



# Probing mechanisms of $0\nu\beta\beta$ at the LHC

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# Outline

- Introduction
- Examples of  $0\nu\beta\beta$  mechanisms
- Information from  $0\nu\beta\beta$  experiments alone
- Complementary information from the LHC
- Single slepton production and  $0\nu\beta\beta$   
[Allanach, CHK, Päs arXiv:0902.4697,0903.0347](#)
- Summary

# What is $0\nu\beta\beta$

- $2n \rightarrow 2p + 2e^-$   $[(Z, A) (\rightarrow (Z + 1, A)^*) \rightarrow (Z + 2, A)]$
- Possible only in even-even nuclei, dominated by  $0^+ \rightarrow 0^+$  transition (e.g.  ${}^{76}_{32}\text{Ge} \rightarrow {}^{76}_{34}\text{Se}$ ).
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    - $\rightarrow$  absolute neutrino mass scale
    - $\rightarrow$  Majorana phases
  - Other LNV physics responsible ?

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  - Other LNV physics responsible ?
- $0\nu\beta\beta$  provides stringent limit to many (1st gen.) LNV processes.

# Neutrino mass measurements

## Neutrino Oscillations:

●  $|\Delta m_{23}^2| \simeq 2.5 \cdot 10^{-3} \text{eV}^2, \sin^2(2\theta_{23}) > 0.9$  SK,MINOS

$\Delta m_{12}^2 = 7.6 \cdot 10^{-5} \text{eV}^2, \tan^2(\theta_{12}) = 0.47$  SNO,KamLAND

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Absolute scale ?  $|m_{\beta\beta}| = |\sum_i U_{ei}^2 m_i|$ :

- $T_{1/2}^{0\nu\beta\beta}({}^{76}\text{Ge}) > 1.9 \cdot 10^{25} \text{yr} \rightarrow |m_{\beta\beta}| < 0.35 \text{eV}$  Heidelberg-Moscow  
(claimed observation  $|m_{\beta\beta}| \sim 0.5 \text{eV}$  Klapdor-Kleingrothaus et. al. 01 )
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## Related mass bounds:

- Cosmology  $\sum_i m_i < (0.17 - 2.0) \text{eV}$  Seljak,Slosar,McDonald 06
- Tritium decay  $m_{\nu_e} < 2.2 \text{eV}$  Mainz,Troitsk 01



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$$J = \bar{u} \mathcal{O}_\beta d, \quad j = \bar{e} \mathcal{O}_\beta \nu, \quad \beta = V \pm A, PS, T \dots$$

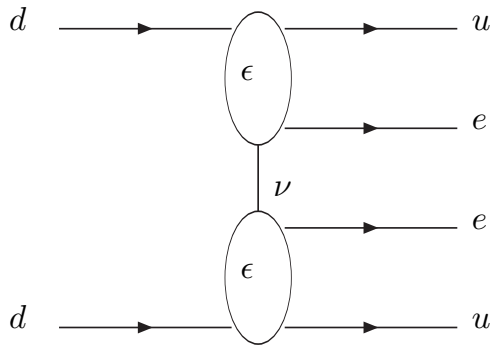
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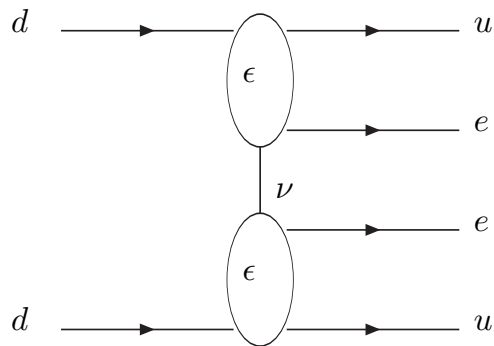
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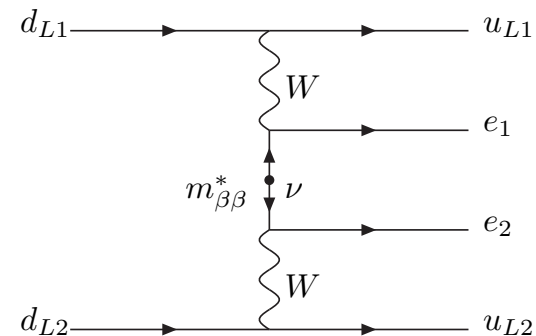
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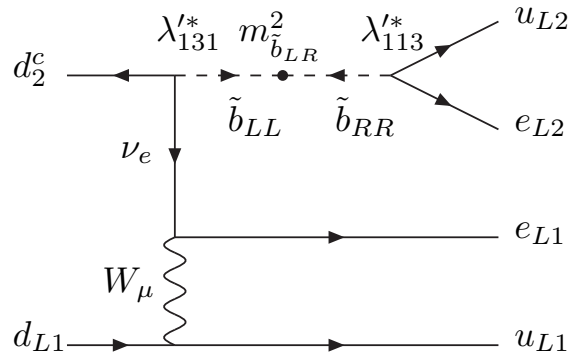
e.g.  $\longrightarrow$



$$\mathcal{L}_{EW}^{eff, \Delta L_e=2}(x) = \frac{G_F^2}{2} \left[ \bar{e}_1 \gamma_\mu (1 - \gamma_5) \frac{m_{\beta\beta}}{q^2} \gamma_\nu e_2^c \right] \times \left[ J_{1, V-A}^\mu(q) J_{2, V-A}^\nu(-q) \right]$$

