_t^n(x))$.

The generator of

$$
\tilde{\zeta}^n_t:=\frac{1}{n^d}\sum_{\mathbf{x}\in\mathbb{T}_n^d}\zeta_t(\mathbf{x})\delta_{\mathbf{x}}
$$

can be expanded as follows

$$
\mathcal{G}_{n}^{FF} f\left(\langle\varphi,\tilde{\zeta}\rangle\right) = \frac{1}{2} f'\left(\langle\varphi,\tilde{\zeta}\rangle\right) \langle\Delta_{n}\varphi,\tilde{\zeta}\rangle + \frac{n^{d}}{4n^{d}} f''\left(\langle\varphi,\tilde{\zeta}\rangle\right) \langle\left|\partial_{n,j}\varphi\right|^{2}, \tau_{j}\tilde{\eta} + \tilde{\eta} - 2\tilde{\eta}\tilde{\tau_{j}}\tilde{\eta}\rangle + O\left(1/n^{\frac{d}{2}+1}\right)
$$

Again

$$
d\langle \varphi, \tilde{\zeta}_t^n \rangle = \frac{1}{2} \langle \Delta_n \varphi, \tilde{\zeta}_t^n \rangle dt + \text{mart.}
$$

$$
d\langle \text{mart.} \rangle_t = \frac{1}{2} \left\langle |\partial_{n,j} \varphi|^2, \tau_j \tilde{\eta}_t^n + \tilde{\eta}_t^n - 2 \tilde{\eta}_t^n \tau_j \eta_t^n \right\rangle dt
$$

Vitalii Konarovskyi (University of Hamburg and Institutive CLT for SSEP Becember 10, 2024 7/22

We now consider the fluctuations of the SSEP around its mean:

 $\zeta_t^n(x) := n^{\frac{d}{2}} (\eta_t^n(x) - \rho_t^n(x))$.

The generator of

$$
\tilde{\zeta}^n_t:=\frac{1}{n^d}\sum_{\mathbf{x}\in\mathbb{T}_n^d}\zeta_t(\mathbf{x})\delta_{\mathbf{x}}
$$

can be expanded as follows

$$
\mathcal{G}_{n}^{FF} f\left(\langle\varphi,\tilde{\zeta}\rangle\right) = \frac{1}{2} f'\left(\langle\varphi,\tilde{\zeta}\rangle\right) \langle\Delta_{n}\varphi,\tilde{\zeta}\rangle + \frac{n^{d}}{4n^{d}} f''\left(\langle\varphi,\tilde{\zeta}\rangle\right) \langle\left|\partial_{n,j}\varphi\right|^{2}, \tau_{j}\tilde{\eta} + \tilde{\eta} - 2\tilde{\eta}\tilde{\tau_{j}}\tilde{\eta}\rangle + O\left(1/n^{\frac{d}{2}+1}\right)
$$

Again

$$
d\langle \varphi, \tilde{\zeta}^n_t \rangle = \frac{1}{2} \langle \Delta_n \varphi, \tilde{\zeta}^n_t \rangle dt + \text{mart.} \to d\langle \varphi, \zeta^{\infty}_t \rangle = \frac{1}{2} \langle \Delta \varphi, \zeta^{\infty}_t \rangle dt + \text{mart.}
$$

$$
d\langle \text{mart.} \rangle_t = \frac{1}{2} \left\langle \left| \partial_{n,j} \varphi \right|^2, \tau_j \tilde{\eta}^n_t + \tilde{\eta}^n_t - 2 \tilde{\eta}^n_t \tau_j \tilde{\eta}^n_t \right\rangle dt \to d\langle \text{mart.} \rangle = \langle \Delta \varphi, \rho_t^{\infty} - \rho_t^{\infty} \rho_t^{\infty} \rangle dt
$$

Vitalii Konarovskyi (University of Hamburg and Institutive CLT for SSEP Becember 10, 2024 7/22

Central limit theorem

Theorem 2 [Galves, Kipnis, Spohn; Ravishankar '90]

Let the initial density profile ρ_0 be smooth. Then the density fluctuation field

$$
\tilde{\zeta}^n_t := \frac{1}{n^d} \sum_{x \in \mathbb{T}_n^d} \zeta_t(x) \delta_x
$$

converges in $D([0, T], \mathcal{D}')$ to the generalized Ornstein-Uhlenbeck process that solves the linear SPDE

$$
d\zeta_t^{\infty} = \frac{1}{2}\Delta \zeta_t^{\infty} dt + \nabla \cdot \left(\sqrt{\rho_t^{\infty} (1 - \rho_t^{\infty})} dW_t\right)
$$

with the centered Gaussian initial condition such that

$$
\mathbb{E}\left[\left\langle \zeta_0^\infty, \varphi \right\rangle^2\right] = \left\langle \rho_0(1-\rho_0)\varphi, \varphi \right\rangle
$$

Vitalii Konarovskyi (University of Hamburg and Institutive CLT for SSEP Becember 10, 2024 8/22

• Consider the semimartingale $\langle \varphi, \tilde{\zeta}_t \rangle$ and its quadratic variation.

$$
d\langle \varphi, \tilde{\zeta}_t^n \rangle = \frac{1}{2} \langle \Delta_n \varphi, \tilde{\zeta}_t^n \rangle dt + \text{mart.}
$$

$$
d\langle \text{mart.} \rangle_t = \frac{1}{2} \left\langle \left| \partial_{n,j} \varphi \right|^2, \tau_j \tilde{\eta}_t^n + \tilde{\eta}_t^n - 2 \tilde{\eta}_t^n \tau_j \eta_t^n \right\rangle dt
$$

э

4 D F

重

• Consider the semimartingale $\langle \varphi, \tilde{\zeta}_t \rangle$ and its quadratic variation.

$$
d\langle \varphi, \tilde{\zeta}_t^n \rangle = \frac{1}{2} \langle \Delta_n \varphi, \tilde{\zeta}_t^n \rangle dt + \text{mart.}
$$

$$
d\langle \text{mart.} \rangle_t = \frac{1}{2} \left\langle \left| \partial_{n,j} \varphi \right|^2, \tau_j \tilde{\eta}_t^n + \tilde{\eta}_t^n - 2 \tilde{\eta}_t^n \tilde{\gamma}_j \eta_t^n \right\rangle dt
$$

 \bullet Convergence of $\langle \text{mart.} \rangle_t$:

$$
\eta_t^n \tau_j \eta_t^n = \rho_t^n \tau_j \rho_t^n + \frac{1}{n^{d/2}} \left(\rho_t^n \tau_j \zeta_t^n + \zeta_t^n \tau_j \rho_t^n \right) + \frac{1}{n^d} \zeta_t^n \tau_j \zeta_t^n,
$$

where $\eta^n_t = \rho^n_t + \frac{1}{(2n+1)^{d/2}} \zeta^n_t$

∍

• Consider the semimartingale $\langle \varphi, \tilde{\zeta}_t \rangle$ and its quadratic variation.

$$
d\langle \varphi, \tilde{\zeta}_t^n \rangle = \frac{1}{2} \langle \Delta_n \varphi, \tilde{\zeta}_t^n \rangle dt + \text{mart.}
$$

$$
d\langle \text{mart.} \rangle_t = \frac{1}{2} \left\langle \left| \partial_{n,j} \varphi \right|^2, \tau_j \tilde{\eta}_t^n + \tilde{\eta}_t^n - 2 \tilde{\eta}_t^n \tilde{\gamma}_j \eta_t^n \right\rangle dt
$$

