

Reminder: no lecture Fri 11th Dec
extra lecture Fri 18th Dec 09:15 if desired

Thm $X = \prod_{s \in S} X_s$, $Y = \prod_{s \in S} Y_s$ (Banach spaces, with sup-norm $L: X \rightarrow Y$ linear
 φ -exp. local, bdd inverse, $y \in Y$ (ζ, λ)-exp.
located around $0 \in S$ (i.e. $|y_s| \leq C \lambda^{d(s,0)}$)
 $\Rightarrow x = \tilde{L}y$ is (WC, μ) -exp. located around 0
for some W, μ functions of $\|L^{-1}\|_{\lambda, \varphi}$

Proof: Let $z \in [1, \zeta]$, $z\lambda < 1$,
 $\tilde{y}_r = y_r z^{d(r,0)}$, $\tilde{x}_r = x_r z^{d(r,0)}$
Then $\|\tilde{y}\| < \infty$. Want to bound $\|\tilde{x}\|$
 $Lx = y \Rightarrow \sum_s L_{rs} z^{d(r,0) - d(s,0)} \tilde{x}_s = \tilde{y}_r$
 $\hookrightarrow (L + L^\circ) \tilde{x} = \tilde{y}$ with $L^\circ_{rs} = L_{rs} (z^{d(r,0) - d(s,0)} - 1)$
Let $\tilde{L}_{rs} = L_{rs} z^{d(r,s)}$ $\|\tilde{L}\| \leq \varphi(z) < \infty$
 $d(r,s) - d(s,0) \leq d(r,s)$ so
 $|L^\circ_{rs}| \leq |\tilde{L}_{rs} - L_{rs}| \approx \|L^\circ\| \leq \|\tilde{L} - L\|$

For $z=1$ $\tilde{L}=L$. $\tilde{L}(z)$ depends continuously
on z , because given $1 \leq z_1 < z_2 < \zeta$
 $(\tilde{L}(z_2) - \tilde{L}(z_1))_{rs} = L_{rs} (z_2^{d(r,s)} - z_1^{d(r,s)})$
 $\hookrightarrow |(\tilde{L}(z_2) - \tilde{L}(z_1))_{rs}| \leq |L_{rs}| d(r,s) z_2^{d(r,s)} \log \frac{z_2}{z_1}$
 $\leq \frac{1}{e} \underbrace{\frac{\log z_2/z_1}{\log z_3/z_2}}_{\text{for any } z_3 > z_2} z_3^{d(r,s)} \|L_{rs}\|$

$\hookrightarrow \|\tilde{L}(z_2) - \tilde{L}(z_1)\| \leq \frac{1}{e} \frac{\log z_2/z_1}{\log z_3/z_2} \|\tilde{L}(z_3)\|$
 $\leq \frac{1}{e} \frac{\log z_2/z_1}{\log z_3/z_2} \varphi(z_3).$

In particular $\|\tilde{L}(z) - L\| \leq \beta(z) = \min_{z_3} \frac{1}{e} \frac{\log z_2/z_1}{\log z_3/z_2} \varphi(z_3)$
 $\rightarrow 0$ as $z \rightarrow 1$, so $< \|\tilde{L}\|^{-1}$ for z near enough to 1. Hence

$$\|(L + L^\circ)^{-1}\| \leq (\|\tilde{L}\|^{-1} - \|\tilde{L} - L\|)^{-1} \text{ for } z \text{ near 1.}$$

$$\hookrightarrow \|\tilde{x}\| \leq \frac{\|\tilde{y}\|}{\|\tilde{L}\|^{-1} - \|\tilde{L} - L\|}. \quad \square$$

Apply this to u.hyp splitting for CML thing

$S = S \times \mathbb{Z}$ by coding $\mathcal{D}F$ exp. loc.

$$\text{wrt } d((s_1, t_1), (s_2, t_2)) = d(s_1, s_2) + |t_2 - t_1|$$

and repeat construction of splitting & see that

scattering projections $P^{\pm}_x : TM_x \rightarrow$ are exp. loc wrt d_S

As for T only, get structure of u.hyp sets in $S \times T$

wrt exp. loc perturbation in space. Also S-T shadowing.

Similarly get Markov partitions for u.hyp sets of CML

(f. Pein & Sinai). A coding $\underline{x} = \underline{x}(\underline{\sigma})$

of the sets $\underline{x} = (x_s^t)_{s \in S}^{t \in \mathbb{Z}}$ of the u.hyp set by

S-T symbol tables $\underline{\sigma}$ from some allowed set Σ st.

every allowed table occurs & $\exists ! \underline{\sigma}$ for each \underline{x} .

Simple case $\sum = \left(\prod_{s \in S} \sum_s \right)^{\mathbb{Z}}$ with each \sum_s

finite e.g. uncoupled lattice of sigmoid maps $\sum_s = \{0, 1\}$

and $x_s^t = \chi_s(\dots \sigma_s^{t-1}, \sigma_s^t, \sigma_s^{t+1}, \dots)$

the unique σ whose bit into R_s, R_t of $\mathcal{D}F^t$

in sequence $(\sigma_s^t)_{t \in \mathbb{Z}}$

Now make a C^1 -small exp. loc coupling. By the

structure of u.hyp sets, can deform coding $\underline{x} = \underline{x}(\underline{\sigma})$

but now x_s^t depends in general on all $\sigma_s^{t'},$ (not just $\sigma_s^{t'})$

though exp. weakly on distant $(s', t').$

Same holds if uncoupled

dynamics has a general finite

graph of allowed transitions

e.g. at map $(\begin{array}{c} 1 \\ 2 \end{array} \rightarrow \begin{array}{c} 2 \\ 3 \end{array})$

(symbols = {5 digits})

$\begin{array}{c} + \\ \vdots \\ - \end{array}$	$\begin{array}{r} 1 \\ 0 \\ 1 \\ 0 \\ 1 \end{array}$
	$\begin{array}{r} 0 \\ 1 \\ 1 \\ 0 \\ 1 \end{array}$
	$\begin{array}{r} 1 \\ 0 \\ 0 \\ 1 \end{array}$

Also can use S-T shadowing then to extract M^{\pm}_{loc} for

locally more u.hyp sets without assuming nearby uncoupling.

Natural measures on u.hyp attractors for CML

Those which arise by starting in distant part
with a prob measure $\nu_M = \mathcal{D}M$ whose marginals on all
finite subsets of S are abscts and Hölder limits.

They are the Gibbs phases for "energy" with contribution

$$P_s^t(\underline{\sigma}) = \text{tr} \left[\log \mathcal{D}F_{-}(x_s^t(\underline{\sigma})) \right]$$

for spacetime site $(s, t).$

By w.h.p theory thus φ_s^t depends exp weakly
on $\sigma_{s'}^{t'}$ wrt $d(s', t'), (s, t)$

The connection with single dyn sys (SRB) is

$$\log |\det A| = \text{tr} \log A = \sum_s \text{tr} [\log A]_{ss}$$

+ cst from chain of bounds of log

see Bricmont & Kupiainen [91]

Thus natural measures for n.hyp attractors of CML

\Leftrightarrow eqm stat mech for a special class of spin systems,
on $S \times T$.

When $\dim(S \times T) \geq 2$ can expect to make
examples with non-unique Gibbs phase if sufficiently
coupled.

Challenge: Make a CML of iterated maps with
n.hyp attractor exhibiting non-unique S-T phase

f. Bardet & Keller got close.

Finished! So no lecture on 18 Dec after all!

Happy Holidays! Come to S-T Phases workshop
6-8 Jan!