Radiative decays at LHCb

Luis Miguel Garcia Martin

EPP Warwick Seminar January 14th, 2021







Introduction

- Theoretical motivation
- LHCb Detector
- Experimental challenges

Analyses

- $B_s \to \phi \gamma$ time-dependent analysis
- $B \to K^* ee$ angular analysis
- $\Lambda_b^0 \to \Lambda \gamma$ (angular) analysis
- Ongoing analyses

3 Conclusions and prospects



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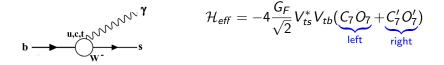
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Introduction: Theoretical motivation

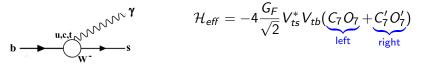
The $b \rightarrow s\gamma$ process is forbidden at tree level in the Standard Model (SM). Indirect searches grant access to larger energy scales than direct ones. At LO in SM only O_7 and O'_7 contribute





Introduction: Theoretical motivation

The $b \rightarrow s\gamma$ process is forbidden at tree level in the Standard Model (SM). Indirect searches grant access to larger energy scales than direct ones. At LO in SM only O_7 and O'_7 contribute



Wilson coefficient can be constrained through measurement of:

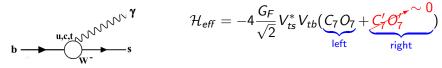
- \bullet Branching ratio: $\mathcal{B}_{\mathsf{rad}} \propto |\mathit{C}_7|^2 + |\mathit{C}_7'|^2$
- Photon polarization: $lpha_{\gamma}^{LO}=rac{1-|rac{C_{\gamma}'}{C_{7}}|^{2}}{1+|rac{C_{\gamma}'}{C_{7}}|^{2}}$

• CP asymmetry:
$$A_{CP} \propto Im rac{C_7 C_7'}{|C_7|^2 + |C_7'|^2}$$

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Introduction: Theoretical motivation

Photons in such transitions are mainly **left-handed in the SM** since the W boson couples to left-handed quarks.

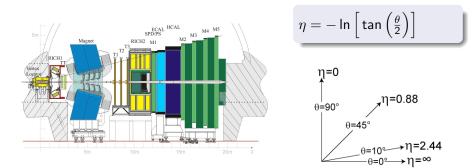


Wilson coefficient can be constrained through measurement of:

- Branching ratio: $\mathcal{B}_{rad} \propto |C_7|^2 + \int C_7^{2} e^{-C_0}$
- Photon polarization: $\alpha_{\gamma}^{LO} = \frac{1 |\frac{C_{\gamma}^{\prime}}{C_{\gamma}}|^2}{1 + |\frac{C_{\gamma}^{\prime}}{C_{\gamma}}|^2}$
- CP asymmetry: $A_{CP} \propto Im \frac{C_7 C_1}{|C_7|^2 + |C_1'|^2}$



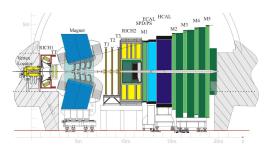
- One of the four detector at LHC
- LHCb is a single-arm (2 $<\eta<$ 5) spectrometer
 - Optimised for beauty and charm decays
- Runs at lower luminosity
 - Optimised for precision measurements



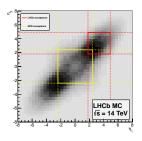
 $\boldsymbol{\theta}:$ Angle between \boldsymbol{p} and positive beam axis.



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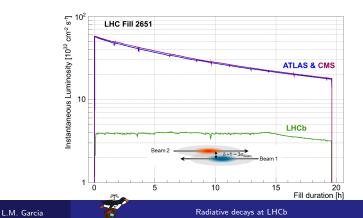


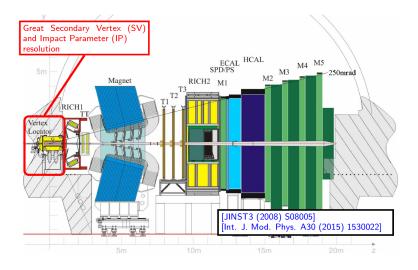
$$\eta = -\ln\left[\tan\left(\frac{\theta}{2}\right)\right]$$





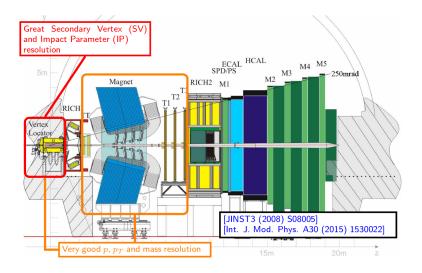
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$\Delta IP = (16+29/p_T \text{ [GeV]}) \ \mu \text{m}$

