

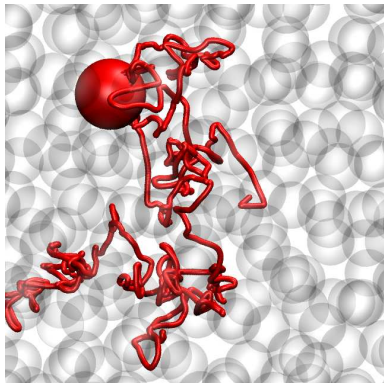
Microscopic flow around a diffusing tagged particle: effects of fluctuations on the hydrodynamics description

R. Vuilleumier

Département de chimie
Ecole normale supérieure – Paris

Workshop on Molecular Dynamics – Warwick, June 5, 2009

Diffusion in a Lennard-Jones fluid



10000 identical LJ particles
($T^* = 1.5$, $\rho^* = 0.5$)

- ▶ Friction from flow around the moving particle
- ▶ Relationship between diffusion and viscosity

$$D = \frac{k_B T}{c\pi\eta R}$$

- ▶ Flow boundary conditions?
- ▶ Direct estimation of the flux?

Stokes' flow

A sphere with a (small) velocity v_1 in an incompressible fluid of viscosity η

- ▶ velocity field
(slip boundary conditions – lab. frame):

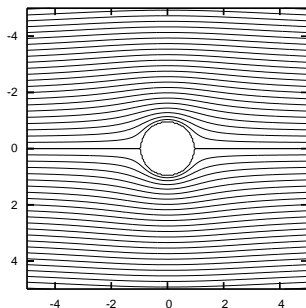
$$\vec{v}(\vec{r}) = \frac{f}{4\pi\eta} \left(\frac{1}{r} \cos\theta \vec{e}_r - \frac{1}{2r} \sin\theta \vec{e}_\theta \right)$$

- ▶ drag force f :

$$f = 4\pi\eta R v_1$$

- ▶ friction $\xi = 4\pi\eta R$

diffusion $D = \frac{k_B T}{\xi} = \frac{k_B T}{4\pi\eta R}$ (Stokes-Einstein relation)



Velocity field: linear response

Infinitesimal force

force f applied on particle 1 since $t = -\infty$ along x

Averages of observables

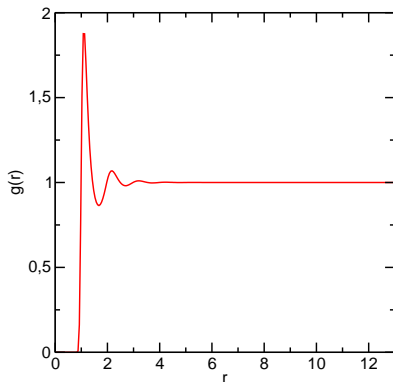
$$\begin{aligned}\langle A \rangle_f &\approx \langle A \rangle_0 + \beta f \int_{-\infty}^0 dt \langle v_{1,x}(t) A \rangle_0 \\ &\approx \langle A \rangle_0 + f \times \beta \langle \Delta r_{1,x} A \rangle_0\end{aligned}$$

$\Delta r_{1,x}$: displacement along x of particle 1 from $t = -\infty$ to $t = 0$

$A \equiv v_1$: $\beta \langle \Delta r_{1,x} v_1 \rangle_0$ is the mobility

Lenard-Jones fluid

Above the critical point: $T^* = 1.5$, $\rho^* = 0.5$ ($\sigma = 1$)



10000 identical LJ spheres $\rightarrow L = 27.14 \sigma$

Velocity field: diffusing particle frame

$$\rho \vec{v}(\vec{r}) = \beta \langle \Delta r_{1,x} (\vec{v}_i - \vec{v}_1) \delta(\vec{r} - \vec{r}_{1i}) \rangle$$

Use of symmetry ?