O Convergence of $\langle \text{mart.} \rangle_t$:

$$
\eta_t^n \tau_j \eta_t^n = \rho_t^n \tau_j \rho_t^n + \frac{1}{n^{d/2}} \left(\rho_t^n \tau_j \zeta_t^n + \zeta_t^n \tau_j \rho_t^n \right) + \frac{1}{n^d} \zeta_t^n \tau_j \zeta_t^n,
$$

where $\eta^n_t = \rho^n_t + \frac{1}{(2n+1)^{d/2}} \zeta^n_t$

 $\textsf{Control of } \mathbb{E}\left[\frac{1}{p^d}\zeta_t^n(\textsf{x})\tau_j\zeta_t^n(\textsf{x})\right] = \mathbb{E}\left[\left(\eta_t^n(\textsf{x})-\rho_t^n(\textsf{x})\right)\left(\eta_t^n(\textsf{x}+\textsf{e}_j)-\rho_t^n(\textsf{x}+\textsf{e}_j)\right)\right] \rightarrow 0$ [Ferrari et al. '91]

• Consider the semimartingale $\langle \varphi, \tilde{\zeta}_t \rangle$ and its quadratic variation.

$$
d\langle \varphi, \tilde{\zeta}_t^n \rangle = \frac{1}{2} \langle \Delta_n \varphi, \tilde{\zeta}_t^n \rangle dt + \text{mart.}
$$

$$
d\langle \text{mart.} \rangle_t = \frac{1}{2} \left\langle |\partial_{n,j} \varphi|^2, \tau_j \tilde{\eta}_t^n + \tilde{\eta}_t^n - 2 \tilde{\eta}_t^n \tau_j \eta_t^n \right\rangle dt
$$

O Convergence of $\langle \text{mart.} \rangle_t$:

$$
\eta_t^n \tau_j \eta_t^n = \rho_t^n \tau_j \rho_t^n + \frac{1}{n^{d/2}} \left(\rho_t^n \tau_j \zeta_t^n + \zeta_t^n \tau_j \rho_t^n \right) + \frac{1}{n^d} \zeta_t^n \tau_j \zeta_t^n,
$$

where $\eta^n_t = \rho^n_t + \frac{1}{(2n+1)^{d/2}} \zeta^n_t$

- $\textsf{Control of } \mathbb{E}\left[\frac{1}{p^d}\zeta_t^n(\textsf{x})\tau_j\zeta_t^n(\textsf{x})\right] = \mathbb{E}\left[\left(\eta_t^n(\textsf{x})-\rho_t^n(\textsf{x})\right)\left(\eta_t^n(\textsf{x}+\textsf{e}_j)-\rho_t^n(\textsf{x}+\textsf{e}_j)\right)\right] \rightarrow 0$ [Ferrari et al. '91]
- Use of tightness argument and the Holley-Stroock theory [Holley, Stroock '79]

Our goal: Obtain the **rate of convergence** of

$$
\tilde{\zeta}_t^n-\zeta_t^\infty=O(?)
$$

造

重

D. \sim

(□) (_□

Our goal: Obtain the **rate of convergence** of

$$
\sup_{t\in[0,\,T]}\left|\mathbb{E} f\left(\langle\varphi,\tilde{\zeta}^{n}_{t}\rangle\right)-\mathbb{E} f\left(\langle\varphi,\zeta^{\infty}_{t}\rangle\right)\right|\rightarrow 0
$$

造

Þ

4 0 F → 母

Our goal: Obtain the **rate of convergence** of

$$
\sup_{t\in[0,\,T]}\left|\mathbb{E} f\left(\langle\varphi,\tilde{\zeta}^{n}_{t}\rangle\right)-\mathbb{E} f\left(\langle\varphi,\zeta^{\infty}_{t}\rangle\right)\right|\rightarrow 0
$$

Difficulty: The tightness argument and the Holley-Stroock theory do not give the rate of convergence.

э

重

4 0 F

Our goal: Obtain the **rate of convergence** of

$$
\sup_{t\in[0,\,T]}\left|\mathbb{E} f\left(\langle\varphi,\tilde{\zeta}^{n}_{t}\rangle\right)-\mathbb{E} f\left(\langle\varphi,\zeta^{\infty}_{t}\rangle\right)\right|\rightarrow 0
$$

Difficulty: The tightness argument and the Holley-Stroock theory do not give the rate of convergence.

Quantitative results:

• [Gess, Wu, Zhang '24]: Higher order fluctuation expansions for nonlinear SPDEs. (SPDEs defined on the same probability space)

∍

 QQQ

ミドマミド

4 0 F

Our goal: Obtain the **rate of convergence** of

$$
\sup_{t\in[0,\,T]}\left|\mathbb{E} f\left(\langle\varphi,\tilde{\zeta}^{n}_{t}\rangle\right)-\mathbb{E} f\left(\langle\varphi,\zeta^{\infty}_{t}\rangle\right)\right|\rightarrow 0
$$

Difficulty: The tightness argument and the Holley-Stroock theory do not give the rate of convergence.

Quantitative results:

- **•** [Gess, Wu, Zhang '24]: Higher order fluctuation expansions for nonlinear SPDEs. (SPDEs defined on the same probability space)
- [Cornalba, Fischer '23, Djurdjevac, Kremp, Perkowski '24]: **Higher order** approximation of Dean-Kawasaki equation (duality of approach, structure of noise is important)

 QQ

Our goal: Obtain the **rate of convergence** of

$$
\sup_{t\in[0,\,T]}\left|\mathbb{E} f\left(\langle\varphi,\tilde{\zeta}^{n}_{t}\rangle\right)-\mathbb{E} f\left(\langle\varphi,\zeta^{\infty}_{t}\rangle\right)\right|\rightarrow 0
$$

Difficulty: The tightness argument and the Holley-Stroock theory do not give the rate of convergence.

Quantitative results:

• [Gess, Wu, Zhang '24]: Higher order fluctuation expansions for nonlinear SPDEs. (SPDEs defined on the same probability space)

• [Cornalba, Fischer '23, Djurdjevac, Kremp, Perkowski '24]: **Higher order** approximation of Dean-Kawasaki equation (duality of approach, structure of noise is important)

[Chassagneux, Szpruch, Tse '22]: Weak quantitative propagation of chaos (mean field limit)

 QQ

イロト イ押ト イヨト イヨト

Our goal: Obtain the **rate of convergence** of

$$
\sup_{t\in[0,\,T]}\left|\mathbb{E} f\left(\langle\varphi,\tilde{\zeta}^{n}_{t}\rangle\right)-\mathbb{E} f\left(\langle\varphi,\zeta^{\infty}_{t}\rangle\right)\right|\rightarrow 0
$$

Difficulty: The tightness argument and the Holley-Stroock theory do not give the rate of convergence.

Quantitative results:

• [Gess, Wu, Zhang '24]: Higher order fluctuation expansions for nonlinear SPDEs. (SPDEs defined on the same probability space)

• [Cornalba, Fischer '23, Djurdjevac, Kremp, Perkowski '24]: Higher order approximation of Dean-Kawasaki equation (duality of approach, structure of noise is important)

- [Chassagneux, Szpruch, Tse '22]: Weak quantitative propagation of chaos (mean field limit)
- [Kolokoltsov '10] Central limit theorem for the Smoluchovski coagulation model (mean field limit, non-local Smoluchowski's coagulation equation)

 \bullet ...