# Examples

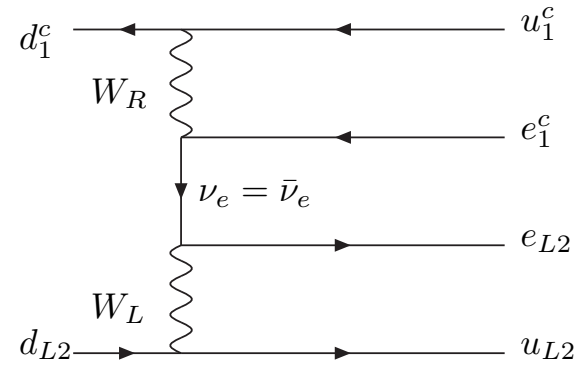
Sbottom exchange in LNV SUSY



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LR symmetric Model

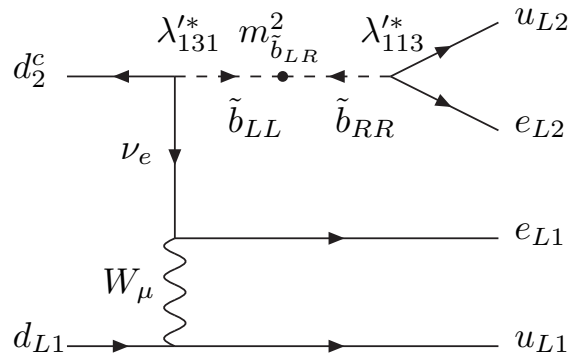


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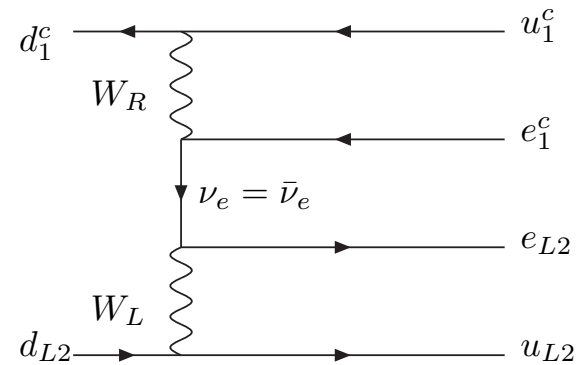
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- (Scalar) lepto-quark contributions similar to sbottom exchange.
- $W_R$ - $W_R$  contributions ...

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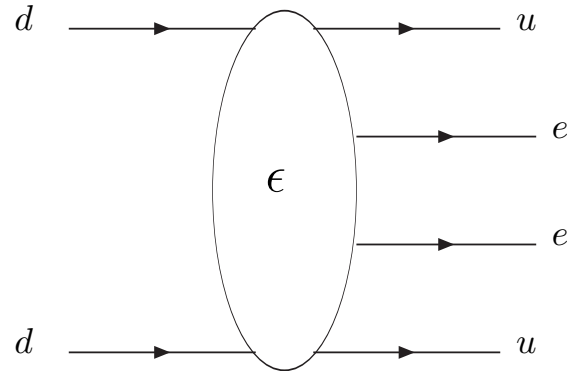
Short range interactions



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- Single point like interaction.

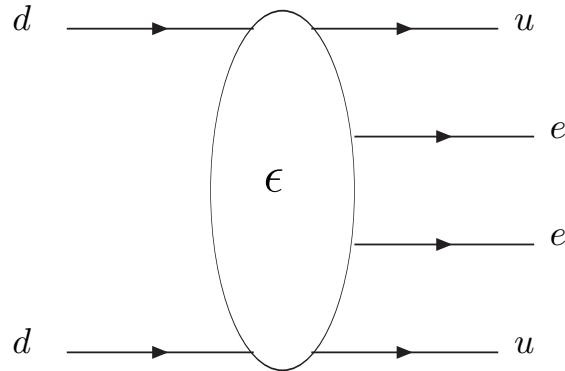


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- New physics (NP): LNV SUSY with gaugino exchange,  $[\mathcal{O}(\text{TeV})]$  heavy neutrinos...

# From NP to rates

Total rate (if dominated by one mechanism)

$$\Gamma_{0\nu\beta\beta} \propto \epsilon^2 G_{0\nu} |M_{0\nu}|^2$$

$\epsilon$ : LNV physics parameter,

$G_{0\nu}$ : precisely calculable phase space factor,

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Bottom line: **take the numerical values with a grind of salt...**

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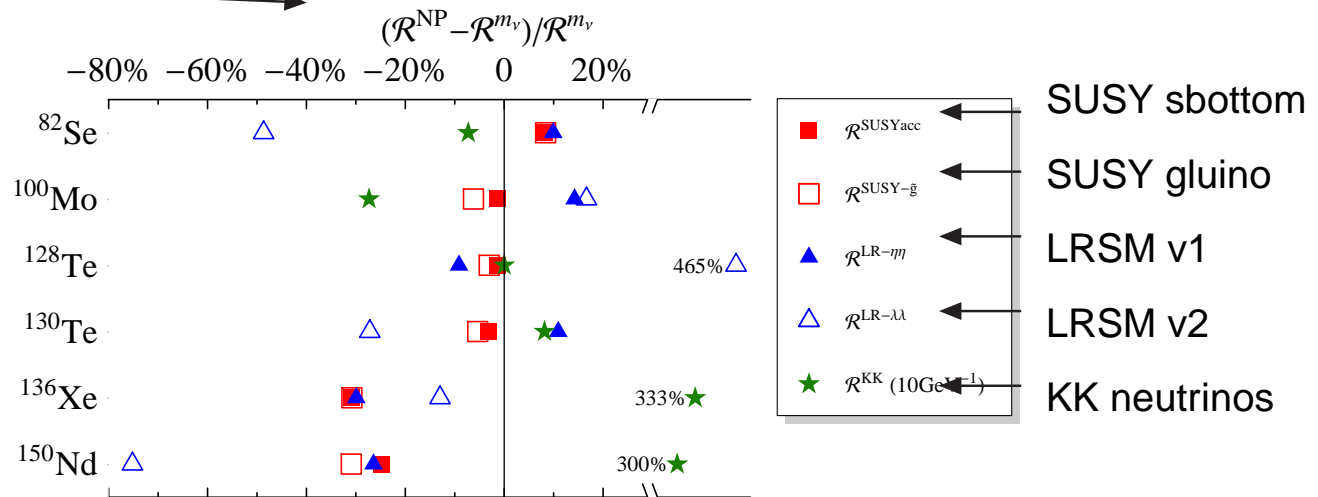
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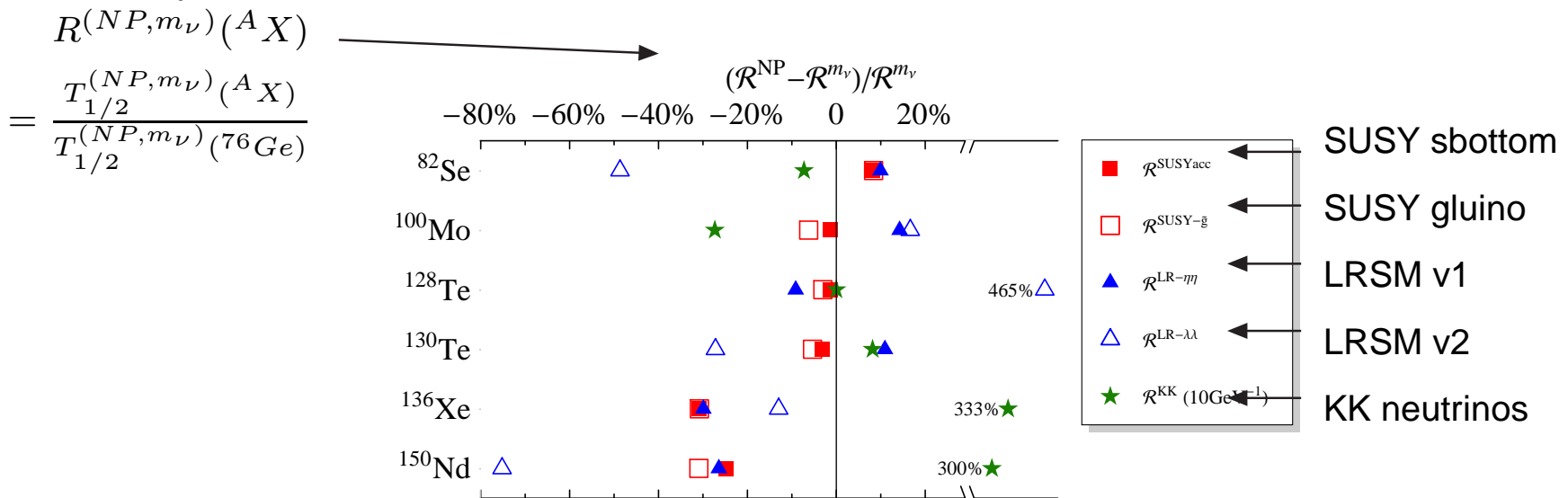
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- Many isotopes required.