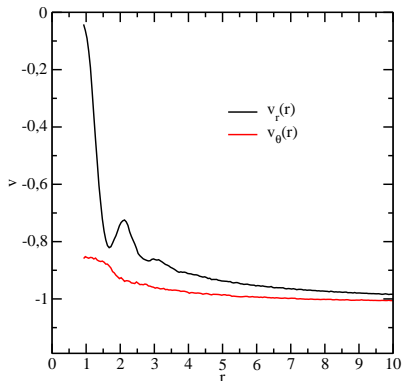
Velocity field: diffusing particle frame

$$\rho \vec{v}(\vec{r}) = \beta \langle \Delta r_{1,x} (\vec{v}_i - \vec{v}_1) \delta(\vec{r} - \vec{r}_{1i}) \rangle = f(r) \cos \theta \cdot \vec{e}_r + g(r) \sin \theta \cdot \vec{e}_\theta$$

Same symmetry as the Stokes flow around a sphere

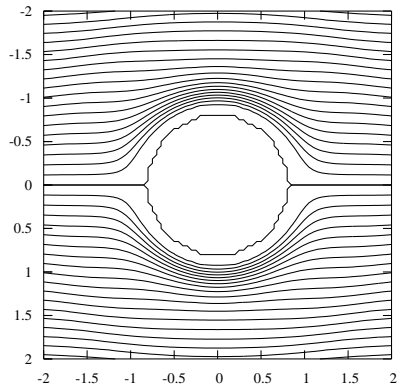
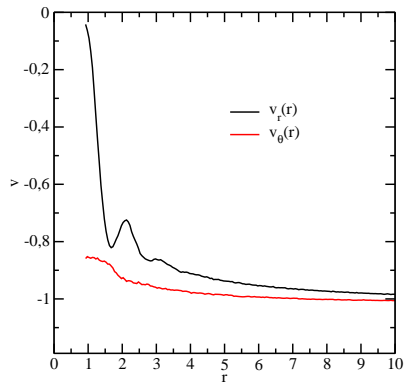
Velocity field: diffusing particle frame

$$\rho \vec{v}(\vec{r}) = \beta \langle \Delta r_{1,x} (\vec{v}_i - \vec{v}_1) \delta(\vec{r} - \vec{r}_{1i}) \rangle = f(r) \cos \theta \cdot \vec{e}_r + g(r) \sin \theta \cdot \vec{e}_\theta$$

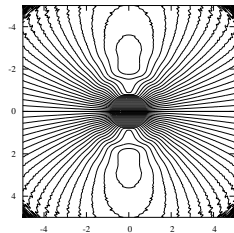


Velocity field: diffusing particle frame

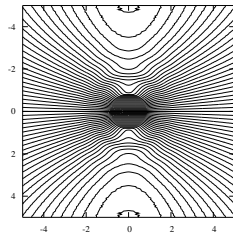
$$\rho \vec{v}(\vec{r}) = \beta \langle \Delta r_{1,x} (\vec{v}_i - \vec{v}_1) \delta(\vec{r} - \vec{r}_{1i}) \rangle = f(r) \cos \theta \cdot \vec{e}_r + g(r) \sin \theta \cdot \vec{e}_\theta$$



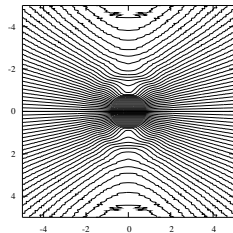
Laboratory frame: *backflow*



$t = 1$



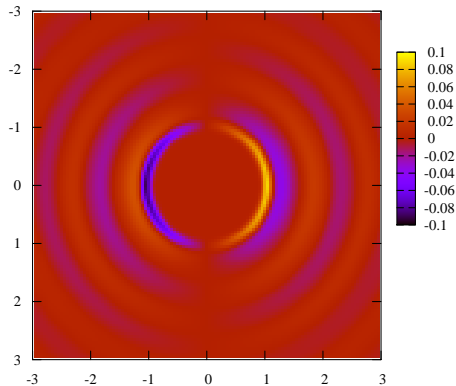
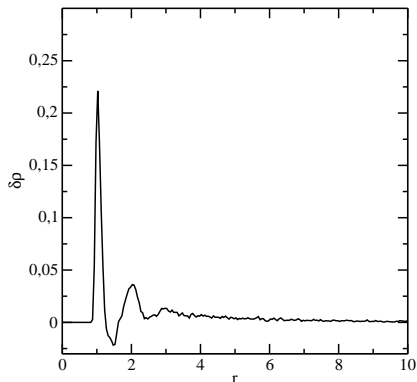
$t = 10$