э

 QQ

Main result

Theorem 3 [Gess, K. '24]

Let

- the initial density profile $\rho_0: \mathbb{T}^d \to [0,1]$ be smooth enough,
- η_t^n be SSEP with $\eta_0^n(x) \sim B(\rho_0(x))$ and independent,

$$
\bullet \ \rho_t^n = \mathbb{E} \eta_t^n, \ \zeta_t^n = n^{d/2} \big(\eta_t^n - \rho_t^n \big)
$$

 ζ_t^∞ solves $d\zeta_t^\infty=\frac{1}{2}\Delta\zeta_t^\infty dt+\nabla\cdot\left(\sqrt{\rho_t^\infty(1-\rho_t^\infty)}dW_t\right)$ with the centered $\mathsf{Gaussian}\;$ initial condition with $\mathbb{E}\left[\langle \zeta_0^\infty,\varphi\rangle^2\right]=\langle \rho_0(1-\rho_0)\varphi,\varphi\rangle$ Then for large enough $I \in \mathbb{N}$

$$
\sup_{t\in[0,T]}\left|\mathbb{E} f\left(\langle \vec{\varphi}, \tilde{\zeta}^n_t \rangle\right) - \mathbb{E} f\left(\langle \vec{\varphi}, \zeta^{\infty}_t \rangle\right)\right| \leq \frac{C}{n^{\frac{d}{2}\wedge 1}} \left\|f\right\|_{C^3_1} \left\|\vec{\varphi}\right\|_{C^1}
$$
\nfor all $n \geq 1$, $f \in C^3_b(\mathbb{R}^m)$ and $\vec{\varphi} \in \left(C^l(\mathbb{T}^d)\right)^m$.

Main result

Theorem 3 [Gess, K. '24]

Let

- the initial density profile $\rho_0: \mathbb{T}^d \to [0,1]$ be smooth enough,
- η_t^n be SSEP with $\eta_0^n(x) \sim B(\rho_0(x))$ and independent,

$$
\bullet \ \rho_t^n = \mathbb{E} \eta_t^n, \ \zeta_t^n = n^{d/2} \big(\eta_t^n - \rho_t^n \big)
$$

 ζ_t^∞ solves $d\zeta_t^\infty=\frac{1}{2}\Delta\zeta_t^\infty dt+\nabla\cdot\left(\sqrt{\rho_t^\infty(1-\rho_t^\infty)}dW_t\right)$ with the centered $\mathsf{Gaussian}\;$ initial condition with $\mathbb{E}\left[\langle \zeta_0^\infty,\varphi\rangle^2\right]=\langle \rho_0(1-\rho_0)\varphi,\varphi\rangle$ Then for large enough $I \in \mathbb{N}$

$$
\sup_{t\in[0,T]}\left|\mathbb{E} f\left(\langle \vec{\varphi}, \vec{\zeta}^n_t \rangle\right) - \mathbb{E} f\left(\langle \vec{\varphi}, \zeta^{\infty}_t \rangle\right)\right| \leq \frac{C}{n^{\frac{d}{2}\wedge 1}}\left\|f\right\|_{C^3_1} \|\vec{\varphi}\|_{C^1}
$$

$$
\text{ for all } n \geq 1, \ f \in \mathrm{C}^3_b(\mathbb{R}^m) \ \text{and } \vec{\varphi} \in \left(\mathrm{C}^{\prime}(\mathbb{T}^d)\right)^m.
$$

The rate $\frac{1}{\frac{d}{2}\wedge 1}$ is optimal: $\frac{1}{n}$ $\frac{1}{n}$ $\frac{1}{n}$ – lattice discretization error, $\frac{1}{n^{\frac{d}{2}}}$ – particle approximation error QQQ Vitalii Konarovskyi (University of Hamburg and Institutive CLT for SSEP December 10, 2024 11/22

Main tool

Idea of proof: Compare two (time-homogeneous) Markov processes X_t , Y_t taking values in the same state space and $X_0 = Y_0 = x$ using

$$
\mathbb{E}F(X_t) - \mathbb{E}F(Y_t) = \int_0^t P_s^X \left(\mathcal{G}^X - \mathcal{G}^Y \right) P_{t-s}^Y F(x) ds,
$$

=
$$
\int_0^t \mathbb{E} \left[\left(\mathcal{G}^X - \mathcal{G}^Y \right) P_{t-s}^Y F(X_s) \right] ds,
$$

[see e.g. Ethier, Kurtz '86]

4 0 8

э

Recall

$$
\mathbb{E} F(X_t) - \mathbb{E} F(Y_t) = \int_0^t \mathbb{E} \left[\left(\mathcal{G}^X - \mathcal{G}^Y \right) P_{t-s}^Y F(X_s) \right] ds,
$$

where $X_0 = Y_0 = x$.

 299

Þ

D

4 D F ∢母 造

Recall

$$
\mathbb{E} F(X_t) - \mathbb{E} F(Y_t) = \int_0^t \mathbb{E} \left[\left(\mathcal{G}^X - \mathcal{G}^Y \right) P_{t-s}^Y F(X_s) \right] ds,
$$

where $X_0 = Y_0 = x$.

We consider the Markov processes:

- particle means and fluctuation field: $(\tilde{\rho}^n_t, \tilde{\zeta}^n_t)$
- solution to heat equation and generalized OU process $(\rho_t^\infty,\zeta_t^\infty).$

 QQ

∍

Recall

$$
\mathbb{E} F(X_t) - \mathbb{E} F(Y_t) = \int_0^t \mathbb{E} \left[\left(\mathcal{G}^X - \mathcal{G}^Y \right) P_{t-s}^Y F(X_s) \right] ds,
$$

where $X_0 = Y_0 = x$.

We consider the Markov processes:

- particle means and fluctuation field: $(\tilde{\rho}^n_t, \tilde{\zeta}^n_t)$
- solution to heat equation and generalized OU process $(\rho_t^\infty,\zeta_t^\infty).$

The processes starts from different initial conditions!

∍

 QQ

Recall

$$
\mathbb{E} F(X_t) - \mathbb{E} F(Y_t) = \int_0^t \mathbb{E} \left[\left(\mathcal{G}^X - \mathcal{G}^Y \right) P_{t-s}^Y F(X_s) \right] ds,
$$

where $X_0 = Y_0 = x$.

We consider the Markov processes:

- particle means and fluctuation field: $(\tilde{\rho}^n_t, \tilde{\zeta}^n_t)$
- solution to heat equation and generalized OU process $(\rho_t^\infty,\zeta_t^\infty).$

The processes starts from different initial conditions!

We will compare:

- $(\tilde{\rho}^n_t, \tilde{\zeta}^n_t)$ an $(\rho^{\infty,n}_t, \zeta^{\infty,n}_t)$ [comparison of dynamics] where the generalized OU process $(\rho_t^{\infty,n}, \zeta_t^{\infty,n})$ started from $(\tilde{\rho}_0^n, \tilde{\zeta}_0^n)$;
- $(\rho_t^{\infty,n}, \zeta_t^{\infty,n})$ and $(\rho_t^{\infty}, \zeta_t^{\infty})$ [comparison of initial conditions] (both are defined by the same equation).

Vitalii Konarovskyi (University of Hamburg and Institutive CLT for SSEP December 10, 2024 13/22

 QQ

Generators

We start from the formal computation for cylindrical functions:

 $\mathcal{F}(\tilde{\rho}, \tilde{\zeta}) := \mathcal{F}\left(\langle \psi, \tilde{\rho} \rangle, \langle \varphi, \tilde{\zeta} \rangle\right)$

重

 \rightarrow \rightarrow \rightarrow

◂**◻▸ ◂⊓▸**

Generators

We start from the formal computation for cylindrical functions:

 $\mathcal{F}(\tilde{\rho}, \tilde{\zeta}) := \mathcal{F}\left(\langle \psi, \tilde{\rho} \rangle, \langle \varphi, \tilde{\zeta} \rangle\right)$

Using Taylor's formula, we get:

$$
G_n^{FF}F(\tilde{\rho},\tilde{\zeta}) = \frac{1}{2}\partial_1 f \langle \Delta_n \varphi, \tilde{\rho} \rangle + \frac{1}{2}\partial_2 f \langle \Delta_n \varphi, \tilde{\zeta} \rangle + \frac{1}{4}\partial_2^2 f \langle |\partial_{n,j}\varphi|^2, \tau_j \tilde{\eta} + \tilde{\eta} - 2\tilde{\eta} \tilde{\tau_j} \tilde{\eta} \rangle + O\left(1/n^{\frac{d}{2}+1}\right),
$$

where $\tilde{\zeta} = n^{d/2}(\tilde{\eta} - \tilde{\rho})$.