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- E.g. SuperNEMO is sensitive to single electron kinematics.

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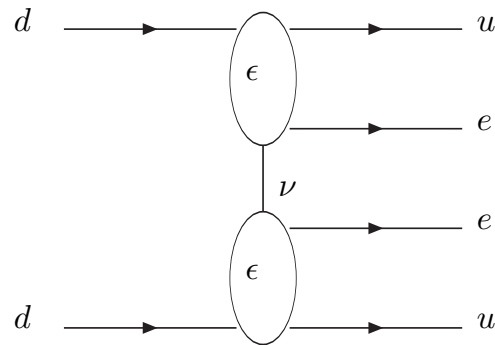
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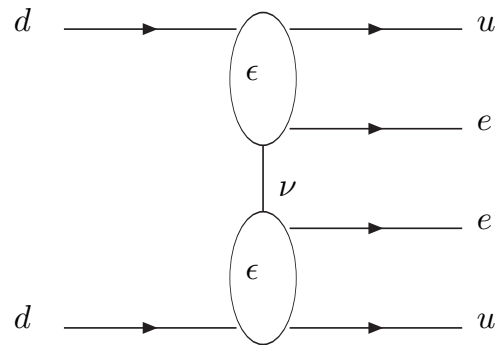
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$0\nu\beta\beta$  ?

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- Particularly suitable for short range interactions
- Focus on  $0\nu\beta\beta$  in LNV SUSY and signals at the LHC.

Allanach, CHK, Päs [arXiv:0902.4697](https://arxiv.org/abs/0902.4697), [0903.0347](https://arxiv.org/abs/0903.0347)

# $0\nu\beta\beta$ in LNV SUSY

LNV SUSY contains superpotential term:

$$W \supset \lambda'_{111} L_1 Q_1 D_1^c \rightarrow$$

$$\mathcal{L} \supset \lambda'_{111} (\bar{l}^c q \tilde{d}^c + \tilde{l} \bar{q}^c d^c + \bar{l}^c \tilde{q} d^c) + \dots$$

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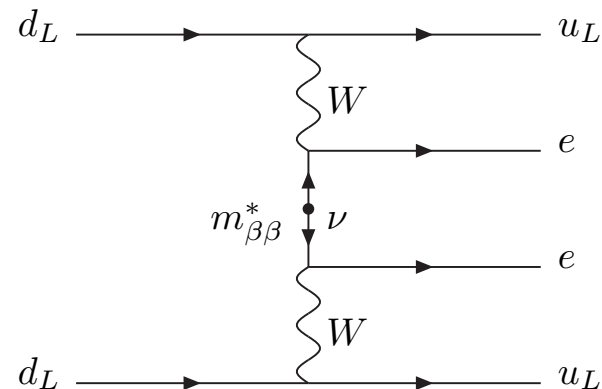
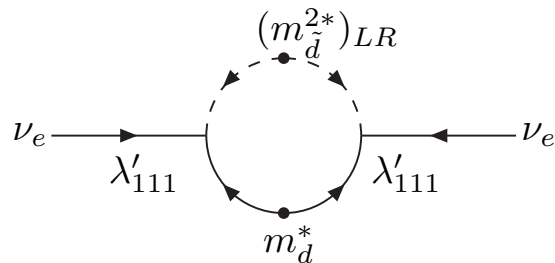
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Contribute to  $0\nu\beta\beta$  via:

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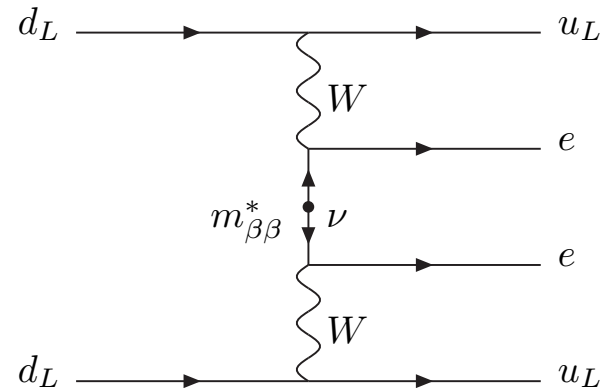
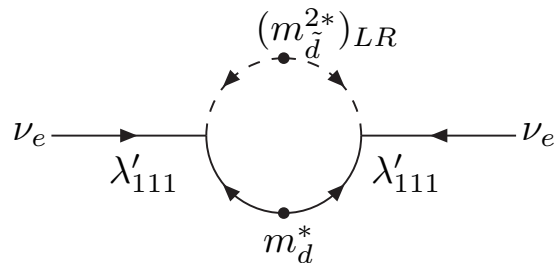
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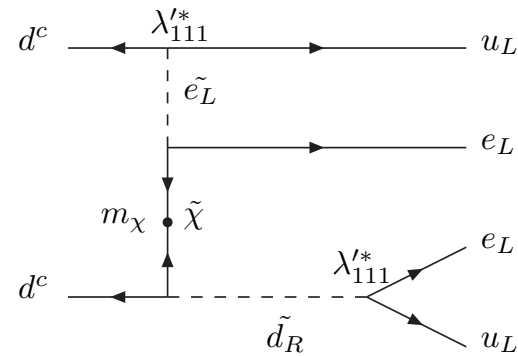
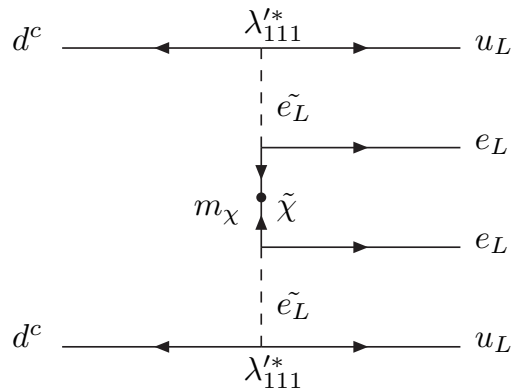
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- Also tree level  $m_\nu$  via RG effect.