$t = 100$

Density map

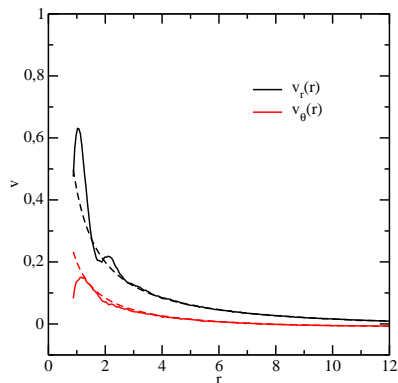
$$\delta\rho(\vec{r}) = \beta\langle\Delta r_{1,x}\delta(\vec{r}-\vec{r}_{1i})\rangle = h(r)\cos\theta$$



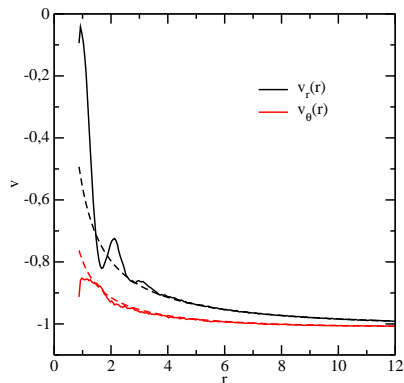
Force felt by the particle = applied force

Hydrodynamic flow – Boundary conditions

Fit of an hydrodynamic flow with **slip** boundary conditions



Labo. frame

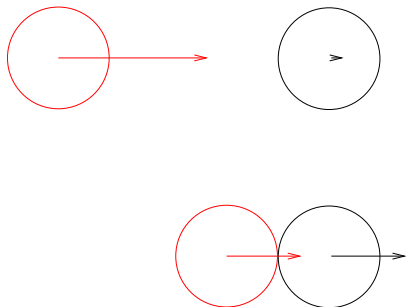


Part. frame

The normal velocity is non zero at contact!

Boundary conditions – normal velocity

Correlation between the particle velocity and the presence of a particle at contact

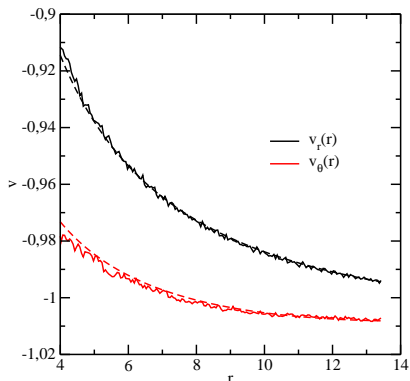


- ▶ velocity lower than the average tagged particle velocity
- ▶ **apparent normal velocity at contact**

Hynes, J. T., Kapral, R. et Weinberg, M., *JCP* **70**, 1456 (1979).

Dissipation through viscosity

Velocity field in the long range part



- ▶ Momentum flux due to viscosity corresponds to only a force $f = 0.76$ injected in the fluid
- ▶ Is there a dissipation term other than viscosity?

Equation of motion for the velocity field

At long times (laboratory frame):

$$\nabla_{\beta} \langle v_{i,\alpha} \cdot v_{i,\beta} \rangle_1(\vec{r}) - \sum_{j \neq 1} \langle f_{j \rightarrow i, \alpha} \rangle_1(\vec{r}) - \nabla_{\beta} \langle v_{i,\alpha} \cdot v_{1,\beta} \rangle_1(\vec{r}) = \langle f_{1 \rightarrow i, \alpha} \rangle_1(\vec{r})$$