э

 298

 \rightarrow \rightarrow \rightarrow

4 D F

Generators

We start from the formal computation for cylindrical functions:

 $\mathcal{F}(\tilde{\rho}, \tilde{\zeta}) := \mathcal{F}\left(\langle \psi, \tilde{\rho} \rangle, \langle \varphi, \tilde{\zeta} \rangle\right)$

Using Taylor's formula, we get:

$$
G_n^{FF}F(\tilde{\rho},\tilde{\zeta}) = \frac{1}{2}\partial_1 f \langle \Delta_n \varphi, \tilde{\rho} \rangle + \frac{1}{2}\partial_2 f \langle \Delta_n \varphi, \tilde{\zeta} \rangle + \frac{1}{4}\partial_2^2 f \langle |\partial_{n,j}\varphi|^2, \tau_j \tilde{\eta} + \tilde{\eta} - 2\tilde{\eta}\tilde{\tau_j}\tilde{\eta} \rangle + O\left(1/n^{\frac{d}{2}+1}\right),
$$

where $\tilde{\zeta} = n^{d/2}(\tilde{\eta} - \tilde{\rho})$. Moreover,

$$
G^{OU}F(\rho^{\infty}, \zeta^{\infty}) = \frac{1}{2}\partial_1 f \langle \Delta \varphi, \rho^{\infty} \rangle + \frac{1}{2}\partial_2 f \langle \Delta \varphi, \zeta^{\infty} \rangle + \frac{1}{2}\partial_2^2 f \langle |\partial_j \varphi|^2, \rho^{\infty} - (\rho^{\infty})^2 \rangle.
$$

Vitalii Konarovskyi (University of Hamburg and Institutive CLT for SSEP Becember 10, 2024 14/22

4 0 F

 299

э

Generators have to be compared on $U:=P_{t-s}^{OU}f(\langle\psi,\cdot\rangle,\langle\varphi,\cdot\rangle)$, that is not cylindrical function.

⇝ We need to work with Frechet derivatives: e.g. D2U instead of *∂*2f · *φ*

∍

4 0 8

Generators have to be compared on $U:=P_{t-s}^{OU}f(\langle\psi,\cdot\rangle,\langle\varphi,\cdot\rangle)$, that is not cylindrical function.

⇝ We need to work with Frechet derivatives: e.g. D2U instead of *∂*2f · *φ*

We compare the generators on particle configurations $\tilde{\rho}^n_t$, $\tilde{\zeta}^n_t$, that are <u>empirica</u>l distributions. But U is not differentiable at $\tilde{\rho}^n_t$ because of the term $\sqrt{\rho(1-\rho)}$ in the SPDE

 \rightsquigarrow Probably, we have to use e.g. linear interpolation of ρ_t^n .

Generators have to be compared on $U:=P_{t-s}^{OU}f(\langle\psi,\cdot\rangle,\langle\varphi,\cdot\rangle)$, that is not cylindrical function.

⇝ We need to work with Frechet derivatives: e.g. D2U instead of *∂*2f · *φ*

We compare the generators on particle configurations $\tilde{\rho}^n_t$, $\tilde{\zeta}^n_t$, that are <u>empirica</u>l distributions. But U is not differentiable at $\tilde{\rho}^n_t$ because of the term $\sqrt{\rho(1-\rho)}$ in the SPDE

 \rightsquigarrow Probably, we have to use e.g. linear interpolation of ρ_t^n .

• The term $\eta\tau_j\eta$ has a part $\zeta\tau_j\rho$. But we do not have enough regularity in linear $\langle \text{interpolation of } \rho \text{ to control } \left\langle \widehat{\zeta_{Tj}\rho}, |\partial_{n,j}\varphi|^2 \right\rangle = \left\langle \tau_j \rho \left| \partial_{n,j}\varphi \right|^2, \widetilde{\zeta} \right\rangle$ \rightsquigarrow Interpolation has to be smooth enough

Generators have to be compared on $U:=P_{t-s}^{OU}f(\langle\psi,\cdot\rangle,\langle\varphi,\cdot\rangle)$, that is not cylindrical function.

⇝ We need to work with Frechet derivatives: e.g. D2U instead of *∂*2f · *φ*

We compare the generators on particle configurations $\tilde{\rho}^n_t$, $\tilde{\zeta}^n_t$, that are <u>empirica</u>l distributions. But U is not differentiable at $\tilde{\rho}^n_t$ because of the term $\sqrt{\rho(1-\rho)}$ in the SPDE

 \rightsquigarrow Probably, we have to use e.g. linear interpolation of ρ_t^n .

- **•** The term $\eta\tau_j\eta$ has a part $\zeta\tau_j\rho$. But we do not have enough regularity in linear $\langle \text{interpolation of } \rho \text{ to control } \left\langle \widehat{\zeta_{Tj}\rho}, |\partial_{n,j}\varphi|^2 \right\rangle = \left\langle \tau_j \rho \left| \partial_{n,j}\varphi \right|^2, \widetilde{\zeta} \right\rangle$ \rightsquigarrow Interpolation has to be smooth enough
- $\mathsf{Control} \text{ of } \mathbb{E}\left[f(\langle \varphi, \tilde{\zeta}^n_t \rangle) \left\langle \widetilde{\zeta^n_t \tau \zeta^n_t}, | \partial_{n,j} \varphi |^2 \right\rangle \right] \text{ via } \mathbb{E}\prod_{i=1}^m \left(\eta^n_t(\mathsf{x}_i) \rho^n_t(\mathsf{x}_i) \right) \text{ does not}$ give the optimal rate.

Replace $\tilde{\rho} = \frac{1}{n^d} \sum_{x \in \mathbb{T}_n^d} \rho(x) \delta_x$ and $\tilde{\zeta} = \frac{1}{n^d} \sum_{x \in \mathbb{T}_n^d} \zeta(x) \delta_x$ by a smooth interpolation.

Vitalii Konarovskyi (University of Hamburg and Institutive CLT for SSEP Becember 10, 2024 16/22

э

 QQQ

Replace $\tilde{\rho} = \frac{1}{n^d} \sum_{x \in \mathbb{T}_n^d} \rho(x) \delta_x$ and $\tilde{\zeta} = \frac{1}{n^d} \sum_{x \in \mathbb{T}_n^d} \zeta(x) \delta_x$ by a smooth interpolation.

Let $L_2(\mathbb{T}_n^d)$ be the Hilbert space of all functions on \mathbb{T}_n^d with inner product

$$
\langle \rho_1, \rho_2 \rangle_n := \frac{1}{n^d} \sum_{x \in \mathbb{T}_n^d} \rho_1(x) \rho_2(x)
$$

Replace $\tilde{\rho} = \frac{1}{n^d} \sum_{x \in \mathbb{T}_n^d} \rho(x) \delta_x$ and $\tilde{\zeta} = \frac{1}{n^d} \sum_{x \in \mathbb{T}_n^d} \zeta(x) \delta_x$ by a smooth interpolation.