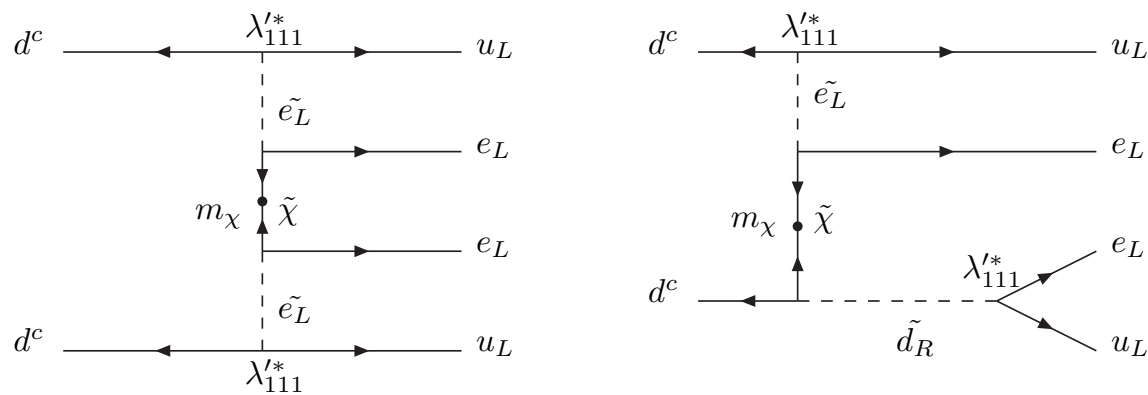
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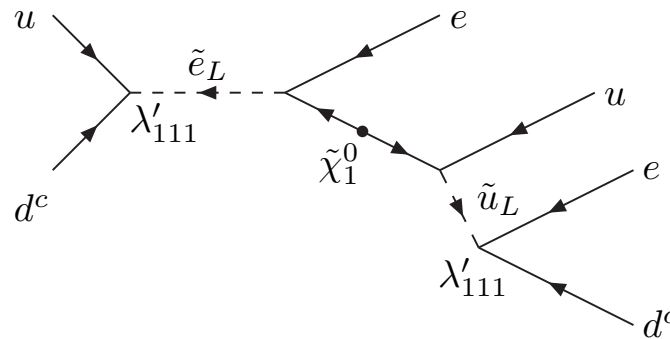
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$$\begin{aligned}
 \mathcal{L}_{\lambda'_{111}\lambda'_{111}}^{eff, \Delta L_e=2}(x) &= \frac{G_F^2}{2} m_p^{-1} [\bar{e}(1 + \gamma_5)e^c] \\
 &\times \left[ (\epsilon_{\tilde{g}} + \epsilon_{\chi})(J_{PS}J_{PS} - \frac{1}{4}J_T^{\mu\nu}J_{T\mu\nu}) + (\epsilon_{\chi}\tilde{e} + \epsilon'_{\tilde{g}} + \epsilon_{\chi\tilde{f}})J_{PS}J_{PS} \right] \\
 \epsilon_i &\sim \pi\alpha_{(S,W)} \frac{\lambda'_{111}{}^2}{G_F^2} \frac{m_P}{m_{(\tilde{g},\tilde{\chi})}} \frac{1}{m_{(\tilde{u},\tilde{d},\tilde{e})}^4}.
 \end{aligned}$$

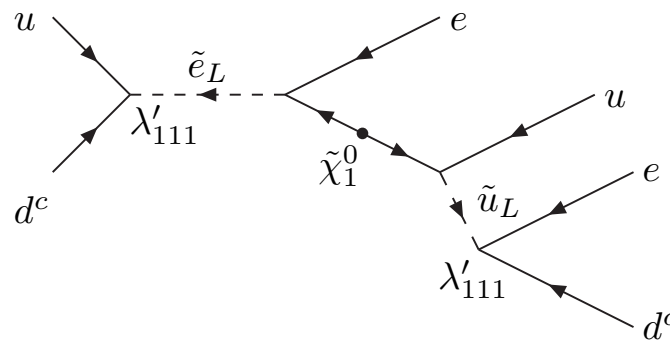
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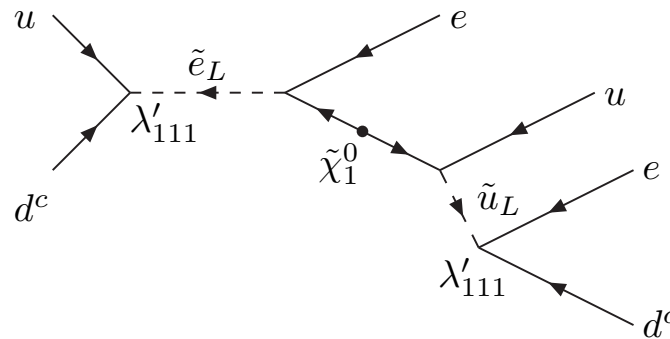
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- Dimension 9 operators:  
 $\lambda'_{111}$  bound relaxes rapidly with increasing  $\Lambda_{SUSY}$ .
- Lower  $T_{1/2}^{0\nu\beta\beta}({}^{76}\text{Ge})$  limit:  $\lambda'_{111} \lesssim 5 \cdot 10^{-4} \left(\frac{\Lambda_{SUSY}}{100\text{GeV}}\right)^{2.5}$ .  
Single slepton production:  $\sigma(pp \rightarrow \tilde{l}) \propto |\lambda'_{111}|^2 / m_{\tilde{l}}^3$   
→ *production upper limit increases with  $\Lambda_{SUSY}$ .*