Equation of motion for the velocity field

At long times (laboratory frame):

$$\nabla_{\beta} \langle v_{i,\alpha} \cdot v_{i,\beta} \rangle_1(\vec{r}) - \sum_{j \neq 1} \langle f_{j \rightarrow i, \alpha} \rangle_1(\vec{r}) - \nabla_{\beta} \langle v_{i,\alpha} \cdot v_{1,\beta} \rangle_1(\vec{r}) = \langle f_{1 \rightarrow i, \alpha} \rangle_1(\vec{r})$$

Equation of motion for the velocity field

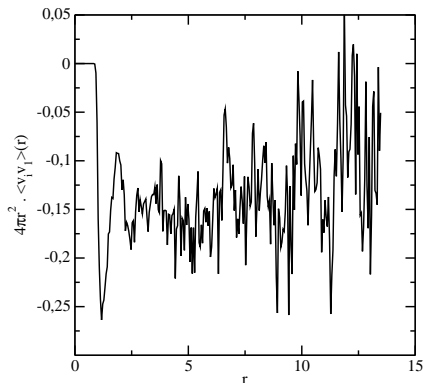
At long times (laboratory frame):

$$\nabla_{\beta} \langle v_{i,\alpha} \cdot v_{i,\beta} \rangle_1(\vec{r}) - \sum_{j \neq 1} \langle f_{j \rightarrow i, \alpha} \rangle_1(\vec{r}) - \nabla_{\beta} \langle v_{i,\alpha} \cdot v_{1,\beta} \rangle_1(\vec{r}) = \langle f_{1 \rightarrow i, \alpha} \rangle_1(\vec{r})$$

Equation of motion for the velocity field

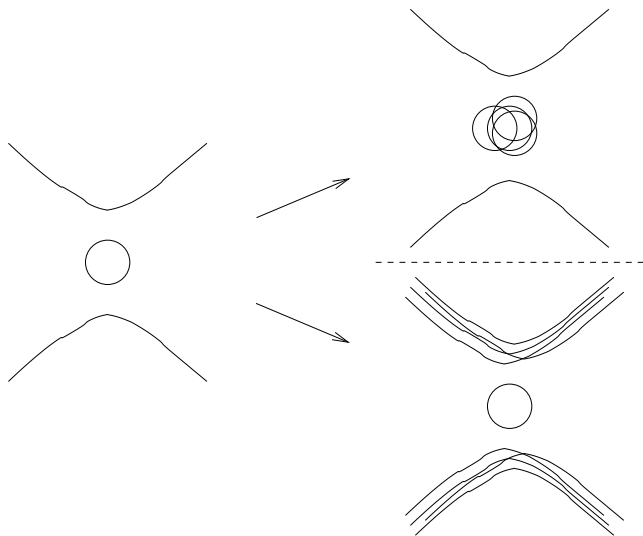
At long times (laboratory frame):

$$\nabla_{\beta} \langle v_{i,\alpha} \cdot v_{i,\beta} \rangle_1(\vec{r}) - \sum_{j \neq 1} \langle f_{j \rightarrow i, \alpha} \rangle_1(\vec{r}) - \nabla_{\beta} \langle v_{i,\alpha} \cdot v_{1,\beta} \rangle_1(\vec{r}) = \langle f_{1 \rightarrow i, \alpha} \rangle_1(\vec{r})$$



- ▶ Momentum flux through a sphere of radius r
- ▶ Contribution from the term $\langle v_{i,\alpha} \cdot v_{1,\beta} \rangle_1(\vec{r})$

Role of the central particle fluctuations



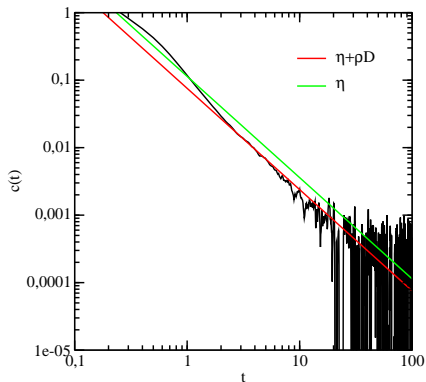
Asymptotic behavior of the velocity autocorrelation function