Let $L_2(\mathbb{T}_n^d)$ be the Hilbert space of all functions on \mathbb{T}_n^d with inner product

$$
\langle \rho_1, \rho_2 \rangle_n := \frac{1}{n^d} \sum_{x \in \mathbb{T}_n^d} \rho_1(x) \rho_2(x)
$$

 $L_2(\mathbb{T}^d)$ be the usual L_2 -space of function on \mathbb{T}^d with

$$
\langle g_1,g_2\rangle:=\int_{\mathbb{T}^d}g_1(x)g_2(x)dx.
$$

Replace $\tilde{\rho} = \frac{1}{n^d} \sum_{x \in \mathbb{T}_n^d} \rho(x) \delta_x$ and $\tilde{\zeta} = \frac{1}{n^d} \sum_{x \in \mathbb{T}_n^d} \zeta(x) \delta_x$ by a smooth interpolation.

Let $L_2(\mathbb{T}_n^d)$ be the Hilbert space of all functions on \mathbb{T}_n^d with inner product

$$
\langle \rho_1, \rho_2 \rangle_n := \frac{1}{n^d} \sum_{x \in \mathbb{T}_n^d} \rho_1(x) \rho_2(x)
$$

 $L_2(\mathbb{T}^d)$ be the usual L_2 -space of function on \mathbb{T}^d with

$$
\langle g_1,g_2\rangle:=\int_{\mathbb{T}^d}g_1(x)g_2(x)dx.
$$

\n- $$
s_k(x) = e^{2\pi i k \cdot x}
$$
, $k \in \mathbb{Z}^d$, $x \in \mathbb{T}^d \supset \mathbb{T}_n^d$
\n- $-$ basis vectors on $L_2(\mathbb{T}_n^d)$ and $L_2(\mathbb{T}^d)$, and $-$ eigenvectors for discrete and continuous diff. operators
\n

Replace $\tilde{\rho} = \frac{1}{n^d} \sum_{x \in \mathbb{T}_n^d} \rho(x) \delta_x$ and $\tilde{\zeta} = \frac{1}{n^d} \sum_{x \in \mathbb{T}_n^d} \zeta(x) \delta_x$ by a smooth interpolation.

Let $L_2(\mathbb{T}_n^d)$ be the Hilbert space of all functions on \mathbb{T}_n^d with inner product

$$
\langle \rho_1, \rho_2 \rangle_n := \frac{1}{n^d} \sum_{x \in \mathbb{T}_n^d} \rho_1(x) \rho_2(x)
$$

 $L_2(\mathbb{T}^d)$ be the usual L_2 -space of function on \mathbb{T}^d with

$$
\langle g_1,g_2\rangle:=\int_{\mathbb{T}^d}g_1(x)g_2(x)dx.
$$

\n- $$
c_k(x) = e^{2\pi i k \cdot x}
$$
, $k \in \mathbb{Z}^d$, $x \in \mathbb{T}^d \supset \mathbb{T}_n^d$
\n- $-$ basis vectors on $L_2(\mathbb{T}_n^d)$ and $L_2(\mathbb{T}^d)$, and $-$ eigenvectors for discrete and continuous diff. operators
\n

$$
L_2(\mathbb{T}_n^d) \ni \rho = \sum_{k \in \mathbb{Z}_n^d} \langle \rho, \varsigma_k \rangle_{n \leq k} \text{ on } \mathbb{T}_n^d, \quad L_2(\mathbb{T}^d) \ni g = \sum_{k \in \mathbb{Z}^d} \langle g, \varsigma_k \rangle_{\varsigma_k} \text{ on } \mathbb{T}^d
$$

Vitalii Konarovskyi (University of Hamburg and Institution CLT for SSEP Becember 10, 2024 16/22

New (smooth) lifting of discrete space

For functions $\rho \in L_2(\mathbb{T}_n^d)$ and $\varphi \in L_2(\mathbb{T}^d)$ define

$$
\mathrm{ex}_n\rho:=\sum_{k\in\mathbb{Z}_n^d}\langle\rho,s_k\rangle_n s_k\quad\text{on }\mathbb{T}^d,\quad\mathrm{pr}_n\varphi:=\sum_{k\in\mathbb{Z}_n^d}\langle\varphi,s_k\rangle s_k\quad\text{on }\mathbb{T}^d
$$

∍

 QQ

New (smooth) lifting of discrete space

For functions $\rho \in L_2(\mathbb{T}_n^d)$ and $\varphi \in L_2(\mathbb{T}^d)$ define

$$
\mathrm{ex}_n\rho:=\sum_{k\in\mathbb{Z}_n^d}\langle\rho,\varsigma_k\rangle_n\varsigma_k\quad\text{on }\mathbb{T}^d,\quad\mathrm{pr}_n\varphi:=\sum_{k\in\mathbb{Z}_n^d}\langle\varphi,\varsigma_k\rangle\varsigma_k\quad\text{on }\mathbb{T}^d
$$

Basic properties of $ex_n f$ **and** $pr_n g$

•
$$
\operatorname{ex}_n \rho = \rho
$$
 on \mathbb{T}_n^d and $\operatorname{ex}_n \rho \in C^\infty(\mathbb{T}^d)$

\n- or
$$
\Pr_n \varphi
$$
 is well defined on \mathbb{T}_n^d for each $\varphi \in H_J$ for $H_J := \left\{ \varphi : \|\varphi\|_{H_J}^2 := \sum_{k \in \mathbb{Z}^d} (1 + |k|^2)^J |\langle \varphi, \varsigma_k \rangle|^2 \right\}, J \in \mathbb{R}$.
\n

$$
\bullet \ \langle \rho_1, \rho_2 \rangle_n = \langle \exp_1, \exp_2 \rangle \text{ and } \langle \rho, \mathrm{pr}_n \mathbf{g} \rangle_n = \langle \exp, \mathbf{g} \rangle
$$

$$
\bullet \ \|\text{pr}_n \mathbf{g} - \mathbf{g}\|_{H_J} \leq \frac{1}{n} \|\mathbf{g}\|_{H_{J+1}}, \ \|\text{ex}_n \varphi - \varphi\|_{H_J} \leq \frac{C}{n} \|\varphi\|_{C^{J+2+\frac{d}{2}}} \dots
$$

New (smooth) lifting of discrete space

For functions $\rho \in L_2(\mathbb{T}_n^d)$ and $\varphi \in L_2(\mathbb{T}^d)$ define

$$
\mathrm{ex}_n\rho:=\sum_{k\in\mathbb{Z}_n^d}\langle\rho,\varsigma_k\rangle_n\varsigma_k\quad\text{on }\mathbb{T}^d,\quad\mathrm{pr}_n\varphi:=\sum_{k\in\mathbb{Z}_n^d}\langle\varphi,\varsigma_k\rangle\varsigma_k\quad\text{on }\mathbb{T}^d
$$

Basic properties of $ex_n f$ **and** $pr_n g$

•
$$
\operatorname{ex}_n \rho = \rho
$$
 on \mathbb{T}_n^d and $\operatorname{ex}_n \rho \in C^\infty(\mathbb{T}^d)$