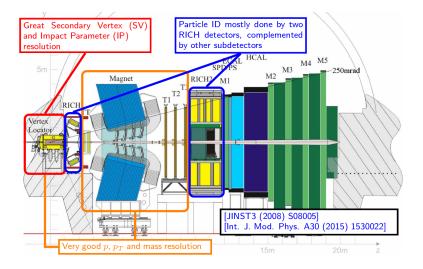




$\Delta p/p = 0.5 - 1.0\%$

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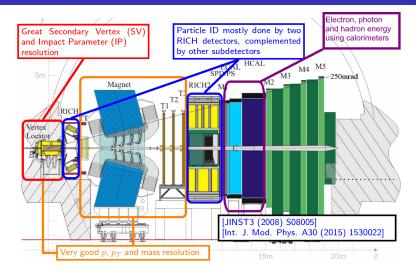




Kaon ID \sim 95% for \sim 5% $\pi \rightarrow K$ mis-id probability



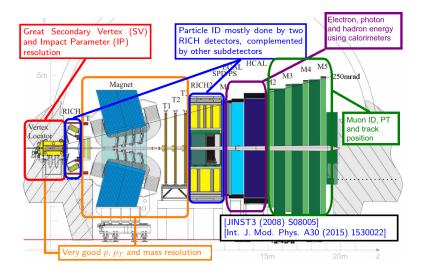




Electron ID $\sim 90\%$ for $\sim 5\%~e \rightarrow h$ mis-id probability $\Delta E/E_{\rm ECAL} = 1\% + 10.0\%/\sqrt{E[{\rm GeV}]}$

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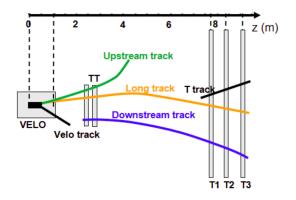




Muon ID \sim 97% for \sim 1-3% $\pi ightarrow \mu$ mis-id probability







Long tracks

• Hits at least in VELO and T stations

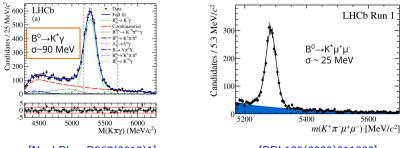
• Used in majority of analyses

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Downstream tracks

- Hits in TT and T stations (not in VELO)
- Decay products of long-lived particles





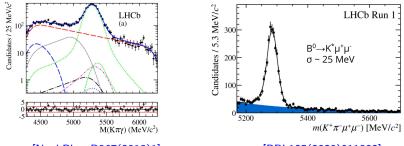
[Nucl.Phys.B867(2013)1]

[PRL125(2020)011802]

Challenges for analysis involving neutrals (γ and π^0):

- Photon direction not reconstructed:
 - Mass resolution dominated by photon momentum
 - Large background ($\sim 10\,\gamma/{\rm events},~{\rm merge}~\pi^0 \to \gamma\gamma)$
- Rare decays \implies low signal yield $\left(\mathcal{B} \sim O(10^{-5})\right)$





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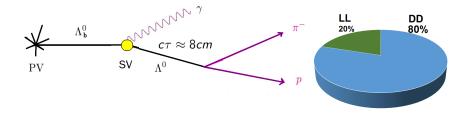
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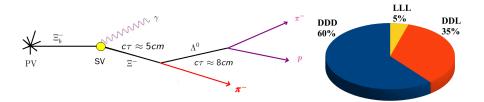


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 - Decay after the VELO
 - $\bullet~$ Worse IP/vertex position resolution
 - Hlt1 (trigger) only selects Long tracks







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Time-dependent decay rates of $B_s \to \phi \gamma$ and $\overline{B_s} \to \phi \gamma$ grant access to photon polarization:

$$\Gamma_{B_{s} \to \phi\gamma}(t) \propto e^{-\Gamma_{s}t} \left[\cosh\left(\frac{\Delta\Gamma_{s}t}{2}\right) - \mathcal{A}_{\phi\gamma}^{\Delta} \sinh\left(\frac{\Delta\Gamma_{s}t}{2}\right) + \mathcal{C}_{\phi\gamma} \cos\left(\Delta m_{s}t\right) - \mathcal{S}_{\phi\gamma} \sin\left(\Delta m_{s}t\right) \right]$$

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- $\mathcal{A}^{\Delta}_{\phi\gamma}$ and $S_{\phi\gamma}$ are sensitive to photon polarization
- $C_{\phi\gamma}$ is related to direct CP violation
- SM prediction close to zero for ${\cal A}^{\Delta}_{\phi\gamma}$, ${\cal C}_{\phi\gamma}$ and ${\cal S}_{\phi\gamma}$
- \mathcal{A}^{Δ} only accessible for B_s decays :
 - $\Delta\Gamma_{s}\sim0.081\pm0.011$

•
$$\Delta\Gamma_d\sim 0$$



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- Previous: $\mathcal{A}^{\Delta}_{\phi\gamma}$ measured in untagged analysis with Run I data at LHCb [LHCb: PRL118(2017)021801]

*Untag: No separation between B_s and $\overline{B_s}$



Time-dependent decay rates of $B_s \to \phi \gamma$ and $\overline{B_s} \to \phi \gamma$ grant access to photon polarization:

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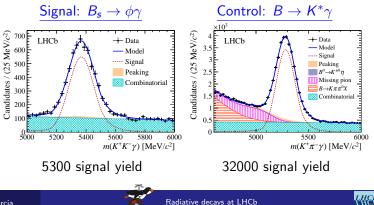
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- Previous: $\mathcal{A}^{\Delta}_{\phi\gamma}$ measured in untagged analysis with Run I data at LHCb [LHCb: PRL118(2017)021801]
- New: $S_{\phi\gamma}$ and $C_{\phi\gamma}$ measurement using tagging [PRL123(2019)081802]

*Tagging: Separation between B_s and $\overline{B_s}$ [JINST11(2016)P05010]



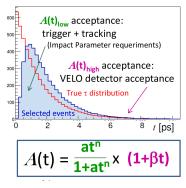


- Mass fit of $B_s o \phi \gamma$ (signal) and $B o K^* \gamma$ (control) decays
- Using Run 1 data at LHCb [PRL123(2019)081802]
- Background subtracted with sPlot technique, fitting the B mass



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- Background subtracted with sPlot technique, fitting the B mass
- Decay time measured from B momentum and flight distance
- Need to control the proper time acceptance

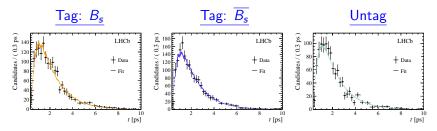




Time-dependent analysis of $B_s \rightarrow \phi \gamma$: Proper time fit

Analysis strategy:

- Simultaneous unbinned ML fit to $B_s \to \phi \gamma$ (signal) and $B \to K^* \gamma$ (control) channels
- Mis-tag probability and resolution evaluated per event



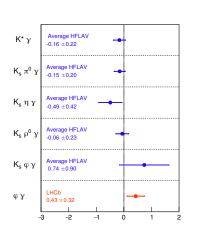
 $egin{aligned} S_{\phi\gamma} &= 0.43 \pm 0.30 \pm 0.11, \ C_{\phi\gamma} &= 0.11 \pm 0.29 \pm 0.11, \ \mathcal{A}^{\Delta}_{\phi\gamma} &= -0.67 \, {}^{+0.37}_{-0.41} \pm 0.17 \end{aligned}$

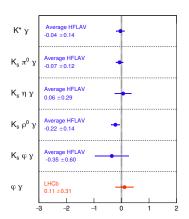
- $\bullet\,$ Compatible with SM at 1.3, 0.3, 1.7 $\sigma\,$
- First measurement of S and C in the $B_{\rm s}
 ightarrow \phi \gamma$ decay [PRL123(2019)081802]



S_{CP} and C_{CP} in $b \rightarrow s\gamma$ transitions

 S_{CP}





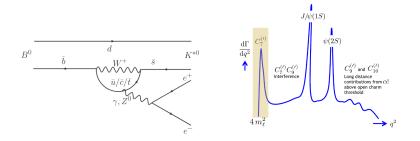
 C_{CP}

• Measurement competitive with other results from *b*-factories



$B \rightarrow K^* ee$ angular analysis

Decay dominated by $b
ightarrow s \gamma$ at very-low $q^2 = m_{
m ee}^2$ pole



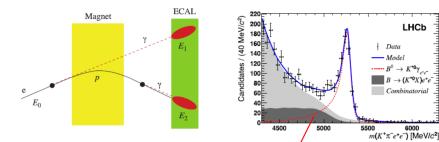
- Virtual γ decaying in an observable $\ell^+\ell^-$ pair
- Previous analysis with Run 1 data [JHEP04(2015)064]
- Recent update including Run 1 + Run 2 data [JHEP12(2020)081]
- All final state particles are charged $(K^+\pi^-e^+e^-)$

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$B \rightarrow K^* e e$ angular analysis

Electrons loose energy by radiation (bremsstrahlung)

- difficult to reconstruct
- need bremsstrahlung recovery



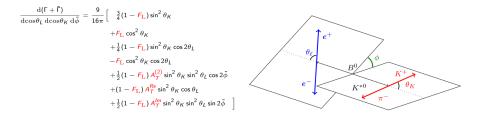
 \rightarrow adding neutral clusters from the ECAL , with $E_T > 75 MeV$

Long radiative tail in the B mass distribution: controlled from $B \rightarrow K^*\gamma$ events ($\gamma \rightarrow e^-e^+$, with bremsstrahlung emission)



$B \rightarrow K^* ee$ angular analysis

Angular distribution:



- Angular distribution with three angles: $\cos \theta_K, \cos \theta_L, \phi$
- Angular observables granting access to the photon polarization:

ve decays at LHCb

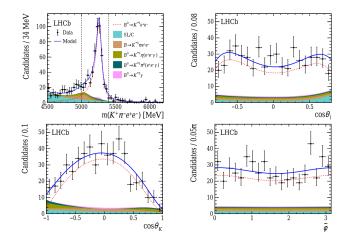
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•
$$A_T^{(2)}(q^2 \to 0) = \frac{2Re(C_7C_7^{/*})}{|C_7|^2 + |C_7^{/}|^2}$$