# Single slepton production 2

- Signal: *Same* sign (SS) di-electron, 2 jets, little  $E_T$ .
- Previous analysis on SS di-muon signals for  $\lambda'_{211}$   
[Dreiner, Richardson, Seymour 99](#) .
- SS di-electron signal for  $\lambda'_{111}$  believed to be mathsize due precisely to 'stringent'  $0\nu\beta\beta$  bound !!
- To estimate  $\epsilon_{\lambda'_{111}}$ , need also  $\tilde{q}$ ,  $\tilde{\chi}$ ,  $\tilde{g}$  masses.  
Assumed to be obtained from other channels

# Model assumptions

MSSM model parameters:

- Masses defined by mSUGRA input  
 $m_0, M_{1/2}, A_0 = 0, \tan\beta = 10, \text{sgn}(\mu) = +1.$
- At SUSY scale, set  $\lambda'_{111}$  to be non-zero.
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NME model parameters ( $^{76}\text{Ge}$ ):

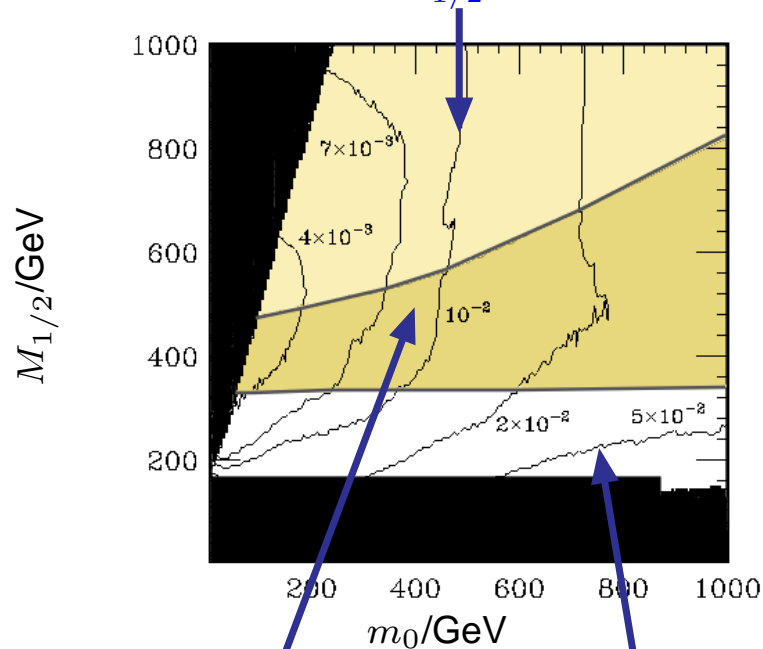
- Pion and nucleon modes included.  
$$M_{\lambda'_{111}} = \epsilon M_{\tilde{g}}^{2N} + \epsilon' M_{\tilde{f}}^{2N} + \left(\epsilon + \frac{5}{8}\epsilon'\right) \left(\frac{4}{3}M^{1\pi} + M^{2\pi}\right)$$
- $M_{\tilde{g}}^{2N} = 283, M_{\tilde{f}}^{2N} = 13.2, M^{1\pi} = -18.2, M^{2\pi} = -601$

Hirsch,Klapdor-Kleingrothaus,Simkovic 96 , Faessler,Kovalenko,Simkovic 98

# Results

Infer  $T_{1/2}^{0\nu\beta\beta}({}^{76}\text{Ge})$  from SS di-electron  $5\text{-}\sigma$  signal at  $10\text{ fb}^{-1}$ , 14 TeV (no  $m_{RR}$  contribution):

$$1 \cdot 10^{27} \text{ yrs} < T_{1/2}^{0\nu\beta\beta}({}^{76}\text{Ge})$$



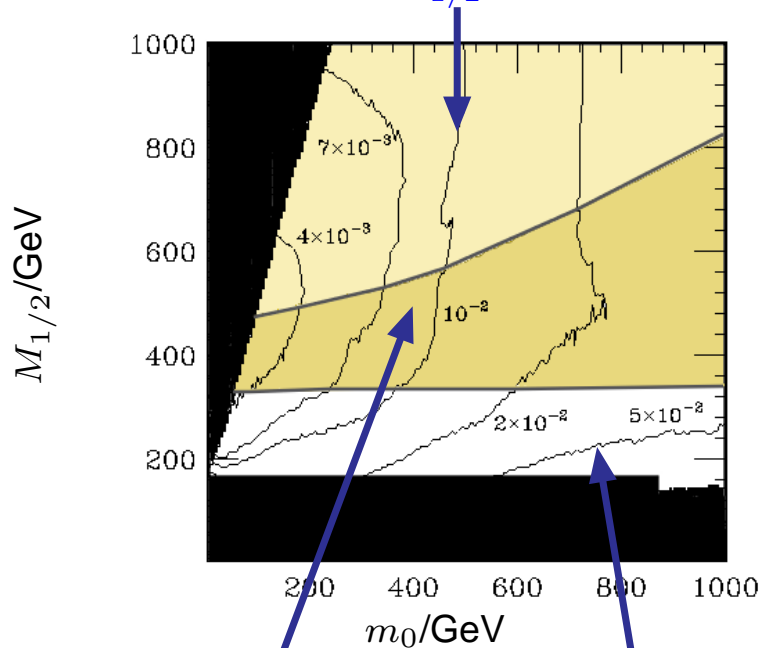
$$1.9 \cdot 10^{25} < T_{1/2}^{0\nu\beta\beta}({}^{76}\text{Ge}) < 1 \cdot 10^{27} \text{ yrs}$$