\n- or
$$
\text{pr}_n\varphi
$$
 is well defined on \mathbb{T}_n^d for each $\varphi \in H_J$ for $H_J := \left\{ \varphi : \|\varphi\|_{H_J}^2 := \sum_{k \in \mathbb{Z}^d} (1 + |k|^2)^J |\langle \varphi, \varsigma_k \rangle|^2 \right\}, J \in \mathbb{R}.$
\n

$$
\bullet \ \langle \rho_1, \rho_2 \rangle_n = \langle \exp_1, \exp_2 \rangle \text{ and } \langle \rho, \mathrm{pr}_n \mathbf{g} \rangle_n = \langle \exp, \mathbf{g} \rangle
$$

$$
\bullet \ \|\text{pr}_n g - g\|_{H_j} \leq \frac{1}{n} \|g\|_{H_{j+1}}, \ \|\text{ex}_n \varphi - \varphi\|_{H_j} \leq \frac{C}{n} \|\varphi\|_{C^{j+2+\frac{d}{2}}} \dots
$$

$$
\langle \varphi, \tilde{\rho} \rangle = \frac{1}{(2n+1)^d} \sum_{x \in \mathbb{T}_n^d} \varphi(x) \rho(x) = \langle \varphi, \rho \rangle_n
$$

$$
= \langle \text{pr}_n \varphi, \rho \rangle_n + O(1/n) = \langle \varphi, \text{ex}_n \rho \rangle + O(1/n)
$$

Vitalii Konarovskyi (University of Hamburg and Institutive CLT for SSEP Becember 10, 2024 17 / 22

Now for $F(\exp_i, \exp_i) := f(\langle \psi, \exp_i \rangle, \langle \varphi, \exp_i \rangle) = f(\langle \mathrm{pr}_i, \psi, \rho \rangle_n, \langle \mathrm{pr}_i, \varphi, \zeta \rangle_n)$

Now for $F(\exp_i, \exp_j) := f(\langle \psi, \exp_i \rangle, \langle \varphi, \exp_j \rangle) = f(\langle \mathrm{pr}_i \psi, \rho \rangle, \langle \mathrm{pr}_i \varphi, \zeta \rangle_n)$, we get

$$
\mathcal{G}_{n}^{FF}F(\exp_{n}\rho,\exp_{n}\zeta) = \frac{1}{2}\partial_{1}f\langle\Delta_{n}\mathrm{pr}_{n}\psi,\rho\rangle_{n} + \frac{1}{2}\partial_{2}f\langle\Delta_{n}\mathrm{pr}_{n}\varphi,\zeta\rangle_{n} \n+ \frac{1}{2}\partial_{2}^{2}f\sum_{j=1}^{d}\langle(\partial_{n,j}\mathrm{pr}_{n}\varphi)^{2},\tau_{j}\eta+\eta-2\eta\tau_{j}\eta\rangle_{n} + O\left(\frac{1}{n^{d/2+1}}\right)
$$

Now for $F(\exp_i, \exp_j) := f(\langle \psi, \exp_i \rangle, \langle \varphi, \exp_j \rangle) = f(\langle \mathrm{pr}_i, \psi, \rho \rangle, \langle \mathrm{pr}_i, \varphi, \zeta \rangle_n)$, we get $\mathcal{G}_n^{FF}F(\exp\cos\alpha\zeta) = \frac{1}{2}\partial_1 f \langle \exp\Delta_n \mathrm{pr}_n \psi, \exp \rangle + \frac{1}{2}$ $\frac{2}{2}\partial_2 f \langle \exp \Delta_n \mathrm{pr}_n \varphi, \exp \langle \rangle$ $+\frac{1}{2}$ $\frac{1}{2}\partial_2^2 f\sum_{i=1}^d\left\langle \exp\left(\partial_{n,j}\text{pr}_n\varphi\right)^2,\exp\left(\tau_j\eta+\eta-2\eta\tau_j\eta\right)\right\rangle +O\left(\frac{1}{n^{d/2}}\right)$ j=1 $rac{1}{n^{d/2+1}}$

つひひ

Now for $F(\exp_i, \exp_i) := f(\langle \psi, \exp_i \rangle, \langle \varphi, \exp_i \rangle) = f(\langle \mathrm{pr}_i, \psi, \rho \rangle_n, \langle \mathrm{pr}_i, \varphi, \zeta \rangle_n)$, we get $\mathcal{G}_n^{FF}F(\exp\cos\alpha\zeta) = \frac{1}{2}\partial_1 f \langle \exp\Delta_n \mathrm{pr}_n \psi, \exp \rangle + \frac{1}{2}$ $\frac{2}{2}\partial_2 f \langle \exp \Delta_n \mathrm{pr}_n \varphi, \exp \langle \rangle$ $+\frac{1}{2}$ $\frac{1}{2}\partial_2^2 f\sum_{i=1}^d\left\langle \exp\left(\partial_{n,j}\text{pr}_n\varphi\right)^2,\exp\left(\tau_j\eta+\eta-2\eta\tau_j\eta\right)\right\rangle +O\left(\frac{1}{n^{d/2}}\right)$ j=1 $rac{1}{n^{d/2+1}}$ $\left|\mathcal{G}_{n}^{FF}\mathcal{F}-\mathcal{G}_{n}^{OU}\mathcal{F}\right|\lesssim\frac{1}{n^{d_{\prime}}}$ $\frac{1}{n^{d/2}}\left|\left\langle \exp\left(\pi_{n}\zeta,\exp\left(\tau_{j}\rho\left|\partial_{n,j}\varphi\right|^{2}\right)\right\rangle \right|+\left\langle \exp\left(\zeta\tau_{j}\zeta\right),\exp\partial_{2}^{2}f\left(\partial_{n,j}\mathrm{pr}_{n}\varphi\right)^{2}\right\rangle ^{2}+\dots\right|$

We now need only to control:

Now for $F(\exp_i, \exp_i) := f(\langle \psi, \exp_i \rangle, \langle \varphi, \exp_i \rangle) = f(\langle \mathrm{pr}_i, \psi, \rho \rangle_n, \langle \mathrm{pr}_i, \varphi, \zeta \rangle_n)$, we get $\mathcal{G}_n^{FF}F(\exp\cos\alpha\zeta) = \frac{1}{2}\partial_1 f \langle \exp\Delta_n \mathrm{pr}_n \psi, \exp \rangle + \frac{1}{2}$ $\frac{2}{2}\partial_2 f \langle \exp \Delta_n \mathrm{pr}_n \varphi, \exp \langle \rangle$ $+\frac{1}{2}$ $\frac{1}{2}\partial_2^2 f\sum_{i=1}^d\left\langle \exp\left(\partial_{n,j}\text{pr}_n\varphi\right)^2,\exp\left(\tau_j\eta+\eta-2\eta\tau_j\eta\right)\right\rangle +O\left(\frac{1}{n^{d/2}}\right)$ j=1 $rac{1}{n^{d/2+1}}$ $\left|\mathcal{G}_{n}^{FF}\mathcal{F}-\mathcal{G}_{n}^{OU}\mathcal{F}\right|\lesssim\frac{1}{n^{d_{\prime}}}$ $\frac{1}{n^{d/2}}\left|\left\langle \exp\left(\pi_{n}\zeta,\exp\left(\tau_{j}\rho\left|\partial_{n,j}\varphi\right|^{2}\right)\right\rangle \right|+\left\langle \exp\left(\zeta\tau_{j}\zeta\right),\exp\partial_{2}^{2}f\left(\partial_{n,j}\mathrm{pr}_{n}\varphi\right)^{2}\right\rangle ^{2}+\dots\right|$

We now need only to control:

• The expectations:

$$
\mathbb{E}\left|\left\langle \exp\zeta^n_t,\exp\left(\tau_j\rho^n_t\left|\partial_{n,j}\varphi\right|^2\right)\right\rangle\right|\lesssim \|\exp\zeta^n_t\|_{C^J}\mathbb{E}\|\exp\zeta^n_t\|_{H_{-I}}
$$