$c(t) \approx a_0 t^{-3/2}$ with

▶ $a_0 = \frac{2k_B T \rho^{1/2}}{4\pi\eta^{3/2}}$: viscosity only

▶ $a_0 = \frac{2k_B T \rho^{1/2}}{4\pi(\eta + \rho D)^{3/2}}$:

including fluctuations



Evolution equations

The evolution equations are modified

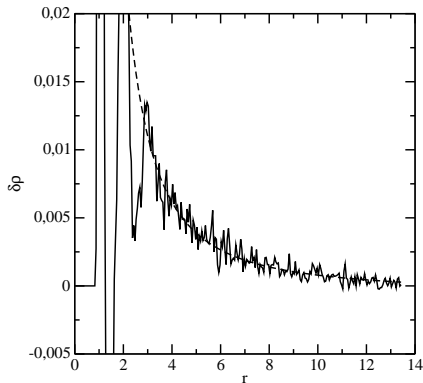
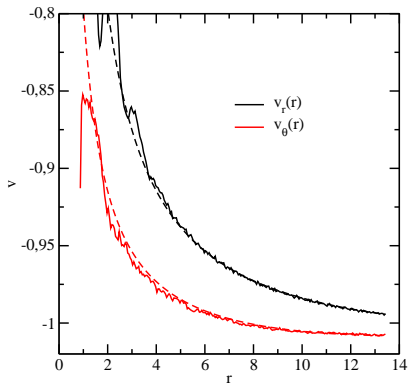
$$\rho \frac{d\vec{v}}{dt} = -\vec{\nabla} P + \eta \nabla^2 \vec{v} + \rho D \nabla^2 \vec{v}$$

$$\frac{d\delta\rho}{dt} = D \nabla^2 \delta\rho - \vec{\nabla} \rho \vec{v}(\vec{r})$$

In regions where $\rho \approx \text{cte.}$, the solution still satisfies $\vec{\nabla} \vec{v} = 0$

Renormalisation of viscosity

Hydrodynamique avec $f = 1$ et $\eta \rightarrow \eta + \rho D$



$$D = D_0 + \frac{k_B T}{4\pi(\eta + \rho D)} \frac{1}{R}$$

Bedeaux, D. et Mazur, P, *Physica* **73**, 431 (1974).

Diffusion mechanism

Most mobile configurations

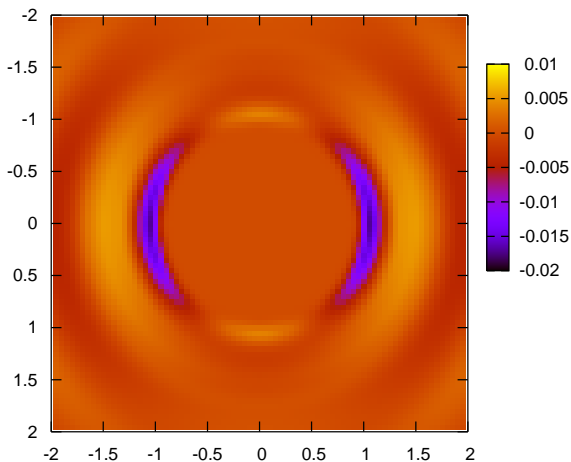
- ▶ Diffusion constant is the average of $\Delta r_{1,x} \cdot v_{1,x}$
- ▶ How variations of this quantity correlate with the local structure ?