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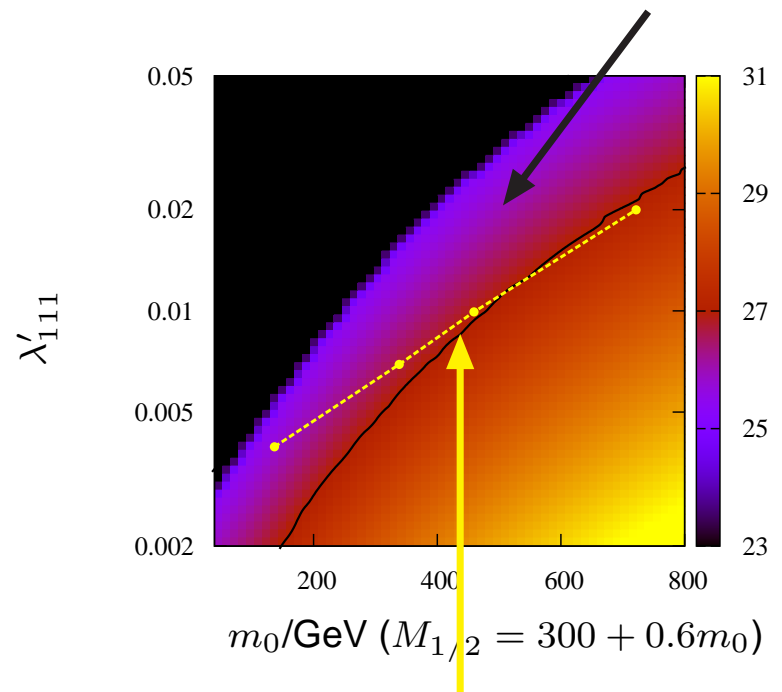
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$$T_{1/2}^{0\nu\beta\beta}({}^{76}\text{Ge}) < 1.9 \cdot 10^{25} \text{ yrs}$$

Sensitive to future  $0\nu\beta\beta$  expts.



LHC single slepton discovery limit

# Results 2

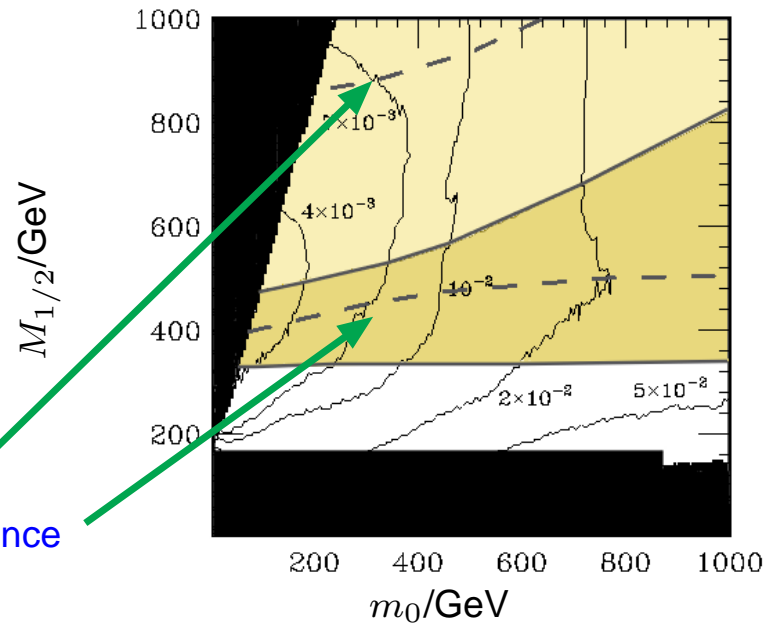
Including

$$|m_{\beta\beta}| = 0.05 \text{ eV}$$

$$(\sim \sqrt{\Delta m_{23}^2})$$

Constructive interference

Destructive interference

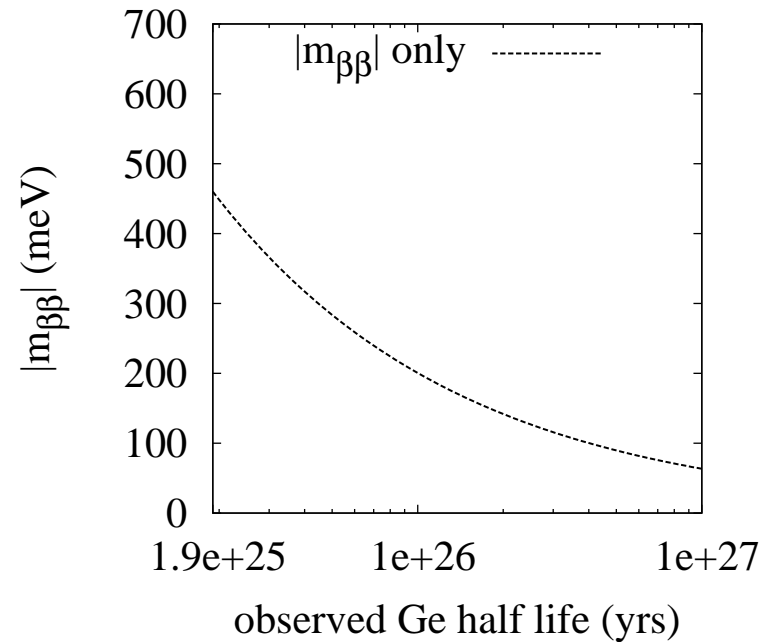
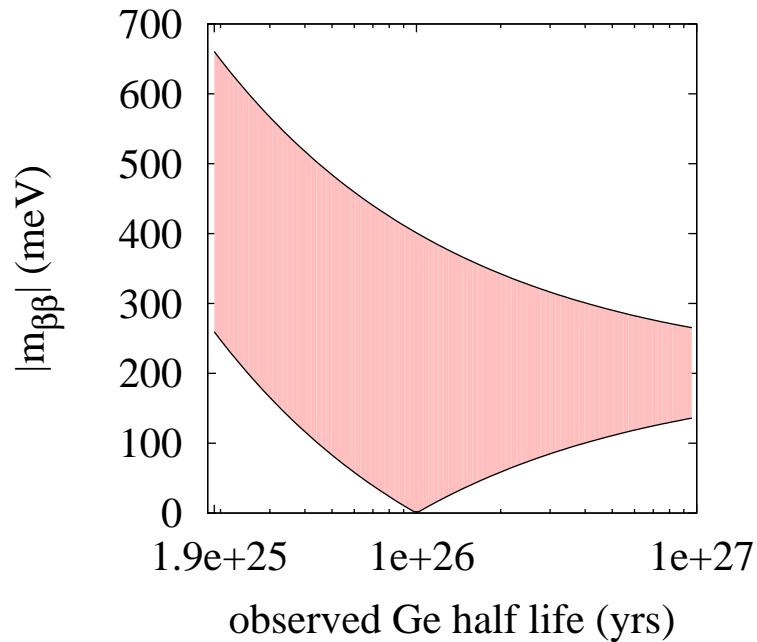


- Destructive interference with  $m_{\beta\beta}$  increases  $T_{1/2}^{0\nu\beta\beta}({}^{76}\text{Ge}) \rightarrow$  dark yellow region shrinks.
- Fixing  $T_{1/2}^{0\nu\beta\beta}({}^{76}\text{Ge})$ , destructive int. with  $m_{\beta\beta}$  increases SSL rate  $\rightarrow$  better SSL discovery prospect.

