



皮影 Shadow play

Source: <http://www.cnhubei.com/ztmjys-pyts>

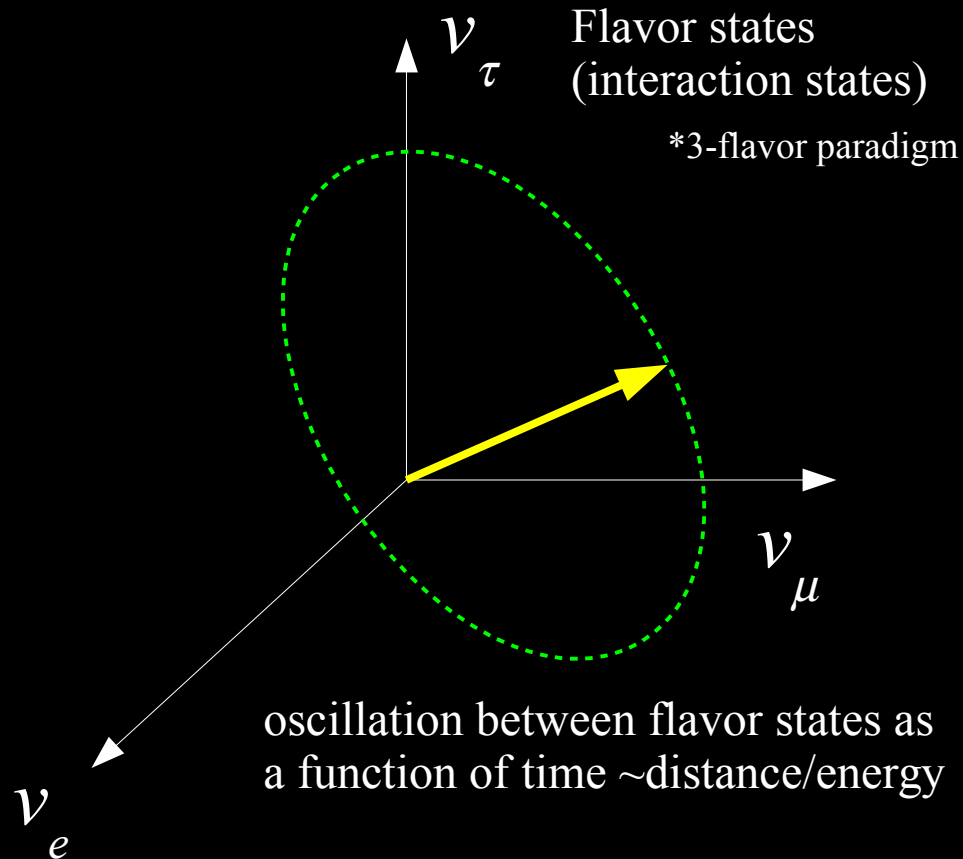
Neutrino Shadow Play

– Kinematic determination of nuclear effects at MINERvA

Xianguo LU/ 卢显国 University of Oxford
on behalf of MINERvA Collaboration
Joint Experimental-Theoretical Physics Seminar
FNAL, 2 March 2018