Computation of

$$\langle \Delta r_{1,x} \cdot v_{1,x} \delta(\vec{r} - \vec{r}_{1i}) \rangle$$

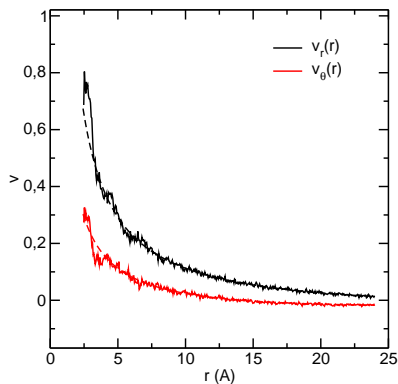
Most mobile configurations in a high density LJ fluid

$$T^* = 0.785, \rho^* = 0.8$$

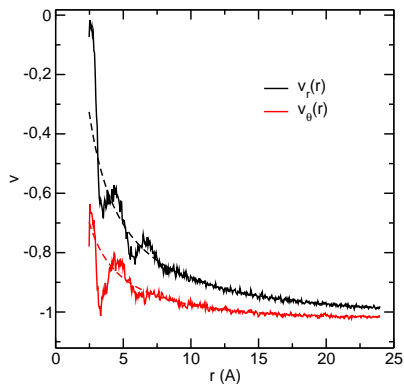


Elongated (de-structured) and squeezed solvation shell

Water self diffusion at ambient conditions



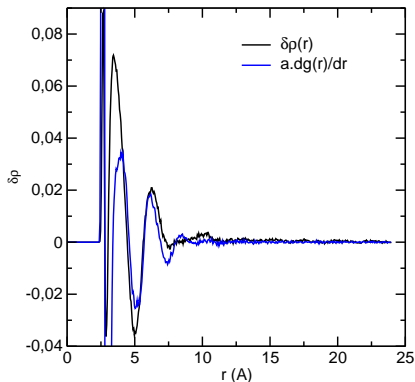
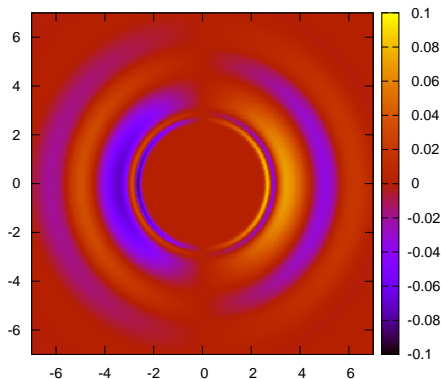
Labo. frame



Part. frame

Slip boundary conditions even for hydrogen bonded system!

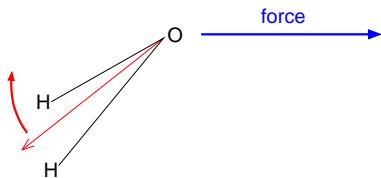
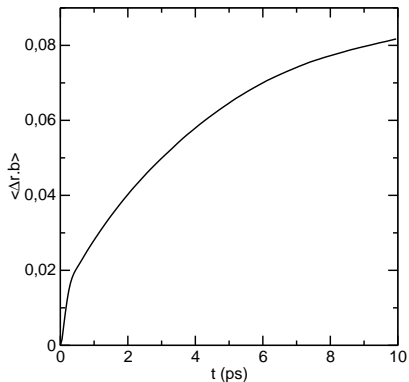
Displacement inside the cage



- ▶ The water molecule seems displaced in its cage.
- ▶ Except in its close vicinity
- ▶ The first shell is also destructured
- ▶ Characteristic retardation time $\tau_{ret} = 350$ fs

Rotation of the water molecule coupled to diffusion

Computation of $\langle \Delta \vec{r}_1 \cdot \vec{b}_1 \rangle$



In liquid water, meanfreepath = 0.04 Å

Conclusions

Summary

- ▶ Hydrodynamic is valid already for $r = 4$ with **slip** boundary conditions
- ▶ Effects of the central particle fluctuations:
 - ▶ **Boundary conditions with non-zero normal velocity**
 - ▶ **An additional dissipation term**

Some perspectives

- ▶ Ions in water, hydrophilic and hydrophobic interactions
- ▶ Solute-solvent coupling in transport phenomena
- ▶ Glasses, silver halides, gaussian core particles, ionic liquids. . .
- ▶ Diffusion mechanism of an excess proton in water, in membranes

Table of contents

Introduction

Diffusion of a particle

Velocity field and density map

Boundary conditions

An extra dissipation

Velocity autocorrelation function

Diffusion mechanism

Water self diffusion

Conclusions