# Inference on $m_{\beta\beta}$

Given  $5\sigma$  SSL observation ( $M_0 = 680\text{GeV}$ ,  $M_{1/2} = 440\text{GeV}$ )

$\rightarrow T_{1/2}^{0\nu\beta\beta}({}^{76}\text{Ge}) = 1 \cdot 10^{26}\text{yrs}$  if direct contribution only.



- Band of  $m_{\beta\beta}$  depending on relative phase.
- Normal hierarchy possible if  $0\nu\beta\beta$  observed.

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- LHC allows direct determination of LNV parameters in  $0\nu\beta\beta$ .
- Discussed interplay between  $0\nu\beta\beta$  and single slepton production at the LHC in LNV SUSY.
- If  $0\nu\beta\beta$  is discovered, single slepton production could test the  $\lambda'_{111}$  hypothesis.

# Planned $0\nu\beta\beta$ experiments

Need  $\sim 1$  ton of active isotopes to investigate the IH.  
An incomplete list of experiments here:

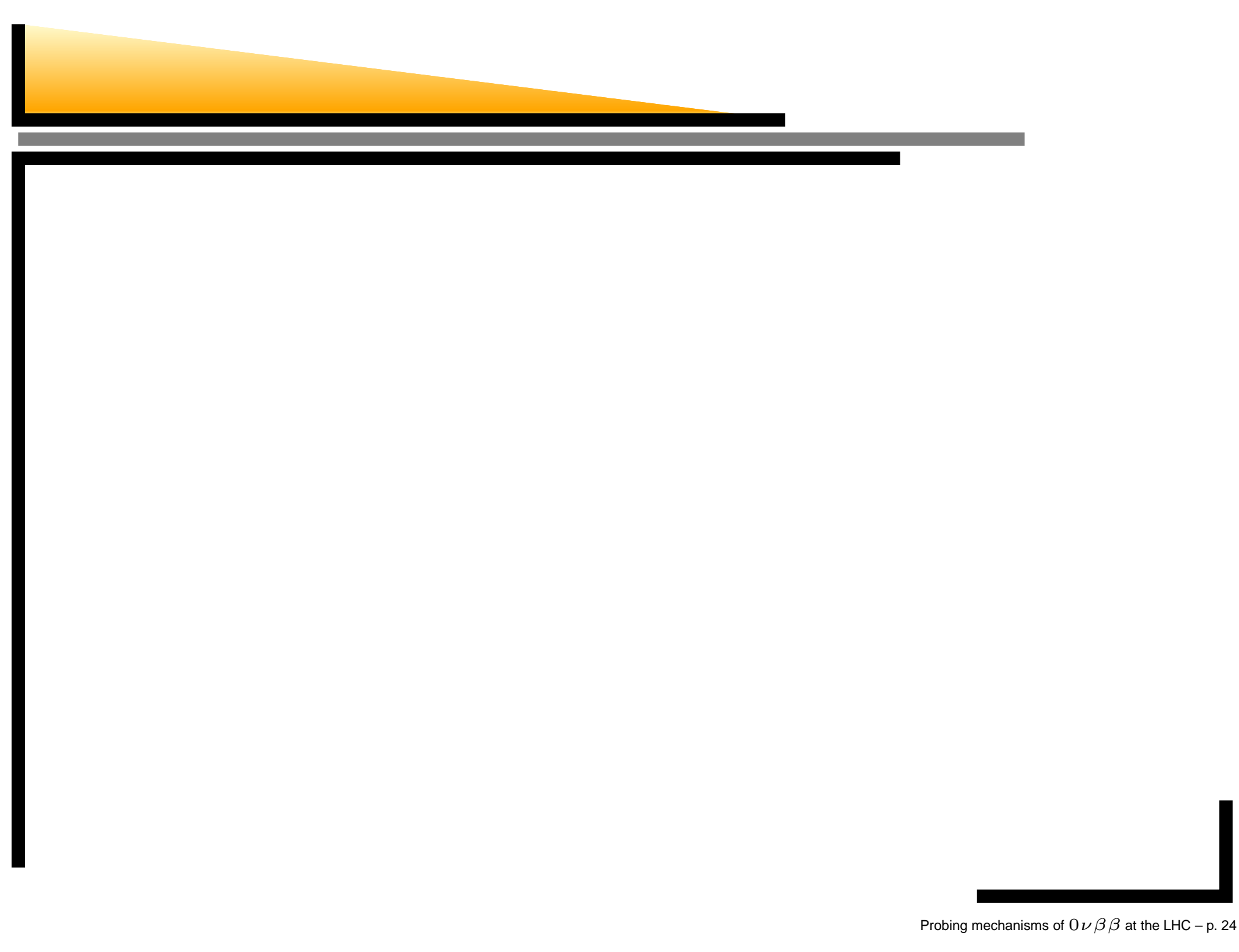
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CUORE( $^{130}\text{Te}$ )	750kg, $\text{TeO}_2$ bolometers	Setup completion in 2012
GERDA( $^{76}\text{Ge}$ )	18kg (Phase I), Ge detector	Construction being completed
MAJORANA( $^{76}\text{Ge}$ )	60kg (initial phase), Ge detector	R&D
EXO( $^{136}\text{Xe}$ )	1-10 ton, TPC, $\text{Ba}^{++}$ tagging	EXO-200 run expected in 2009
SuperNEMO( $^{82}\text{Se}$ )	100kg, track + calor.	1st module in 2011. Rest in 2012