Neutrino

– Oscillation



The Nobel Prize in Physics 2015



Photo: A. Mahmoud
Takaaki Kajita
Prize share: 1/2



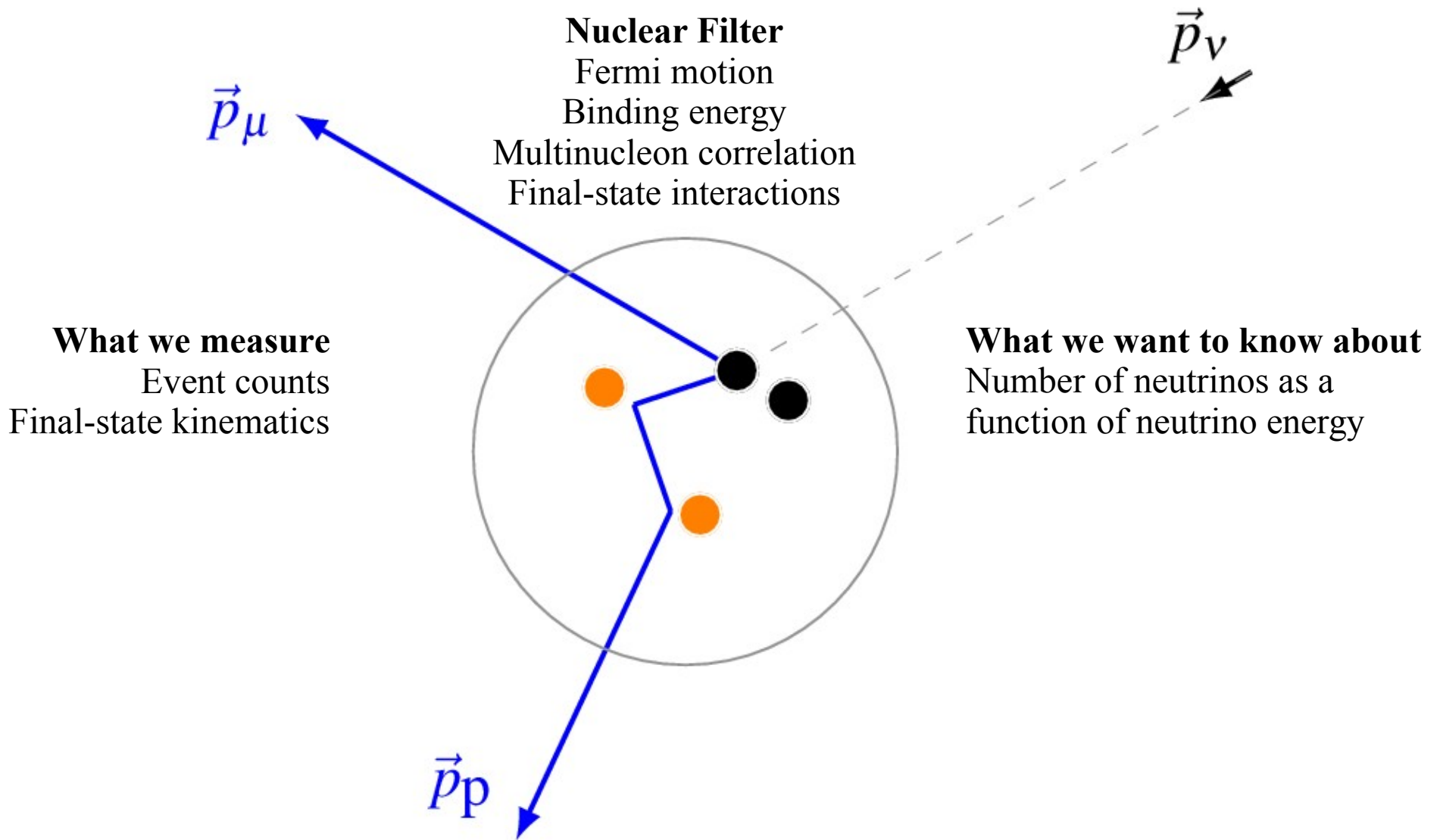
Photo: A. Mahmoud
Arthur B. McDonald
Prize share: 1/2

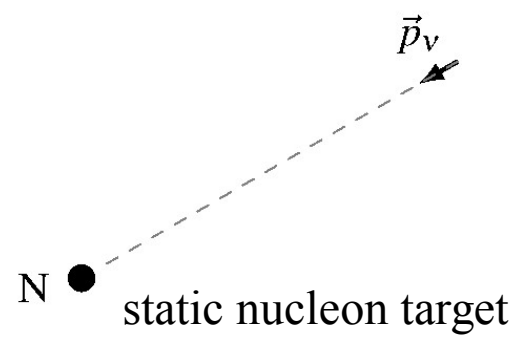
The Nobel Prize in Physics 2015 was awarded jointly to Takaaki Kajita and Arthur B. McDonald *"for the discovery of neutrino oscillations, which shows that neutrinos have mass"*

https://www.nobelprize.org/nobel_prizes/physics/laureates/2015/

Neutrino mass:
shift between interaction and propagation states

The big picture of neutrino detection in oscillation experiment





Neutrino energy in GeV regiem

Quasi-elastic scattering (QE)

$$\nu n \rightarrow \ell^- p$$

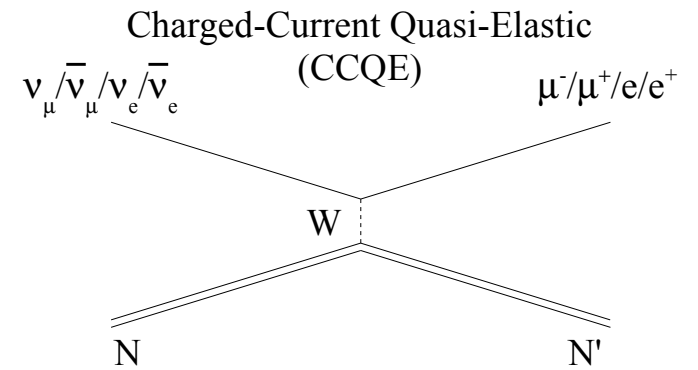
charged current (CC) $\nu \rightarrow \ell'$