Now for $F(\exp_{n} \rho, \exp_{n} \zeta) := f(\langle \psi, \exp_{n} \rangle, \langle \varphi, \exp_{n} \zeta \rangle) = f(\langle \mathrm{pr}_{n} \psi, \rho \rangle_{n}, \langle \mathrm{pr}_{n} \varphi, \zeta \rangle_{n})$, we get $\mathcal{G}_n^{FF}F(\exp\cos\alpha\zeta) = \frac{1}{2}\partial_1 f \langle \exp\Delta_n \mathrm{pr}_n \psi, \exp \rangle + \frac{1}{2}$ $\frac{2}{2}\partial_2 f \langle \exp \Delta_n \mathrm{pr}_n \varphi, \exp \langle \rangle$ $+\frac{1}{2}$ $\frac{1}{2}\partial_2^2 f\sum_{i=1}^d\left\langle \exp\left(\partial_{n,j}\text{pr}_n\varphi\right)^2,\exp\left(\tau_j\eta+\eta-2\eta\tau_j\eta\right)\right\rangle +O\left(\frac{1}{n^{d/2}}\right)$ j=1 $rac{1}{n^{d/2+1}}$ $\left|{\cal G}^{FF}_n F - {\cal G}^{OU}_n F\right| \lesssim \frac{1}{n^{d_n}}$ $\frac{1}{n^{d/2}}\left|\left\langle \exp\left(\pi_{n}\zeta,\exp\left(\tau_{j}\rho\left|\partial_{n,j}\varphi\right|^{2}\right)\right\rangle \right|+\left\langle \exp\left(\zeta\tau_{j}\zeta\right),\exp\partial_{2}^{2}f\left(\partial_{n,j}\mathrm{pr}_{n}\varphi\right)^{2}\right\rangle ^{2}+\dots\right|$

We now need only to control:

O The expectations:

$$
\mathbb{E}\left|\left\langle \exp_{n} \zeta_{t}^{n}, \exp\left(\tau_{j} \rho_{t}^{n} |\partial_{n,j} \varphi|^{2}\right)\right\rangle\right| \lesssim \|\exp_{n} \rho_{t}^{n}\|_{C^{J}} \mathbb{E}\|\exp_{n} \zeta_{t}^{n}\|_{H_{-I}}
$$

Using the Fourier analysis, the term $\mathbb{E} \left< \text{ex}_n \left(\zeta_t^n \tau_j \zeta_t^n \right), \text{ex}_n \partial_2^2 f (\ldots)^2 \right>$ can be controlled via

$$
\mathbb{E}\prod_{i=1}^4\left(\eta_t^n(x_i)-\rho_t^n(x_i)\right)\lesssim\frac{1}{n}
$$

Now for $F(\exp_i, \exp_i) := f(\langle \psi, \exp_i \rangle, \langle \varphi, \exp_i \rangle) = f(\langle \mathrm{pr}_i, \psi, \rho \rangle_n, \langle \mathrm{pr}_i, \varphi, \zeta \rangle_n)$, we get $\mathcal{G}_n^{FF}F(\exp\cos\alpha\zeta) = \frac{1}{2}\partial_1 f \langle \exp\Delta_n \mathrm{pr}_n \psi, \exp \rangle + \frac{1}{2}$ $\frac{2}{2}\partial_2 f \langle \exp \Delta_n \mathrm{pr}_n \varphi, \exp \langle \rangle$ $+\frac{1}{2}$ $\frac{1}{2}\partial_2^2 f \sum_{i=1}^d$ j=1 $\langle \exp \left(\partial_{n,j} \text{pr}_n \varphi \right)^2, \exp \left(\tau_j \eta + \eta - 2 \eta \tau_j \eta \right) \rangle + O\left(\frac{1}{n^{d/2}} \right)$ $rac{1}{n^{d/2+1}}$ $\left|{\cal G}^{FF}_n F - {\cal G}^{OU}_n F\right| \lesssim \frac{1}{n^{d_n}}$ $\frac{1}{n^{d/2}}\left|\left\langle \exp\left(\pi_{n}\zeta,\exp\left(\tau_{j}\rho\left|\partial_{n,j}\varphi\right|^{2}\right)\right\rangle \right|+\left\langle \exp\left(\zeta\tau_{j}\zeta\right),\exp\partial_{2}^{2}f\left(\partial_{n,j}\mathrm{pr}_{n}\varphi\right)^{2}\right\rangle ^{2}+\dots\right|$

We now need only to control:

• The expectations:

$$
\mathbb{E}\left|\left\langle \exp_{n}\zeta_{t}^{n}, \exp_{n}\left(\tau_{j}\rho_{t}^{n}|\partial_{n,j}\varphi|^{2}\right)\right\rangle\right|\lesssim\left\|\exp_{n}\rho_{t}^{n}\right\|_{C^{J}}\mathbb{E}\left\|\exp_{n}\zeta_{t}^{n}\right\|_{H_{-I}}
$$

Using the Fourier analysis, the term $\mathbb{E} \left< \text{ex}_n \left(\zeta_t^n \tau_j \zeta_t^n \right), \text{ex}_n \partial_2^2 f (\ldots)^2 \right>$ can be controlled via

$$
\mathbb{E}\prod_{i=1}^4\left(\eta_t^n(x_i)-\rho_t^n(x_i)\right)\lesssim\frac{1}{n}
$$

 \bullet We can compare generators on Frechet diff. functio[ns](#page-59-0) [on](#page-61-0) $H_J \times H_{-J}$ $H_J \times H_{-J}$

Differentiability of $P_t^{OU}F(\text{ex}_n\rho,\text{ex}_n\zeta)$

A solution to

$$
d\rho_t^{\infty} = \frac{1}{2} \Delta \rho_t^{\infty} dt
$$

$$
d\zeta_t^{\infty} = \frac{1}{2} \Delta \zeta_t^{\infty} dt + \nabla \cdot \left(\sqrt{\rho_t^{\infty} (1 - \rho_t^{\infty})} dW_t \right)
$$

exists for all $\rho_0^{\infty} \in L_2(\mathbb{T}^d; [0,1])$ and $\zeta_0^{\infty} \in H_{-1}$ for $l > \frac{d}{2} + 1$.

 QQ

Differentiability of $P_t^{OU}F(\text{ex}_n\rho,\text{ex}_n\zeta)$

A solution to

$$
d\rho_t^{\infty} = \frac{1}{2} \Delta \rho_t^{\infty} dt
$$

$$
d\zeta_t^{\infty} = \frac{1}{2} \Delta \zeta_t^{\infty} dt + \nabla \cdot \left(\sqrt{\rho_t^{\infty} (1 - \rho_t^{\infty})} dW_t \right)
$$

exists for all $\rho_0^{\infty} \in L_2(\mathbb{T}^d; [0,1])$ and $\zeta_0^{\infty} \in H_{-1}$ for $l > \frac{d}{2} + 1$.

 $\mathsf{For} \ \mathsf{F} \in \mathrm{C}(\mathsf{H}_{J} \times \mathsf{H}_{-I}) \ (\text{e.g.} \ \mathsf{F} = f(\langle \psi, \cdot \rangle, \langle \varphi, \cdot \rangle)) \ \text{define} \ \mathsf{U}_{t}(\rho^{\infty}_{0}, \zeta^{\infty}_{0}) := \mathbb{E} \mathsf{F}\left(\rho^{\infty}_{t}, \zeta^{\infty}_{t}\right)$

 QQ

э

Differentiability of $P_t^{OU}F(\text{ex}_n\rho,\text{ex}_n\zeta)$

A solution to

$$
d\rho_t^{\infty} = \frac{1}{2} \Delta \rho_t^{\infty} dt
$$

$$
d\zeta_t^{\infty} = \frac{1}{2} \Delta \zeta_t^{\infty} dt + \nabla \cdot \left(\sqrt{\rho_t^{\infty} (1 - \rho_t^{\infty})} dW_t \right)
$$

exists for all $\rho_0^{\infty} \in L_2(\mathbb{T}^d; [0,1])$ and $\zeta_0^{\infty} \in H_{-1}$ for $l > \frac{d}{2} + 1$.