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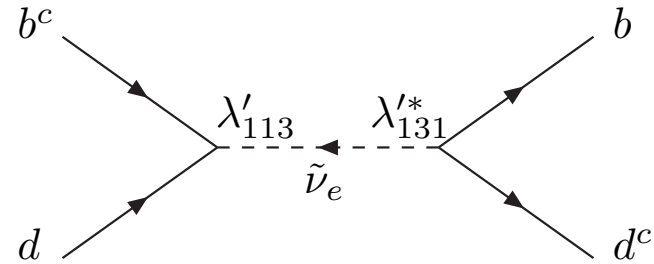
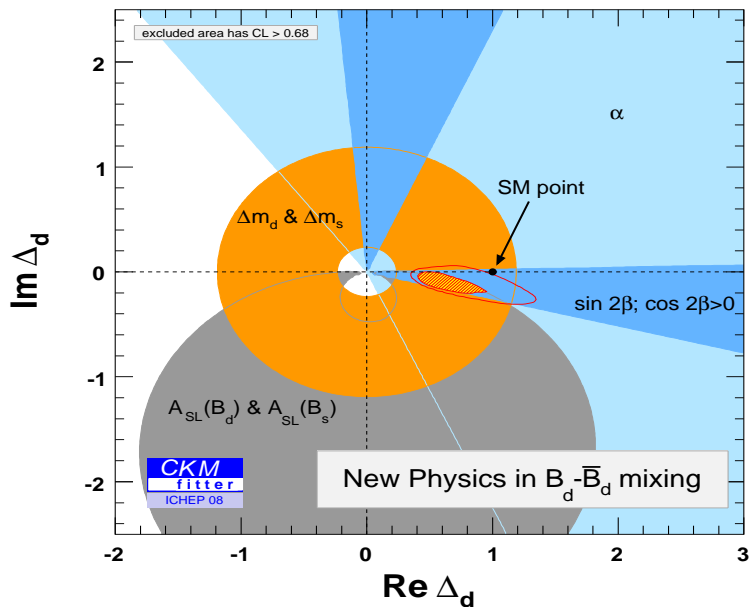


# Thankyou



# $B_d^0-\bar{B}_d^0$ mixing and $0\nu\beta\beta$

$$B_d^0-\bar{B}_d^0 \text{ mixing limit: } \langle B_d | M_{12}^{\text{SM+New Physics}} | \bar{B}_d \rangle = \Delta_d \langle B_d | M_{12}^{\text{SM}} | \bar{B}_d \rangle$$



$$\lambda'_{113} \lambda'_{131} \leq 4.0 \cdot 10^{-8} \frac{m_{\tilde{\nu}_e}^2}{(100\text{GeV})^2}$$

- For  $0\nu\beta\beta$ ,  $\lambda'_{113} \lambda'_{131} \lesssim 2 \cdot 10^{-8} \left( \frac{\Lambda_{\text{SUSY}}}{100\text{GeV}} \right)^3$
- $B_d^0-\bar{B}_d^0$  bound more stringent than  $0\nu\beta\beta$  with high  $\Lambda_{\text{SUSY}}$ , though with different SUSY mass dependence.

# Nuclear structure model

Examples of approximations used:

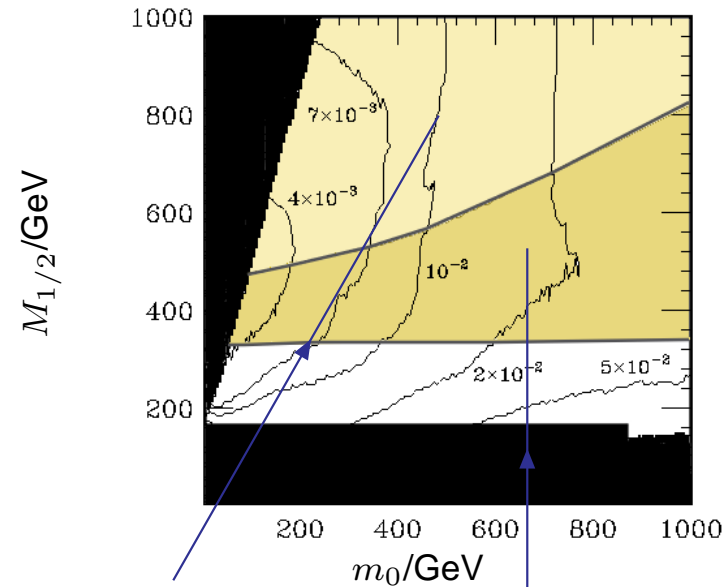
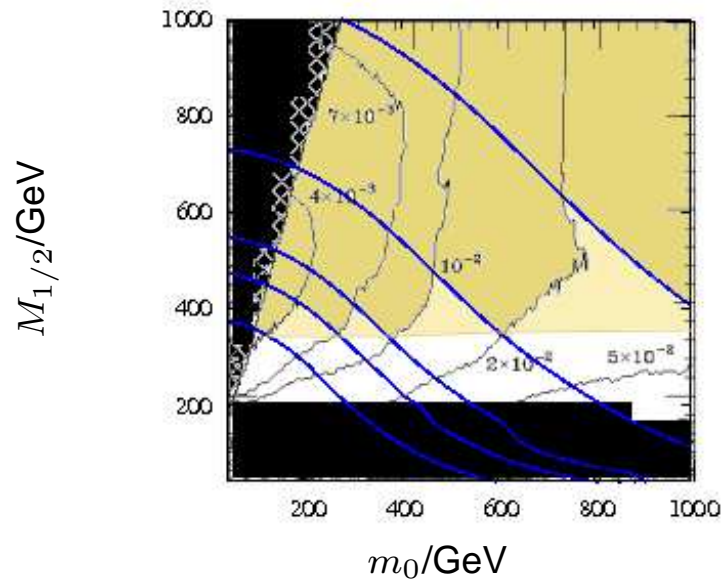
- Mean field theory, truncated single particle basis.
- Energy of states, interactions adjusted to reproduce data.
- Quenching of the axial vector coupling  $g_A$ .
- Short range correlations included explicitly.
- QRPA: quasiboson approx., pairing parameter  $g_{pp}$  adjusted...

More advanced tools available in 'simpler' nuclear properties ...



# Comparing $\lambda'_{111}$ bounds

- Infer  $T_{1/2}^{0\nu\beta\beta}({}^{76}\text{Ge})$  from SS di-electron 5- $\sigma$  discovery reach at 10 fb<sup>-1</sup>.



$$1 \cdot 10^{27} \text{ yrs} < T_{1/2}^{0\nu\beta\beta}({}^{76}\text{Ge})$$

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# LHC SS di-lepton cuts

From [Dreiner, Richardson, Seymour 99](#)

- Lepton  $|\eta| < 2.0$ .
- Lepton  $p_T > 40$  GeV.
- Isolation cut:  $E_T < 5$  GeV in cone  $R=0.4$ .
- Reject  $65 < M_T < 80$  GeV.
- $\cancel{E}_T < 20$  GeV.
- OSSF lepton veto.
- No more than 2 jets, each with  $p_T > 50$  GeV.