•
$$A_T^{lm}(q^2 \to 0) = rac{2lm(C_7C_7^{\prime*})}{|C_7|^2 + |C_7^{\prime}|^2}$$

$B \rightarrow K^* ee$ angular analysis

Fit to the B mass and angles (in reduced mass region)



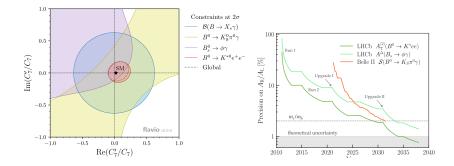
[JHEP12(2020)081]

Radice decays at LHCb

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$B \rightarrow K^* ee$ angular analysis : Constraint to C_7'

- This analysis set strong constraints in the C_7 C'_7 plane
- Still statistically limited
- precision will improve with Upgrade

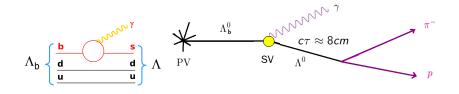


[JHEP12(2020)081]

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Radiative b-baryon decays:

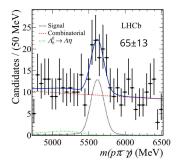
- Non-zero spin grants access to more observables
- Two spectator quarks \implies different form factors
- Photon polarization has never been measured!!
- b-baryons only at accesible pp colliders (LHC)



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First observation: $\Lambda^0_b \to \Lambda \gamma$

First observation of a radiative *b*-baryon decay using 2016 data ($\mathcal{L} = 1.67 \text{ fb}^{-1}$) [PRL123(2019)031801]:



$$\mathcal{B}(\Lambda_b^0
ightarrow \Lambda\gamma) = (7.1 \pm 1.5 \pm 0.6 \pm 0.7) imes 10^{-6}$$

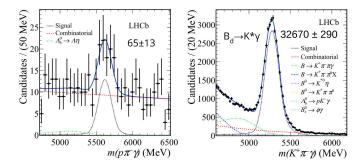
Radiative decent at LHCb

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- Observed with 5.6 σ significance
- Open door to photon polarization measurement

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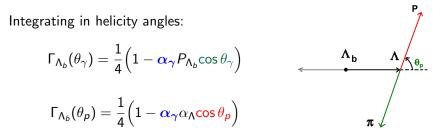
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Angular distribution: $\Lambda_b \rightarrow \Lambda \gamma$

Angular distribution for Λ_b decay:

$$\Gamma_{\Lambda_b}(\theta_{\gamma},\theta_{p}) = 1 - \alpha_{\Lambda} P_{\Lambda_b} \cos \theta_{p} \cos \theta_{\gamma} - \alpha_{\gamma} \left(\alpha_{\Lambda} \cos \theta_{p} - P_{\Lambda_b} \cos \theta_{\gamma} \right)$$



The decay parameters are:

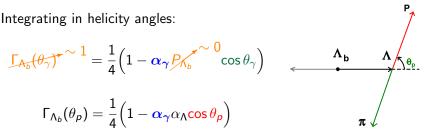
- $P_{\Lambda_b} = 0.00 \pm 0.06(stat) \pm 0.06(sys)$ [PhysRevD.97(2018)072010]
- $\alpha_{\Lambda} = 0.732 \pm 0.010$ [PDG 2020]



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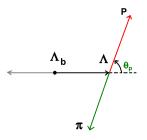
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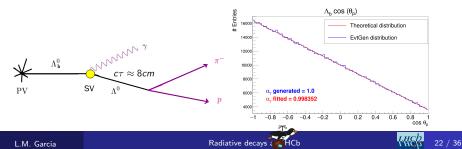
Angular distribution for Λ_b decay:

$$\Gamma_{\Lambda_b} = \frac{1}{4} \left(1 - \frac{\alpha_{\gamma} \alpha_{\Lambda} \cos \theta_{p}}{4} \right)$$

The value of the Λ decay parameter is:

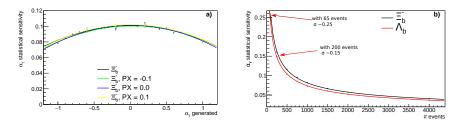
• α_Λ = 0.732 ± 0.010 [PDG 2020]





Statistical sensitivity to photon polarization

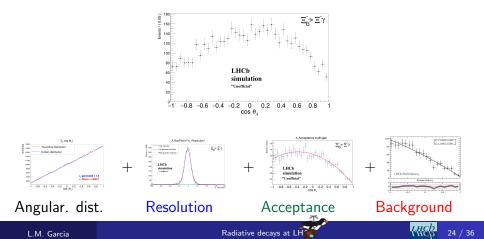
Sensitivity to the photon polarization (α_{γ}) [EPJC(2019)79:634]:



- Measuring P_{H_b} as well does not reduce the sensitivity to α_γ
 - For a large data sample
- Similar α_{γ} sensitivity using Λ_b and Ξ_b^- systems
- Only considering the theoretical angular polarization:
 - $\sigma_{\alpha} \sim 0.25$ with $\Lambda_b^0 \to \Lambda \gamma$ 2016 data sample (65 events)
 - $\sigma_{\alpha} \sim 0.15$ with $\Lambda_b^0 \to \Lambda \gamma$ Run2 data sample (assuming 200 events)



$$\Gamma(\theta_{\Lambda},\theta_{\rho};\boldsymbol{\alpha}_{\gamma}) = \left(Signal(\theta_{\Lambda},\theta_{\rho};\boldsymbol{\alpha}_{\gamma}) \times A(\theta_{\Lambda},\theta_{\rho})\right) * Res(\theta_{\Lambda},\theta_{\rho}) + \frac{S}{B}\left(BKG(\theta_{\Lambda},\theta_{\rho})\right)$$

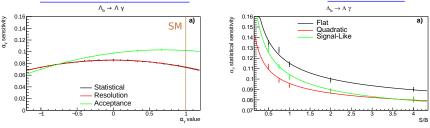


Sensitivity studies

Independently testing the effects:

- Statistical: Shape of theoretical angular distribution
- Resolution: Effect of the detector
- Acceptance: Effect of the selection
- Studying the effect of the background:
 - with different angular background shapes
 - for several signal over background ratios

Acceptance/Resolution



Tested for 1k events [EPJC(2019)79:634]

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Radiative decays at LHCb



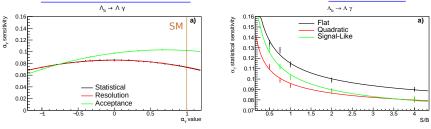
Background

Sensitivity studies

Independently testing the effects:

- Statistical: Shape of theoretical angular distribution
- Resolution: Effect of the detector (negligible effect)
- Acceptance: Effect of the selection (asymmetric with α_{γ})
- Studying the effect of the background:
 - with different angular background shapes
 - for several signal over background ratios

Acceptance/Resolution



Tested for 1k events [EPJC(2019)79:634]

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Background

Fit strategy: $\Lambda_b \rightarrow \Lambda \gamma$

The photon polarization will be extracted using the PDF:

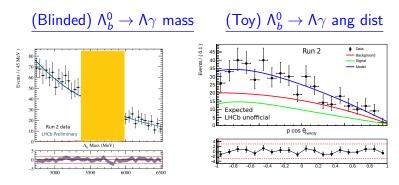
$$\Gamma(\theta_{p}; \alpha_{\gamma}) = \frac{S}{S+B} \Big(Signal(\theta_{p}; \alpha_{\gamma}) \times A(\theta_{p}) \Big) + \frac{B}{S+B} \Big(BKG(\theta_{p}) \Big)$$

• Using Run 2 data in Λ_b^0 mass_{PDG} $\pm\,2.5\sigma$

• Signal:
$$\frac{1}{4} \left(1 - \boldsymbol{\alpha}_{\gamma} \alpha_{\Lambda} \cos \theta_{p} \right)$$

- Resolution: Negligible effect
- Acceptance: Extracted from MC, controlled from data $(\Lambda_b \rightarrow \Lambda J/\psi)$
- Background: Extracted from data mass side-bands
- Signal fraction is Gaussian constrained
- Sensitivity to α_{γ} studied in [EPJC(2019)79:634]
- Expected $\sigma_{\gamma}(\text{stat.}) \sim 0.32 \text{ (WIP)}$

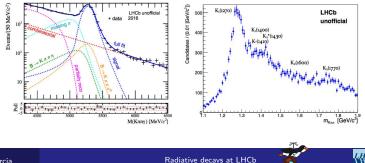
$\Lambda^0_b ightarrow \Lambda\gamma$ analysis with Run 2 data



- Reconstruction and selection strategy defined using Run 2
- Background sources (other than combinatorial) found negligible
- Angular fit procedure established
 - Fitting signal and background, and including acceptance effects
- Systematics evaluated and found to be below 0.15
- Results (still blinded) will be delivered very soon!

$B \rightarrow K^+ \pi^+ \pi^- \gamma$ amplitude analysis

- First analysis observing a non-zero photon polarization at LHCb [PRL112(2014)161801]
- Inteference with reasonance grant access to the photon polarization [EPJC79(2019)622]
- Using amplitude analysis to provide a direct measurement
- Expecting an statistical uncertainty of 0.014 [EPJC79(2019)622]



Search for $\Xi_b^- \to \Xi^- \gamma$

First search of the $\Xi_b^- \to \Xi^- \gamma$ decay

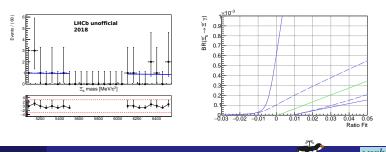
• No previous limit:

•
$$\mathcal{B}(\Xi_b^- \to \Xi^- \gamma)_{\text{theo}} = (3.03 \pm 0.10) \times 10^{-4} \text{ [PRD83('11)054007]}$$

• $\mathcal{B}(\Xi_b^- \to \Xi^- \gamma)_{\text{theo}} = (1.23 \pm 0.64) \times 10^{-5} \text{ [arXiv:2008.06624]}$

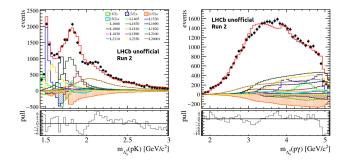
• Only 5% events accessible (HLT1)

- Use of $\Xi_b^-\to \Xi^- J/\psi$ for normalization removes dependence with ${\it f}_{\Xi_b}$
- Still blinded
- Expected to set an upper limit of $\mathcal{O}(10^{-4})$



$\Lambda^0_b ightarrow (\Lambda^* ightarrow p K^-) \gamma$ amplitude analysis

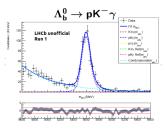
- Goal: Gain knowledge on the pK spectrum
 - Aid in the interpretation of the LFU measurement (R_{pK})
 - Pentaquarks: Set an upper limit on $\mathcal{B}(P_c o p\gamma)$
 - $P_c \rightarrow pJ/\psi$ was observed in $\Lambda_b^0 \rightarrow pK^-J/\psi$ [PRL115(2015)072001]
- 2D fit in m_{pK} and $m_{p\gamma}$
- Isobar model with parameters from PDG when available



Other radiative decays

Search for $\Lambda_b \rightarrow p \pi \gamma$

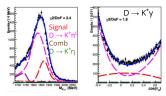
- $b \rightarrow d\gamma$ transition:
 - More suppressed
 - Larger CPV
- No previous limit
- *pK*γ and *K*πγ contaminations controlled by simultaneous fit



Radiative charm decays:

- $D \to K^*\gamma, \phi\gamma, \rho\gamma$
 - $c \rightarrow u\gamma$ transition
 - $\mathcal{B} \sim 10^{-4} 10^{-5}$
 - Cleaner NP probes (*A*_{CP}, *γ* pol.) [JHEP08(2017)091]
 - More background
 - Softer γ

•
$$\mathcal{B}(D o V^0 \pi^0) >> \mathcal{B}(V^0 \gamma)$$



~

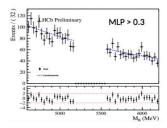
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Non-hadronic radiative decays

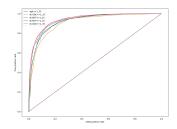
Search for $B_{\rm s} \rightarrow \mu \mu \gamma$ with Run 2 data

- Sensitive to C_7 , C_9 and C_{10}
- $\mathcal{B}_{SM} \sim 10^{-9} 10^{-10}$
- BaBar limit $\mathcal{B} < 10^{-7}$



Search for $\gamma\gamma$ final states (B_s and ALPs)

- Light ALPs not reachable for ATLAS/CMS
- $\mathcal{B}_{\mathsf{SM}} \sim 10^{-7}$
- Belle limit $\mathcal{B} < 3 imes 10^{-6}$
- Challenging topology



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Introduction

- Theoretical motivation
- LHCb Detector
- Experimental challenges

Analyses

- $B_s \rightarrow \phi \gamma$ time-dependent analysis
- $B \to K^* ee$ angular analysis
- $\Lambda_b^0 \to \Lambda \gamma$ (angular) analysis
- Ongoing analyses

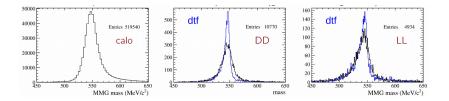
3 Conclusions and prospects



Prospects: Converted photons

When photons convert before the magnet ($\gamma
ightarrow e^+e^-$):

- e^+e^- tracks can be reconstructed \implies resolution 3 times better
 - ullet Allow to access $|V_{\rm td}/V_{\rm ts}|$ suppressed modes like ${\it B_s} \to {\it K^*\gamma}$
- Its rate is 20 times lower than for calo photons



For Run 3:

- Converted photon analyses feasible with more luminosity
- Improvements in electron tracking for Upgrade

Conclusions

- The **photon polarization** is being measured at LHCb using several channels and different observables
- Important constraints to the Wilson coefficient $C_7^{(\prime)}$
- Nice competition ahead with Belle II

Preparations for Run 3:

- Developing fast downstream tracking algorithms fitting Hlt1 time-budget
 - Can increase statistics in b-baryon analysis up to 20 times
- Including BDT methods to HIt2 lines:
 - Better sgl/bkg separation
 - Add downstream to channels with huge bkg