 $\mathsf{For} \ \mathsf{F} \in \mathrm{C}(\mathsf{H}_{J} \times \mathsf{H}_{-I}) \ (\text{e.g.} \ \mathsf{F} = f(\langle \psi, \cdot \rangle, \langle \varphi, \cdot \rangle)) \ \text{define} \ \mathsf{U}_{t}(\rho^{\infty}_{0}, \zeta^{\infty}_{0}) := \mathbb{E} \mathsf{F}\left(\rho^{\infty}_{t}, \zeta^{\infty}_{t}\right)$

Proposition 1 [Gess, K. '24]

Let $I > \frac{d}{2}+1$ and $F \in \mathrm{C}^{2,4}_b(H_{-I})$. Then $U_t(\rho_0^{\infty}, \zeta_0^{\infty}) = \mathbb{E} F\left(\zeta_t^{\infty}\right) \in \mathrm{C}^{1,3}_b(H_J \times H_{-I})$ for $J > \frac{d}{2}$. Moreover,

$$
D_1U_t(\rho_0^{\infty},\zeta_0^{\infty})[h] = \frac{1}{2}\mathbb{E}\left[D^2F(\zeta_t^{\infty}) : DV_t(\rho_0^{\infty})[h]\right]
$$

with

$$
V_t(\rho_0^{\infty})(\varphi, \psi) = \text{Cov}(\langle \varphi, \zeta_t^{\infty} \rangle, \langle \psi, \zeta_t^{\infty} \rangle)
$$

=
$$
\frac{1}{2} \int_0^t \left\langle \nabla P_{t-s}^{HE} \varphi \cdot \nabla P_{t-s}^{HE} \psi, \rho_s^{\infty} (1 - \rho_s^{\infty}) \right\rangle ds
$$

Vitalii Konarovskyi (University of Hamburg and Institutive CLT for SSEP December 10, 2024 19/22

Comparison of dynamics

Recall that (ρ^n_t, ζ^n_t) is the mean process together with the fluctuation field of SSEP.

目

4 0 F

Comparison of dynamics

Recall that (ρ^n_t, ζ^n_t) is the mean process together with the fluctuation field of SSEP. $(\rho_t^{\infty,n}, \zeta_t^{\infty,n})$ is a solution to

$$
d\rho_t^{\infty} = \frac{1}{2} \Delta \rho_t^{\infty} dt
$$

$$
d\zeta_t^{\infty} = \frac{1}{2} \Delta \zeta_t^{\infty} dt + \nabla \cdot \left(\sqrt{\rho_t^{\infty} (1 - \rho_t^{\infty})} dW_t \right)
$$

started from $(\mathrm{ex}_n\rho_0^n,\mathrm{ex}_n\zeta_0^n)$.

4 D F

 298

э

∋ x a ∋ x

Comparison of dynamics

Recall that (ρ^n_t, ζ^n_t) is the mean process together with the fluctuation field of SSEP. $(\rho_t^{\infty,n}, \zeta_t^{\infty,n})$ is a solution to

$$
d\rho_t^{\infty} = \frac{1}{2} \Delta \rho_t^{\infty} dt
$$

$$
d\zeta_t^{\infty} = \frac{1}{2} \Delta \zeta_t^{\infty} dt + \nabla \cdot \left(\sqrt{\rho_t^{\infty} (1 - \rho_t^{\infty})} dW_t \right)
$$

started from $(\mathrm{ex}_n\rho_0^n,\mathrm{ex}_n\zeta_0^n)$.

Then for each I, J large enough and $F \in C_b(H_J \times H_{-I})$

$$
|\mathbb{E} F(\exp_n\rho_t^n, \exp_n\zeta_t^n) - \mathbb{E} F(\rho_t^{\infty,n}, \zeta_t^{\infty,n})| \leq \int_0^t \left| \mathbb{E} \left[\left(\mathcal{G}^{FF} - \mathcal{G}^{OU} \right) P_{t-s}^{OU} F(\exp_n\rho_t^n, \exp_n\zeta_t^n) \right] \right| ds
$$

\$\lesssim ||F||_{C^{1,3}} \int_0^t \left(\frac{1}{n} + \frac{1}{n^{d/2}} ||\exp_n^{\theta}||_{C^J} \mathbb{E} ||\exp_n\zeta_s^n||_{H_{-I}} + \mathbb{E} \left\langle \exp_n\left(\zeta_s^n \tau_j \zeta_s^n \right), \ldots \right\rangle^2 + \ldots \right) ds

Vitalii Konarovskyi (University of Hamburg and Institutive CLT for SSEP Becember 10, 2024 20/22

4 D F

э

Berry-Esseen bound for the initial fluctuations

It remains only to compare

$$
\mathbb{E} F(\rho_t^{\infty,n},\zeta_t^{\infty,n}) - \mathbb{E} F(\rho_t^{\infty},\zeta_t^{\infty}) = P_t^{OU} F(\exp_n\rho_0^n,\exp_n\zeta_0^n) - P_t^{OU} F(\rho_0,\zeta_0)
$$

where $\rho_{\bm{t}}^\infty$ started from the initial profile ρ_0 and $\zeta_{\bm{t}}$ started from the centered Gaussian distribution with

$$
\mathbb{E}\langle \zeta_0,\varphi\rangle^2=\langle \rho_0(1-\rho_0)\varphi,\varphi\rangle.
$$

It is enough to compare only

$$
\mathbb{E} F(\mathrm{ex}_n\zeta_0^n)-\mathbb{E} F(\mathrm{pr}_n\zeta_0),
$$

for $\mathit{F}\in C^{3}(\mathit{H}_{-l})$, where

$$
\mathrm{ex}_n\zeta_0^n := \sum_{k \in \mathbb{Z}_n^d} \langle \zeta_0^n, \varsigma_k \rangle_n \varsigma_k, \quad \mathrm{pr}_n\zeta_0 := \sum_{k \in \mathbb{Z}_n^d} \langle \zeta_0, \varsigma_k \rangle \varsigma_k
$$

Is enough to compare for $f\in\mathrm{C}^3\left(\mathbb{R}^{\mathbb{Z}_n^d}\right)$

$$
\mathbb{E} f\left(\left((1+|k|^2)^{-1/2}\langle \zeta_0^n,\varsigma_k\rangle_n\right)_{k\in\mathbb{Z}_n^d}\right)-\mathbb{E} f\left(\left((1+|k|^2)^{-1/2}\langle \zeta_0,\varsigma_k\rangle\right)_{k\in\mathbb{Z}_n^d}\right).
$$

Apply multidimensional Berry-Essen theorem [Meck[es](#page-66-0) '[09](#page-68-0)[\]](#page-66-0)

References

- [1] Benjamin Gess and Vitalii Konarovskyi. A quantitative central limit theorem for the simple symmetric exclusion process (2024), arXiv:2408.01238
- [2] Claude Kipnis and Claudio Landim. Scaling limits of interacting particle systems (1999)
- [3] K. Ravishankar. Fluctuations from the hydrodynamical limit for the symmetric simple exclusion in \mathbb{Z}^d (1992)
- [4] P. A. Ferrari, E. Presutti, E. Scacciatelli, and M. E. Vares. The symmetric simple exclusion process. I. Probability estimates (1991)
- [5] Elizabeth Meckes. On Stein's method for multivariate normal approximation (2009)

Thank you!

∢ ロ ▶ - ∢ 母 ▶ - ∢ ヨ ▶ -∢ ヨ ▶

 QQ