Stay TUNED FOR something AWESOME



Thanks for your attention





Ratio of branching ratios

The branching ratio can be extracted from the signal yield:

$$N = 2 \times \mathcal{L} \times \sigma_{b\overline{b}} \times f_{\Xi_{b}} \times \mathcal{B}_{\Xi_{b}^{-} \to \Xi^{-} \gamma} \times \mathcal{B}_{\text{sub-decays}} \times \epsilon_{sel}$$

The \mathcal{B} is extracted as a ratio of \mathcal{B} using dimuonic channels

• Same hadronic decay chain

$$\frac{\mathcal{N}(\Xi_{b}^{-} \to \Xi^{-} \gamma)}{\mathcal{N}(\Xi_{b}^{-} \to \Xi^{-} J/\psi)} = \frac{\mathcal{B}(\Xi_{b}^{-} \to \Xi^{-} \gamma)}{\mathcal{B}(\Xi_{b}^{-} \to \Xi^{-} J/\psi)} \times \frac{1}{\mathcal{B}(J/\psi \to \mu^{+} \mu^{-})} \times \frac{\epsilon_{sel}(\Xi_{b}^{-} \to \Xi^{-} \gamma)}{\epsilon_{sel}(\Xi_{b}^{-} \to \Xi^{-} J/\psi)}$$

where:

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• $\epsilon_{sel} = \epsilon_{acc} \times \epsilon_{reco+strip} \times \epsilon_{trigger} \times \epsilon_{Presel} \times \epsilon_{PID} \times \epsilon_{BDT} \times \epsilon_{IsPhoton}$



Why $\Xi_b^- \to \Xi^- J/\psi$ as normalization channel

 $\begin{aligned} \mathcal{B}(\Xi_b^- \to \Xi^- J/\psi) &= (5.0 \pm 2.4) \times 10^{-4} \\ \mathcal{B}(\Lambda_b^0 \to \Lambda J/\psi) &= (3.29 \pm 1.11) \times 10^{-4} \\ \mathcal{B}(\Lambda_b^0 \to \rho K^- J/\psi) &= (3.2 \pm 0.6) \times 10^{-4} \\ \frac{f_{\Xi_b}}{f_{\Lambda_b}} &= (8.2 \pm 2.6) \times 10^{-4} \end{aligned}$

$$\begin{split} \mathcal{B}(\Xi_b^- \to \Xi^- \gamma) &\propto \mathcal{B}(\Xi_b^- \to \Xi^- J/\psi) & \implies \sigma = 48\% \\ \mathcal{B}(\Xi_b^- \to \Xi^- \gamma) &\propto \frac{1}{f_{\Xi_b}/f_{\Lambda_b}} \mathcal{B}(\Lambda_b^0 \to \Lambda \gamma) & \implies \sigma = 46\% \\ \mathcal{B}(\Xi_b^- \to \Xi^- \gamma) &\propto \frac{1}{f_{\Xi_b}/f_{\Lambda_b}} \mathcal{B}(\Lambda_b^0 \to pK^- J/\psi) & \implies \sigma = 37\% \end{split}$$

• When $\mathcal{B}(\Lambda_b^0 \to \Lambda J/\psi)$ is measured by LHCb $\sigma = 37\%$

• When $\mathcal{B}(\Xi_b^- \to \Xi^- J/\psi)$ is measured by LHCb, no f_{H_b} dependency and $\sigma = 20\%$



Expected yields: $\Xi_b^- \to \Xi^- \gamma$

$$\textit{N} = 2 \times \mathcal{L} \times \sigma_{b\overline{b}} \times \textit{f}_{\Xi_b} \times \mathcal{B}_{\Xi_b^- \to \Xi^- \gamma} \times \mathcal{B}_{\text{sub-decays}} \times \epsilon_{\textit{sel}}$$

| Variable | Value | |
|--|-----------------------------|-----------------------------|
| $\int \mathcal{L} [fb^{-1}]$ | 2.19 | |
| $\sigma_{b\overline{b}} \; [\mu b]$ | ~ 600 | |
| f_{Ξ_b} [%] | 0.021 ± 0.007 | |
| $\mathcal{B}(\Xi_b^- \to \Xi^- \gamma)$ | $(1.1\pm0.3)	imes10^{-5}$ (| $(3.03\pm0.10)	imes10^{-4}$ |
| $\mathcal{B}(\Xi^- \to \Lambda^0 \pi^-)$ [%] | 99.887 ± 0.035 | |
| ${\cal B}(\Lambda^0 	o p\pi)$ [%] | 63.9 ± 0.5 | |
| ϵ_{sel} | $(4.3\pm 0.2)	imes 10^{-6}$ | |
| Expected yield | 0.8 ± 0.4 | 22 ± 7 |

Assumming:

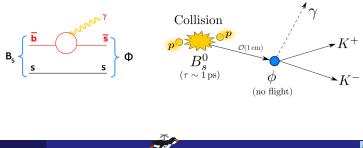
- $\mathcal{B}(\Xi_b^- \to \Xi^- \gamma) = \frac{3}{2} \mathcal{B}(\Lambda_b^0 \to \Lambda \gamma)$ (as for dimuonic channels)
- Theoretical prediction



Introduction: Meson decays

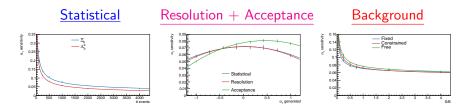
Radiative b-meson decays:

- Can oscillate
- SV can be determined from resonances decays
- Quite studied (lot of experience)
- Large production
- Accessible at B-factories



Sensitivity to photon polarization

Sensitivity to the photon polarization in the $\Lambda_b \rightarrow \Lambda \gamma$ decay using angular distribution [EPJC(2019)79:634]:



• Statistical uncertainty: Goes as $1/\sqrt{N}$ with number of events

decays at LHCb

- Resolution: Effect neglible
- Acceptance: Asymmetric in α_{γ}
- Background: Important dilution

Simultaneous Mass fit: Validation

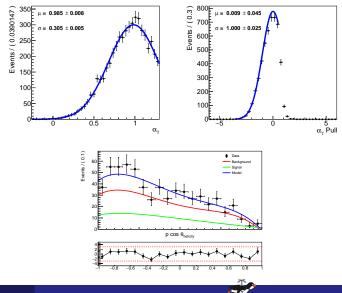
Mass fit validated throught MC Study:

- 1000 pseudo-experiments
- Testing $\mathcal{B}(\Xi_b^- \to \Xi^- \gamma)$ hypotheses in range $(10^{-5}, 10^{-3})$
- Measurement significance using Wilk's theorem

| $\mathcal{B}(\Xi_b^- \to \Xi^- \gamma)$ | Evidence Prob [%] ($\sigma \geq$ 3) | Observation Prob [%] ($\sigma \ge 5$) |
|---|--------------------------------------|---|
| $1	imes 10^{-3}$ | 100.0 | 100.0 |
| $5	imes 10^{-4}$ | 100.0 | 98.0 |
| $3.3	imes10^{-4}$ | 96.7 | 57.1 |
| $1	imes 10^{-4}$ | 26.2 | 0.8 |
| $5	imes 10^{-5}$ | 5.5 | 0.0 |
| $1.1	imes10^{-5}$ | 0.9 | 0.0 |



MC-Toys result (S=197, B=412): α_{γ}



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Radiative decays at

LHCD

Amplitud Model

Helicity formalism + isobar lineshape in m_{pK}

Amplitude for a defined set of helicities in $\Lambda_b \to \Lambda^*(\to pK)\gamma$

To obtain full decay rate

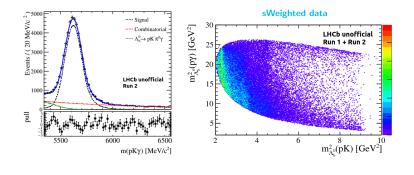
Coherently sum possible Λ^* helicities and Λ^* resonances Incoherently sum possible proton, photon, and Λ_b helicities

To check for pentaquarks in $m_{p\gamma}$

Need for second decay chain amplitude Non-trivial relation between the decay planes of $\Lambda_b \to \Lambda^*(\to pK)\gamma$ and $\Lambda_b \to P_c(\to p\gamma)K$

Mass fit and Dalitz plot

- Signal: Bifurcated Crystal Ball
- Background:
 - Combinatorial: Exponential
 - Partially reconstructed: Argus convoluted with signal shape
- Tails parameters fixed in MC



Introduction: Experimental status

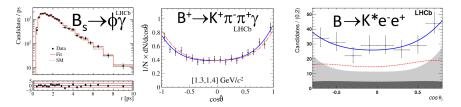
Photon polarization has only been measured using radiative *b*-meson decays:

Proper time distribution:

• $B - \overline{B}$ interference: $B_s \rightarrow \phi \gamma$, [PRL118('17)109901]

Angular distribution:

- $B^+ \to K^- \pi^+ \pi^+ \gamma$, [PRL 112('14)161801]
- $B \rightarrow K^* e^+ e^-$ at low q^2 , [JHEP04('15)064]





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Angular distribution: $\Xi_b^- o \Xi^- \gamma$

Angular distribution for unpolarized Ξ_b^- decay (integrating on azimuthal angles):

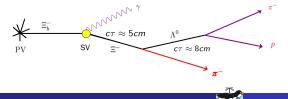
$$\Gamma_{\Xi_{b}}(\theta_{\Lambda},\theta_{p}) = \frac{1}{4} \left(1 - \alpha_{\gamma} \alpha_{\Xi} \cos \theta_{\Lambda} + \alpha_{\Lambda} \cos \theta_{p} \left(\alpha_{\Xi} - \alpha_{\gamma} \cos \theta_{\Lambda} \right) \right)$$

The values of the decay parameters are:

- α_Ξ = −0.401 ± 0.010 [PDG 2020]
- α_Λ = 0.732 ± 0.010 [PDG 2020]

Advantages of Ξ_b^- over Λ_b

- Extra decay ⇒ richer angular distribution
- Charged particle: Ξ⁻



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