$\vec{p}_{\ell'}$

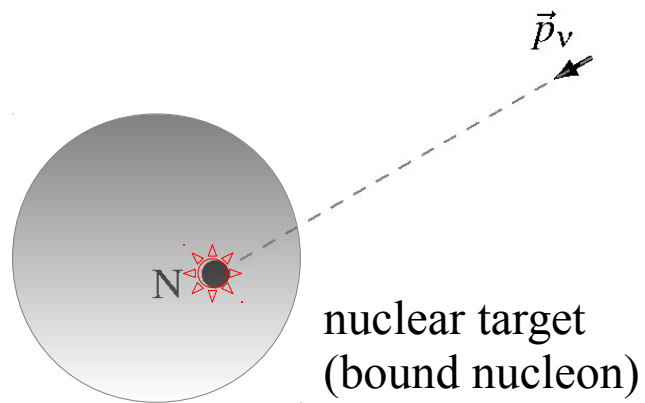
static nucleon target

$\vec{p}_{N'}$

quasi-elastic (QE) $N \rightarrow N'$

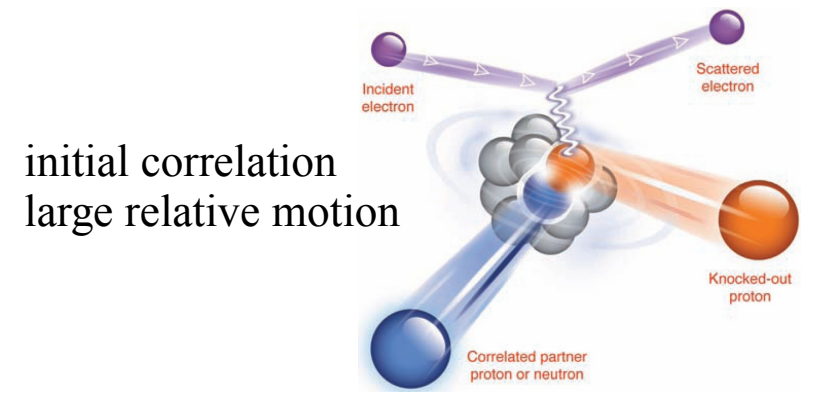


Fermi motion (FM) biases E_ν reconstruction

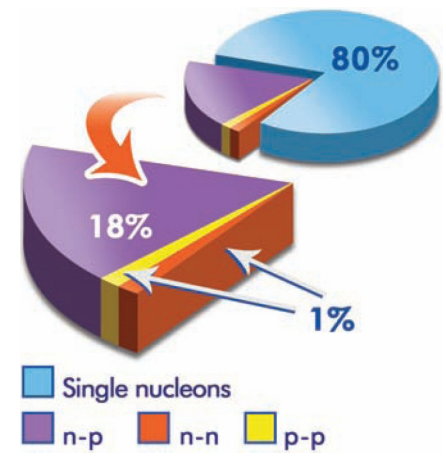


Fermi motion (FM) biases E_v reconstruction

Multinucleon correlations:



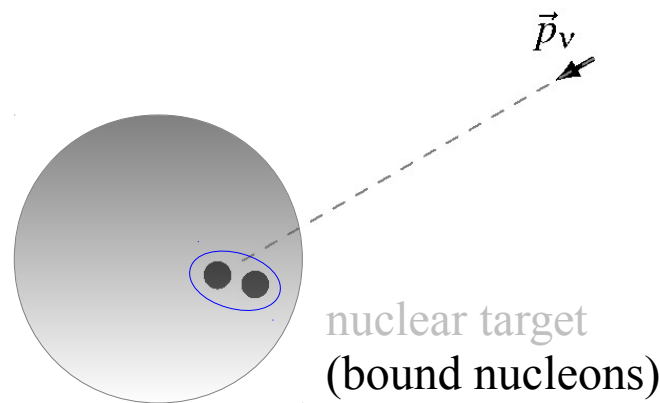
Science 320 (2008) 1476-1478



Fermi motion (FM) biases E_ν reconstruction

Multinucleon correlations:

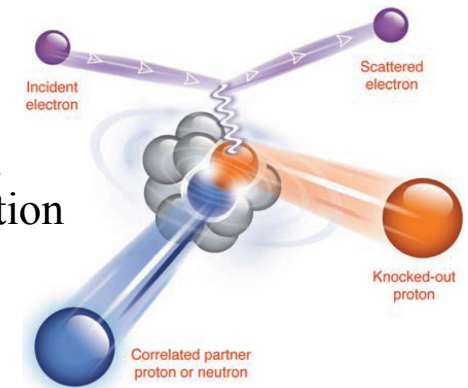
cross section unknown, strong bias to *all* final-state kinematics



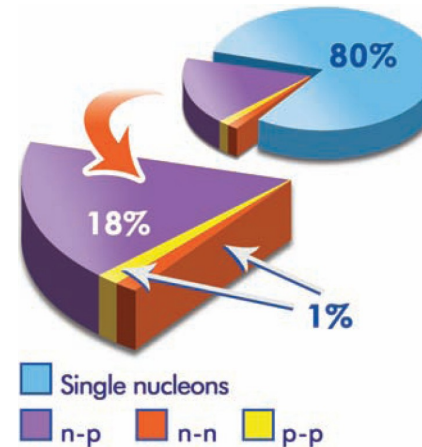
- Impulse approximation: independent particles
- In **p**article-**h**ole excitation:
 - RPA (random phase approximation): sum of 1p1h excitation (over all pairs) \sim ground state correlations (long range)
 - npnh ($n \geq 2$): sub-leading terms in ph expansion \sim multinucleon correlations (short range)

Xianguo Lu, Oxford

initial correlation
large relative motion



Science 320 (2008) 1476-1478



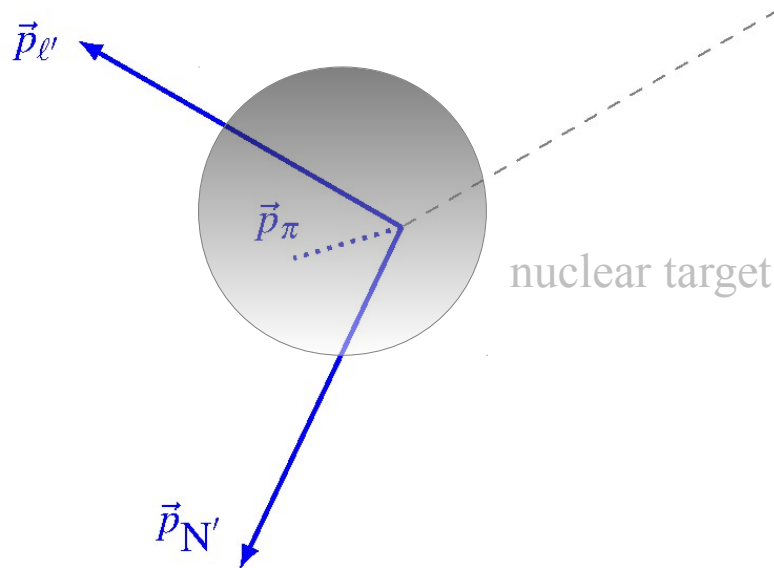
Fermi motion (FM) biases E_ν reconstruction

Multinucleon correlations:

cross section unknown, strong bias to *all* final-state kinematics

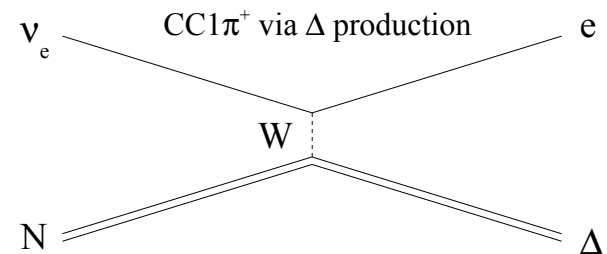
QE-like: π absorbed in nucleus \leftarrow final-state interaction (FSI)

charged current (CC) $\nu \rightarrow l'$



Resonance production (RES)

$$\nu p \rightarrow l^- \Delta^{++} \rightarrow l^- p \pi^+$$



QE-like $N \rightarrow N'$

including resonance production (RES) $\Delta \rightarrow N'\pi$ followed by π absorption

Fermi motion (FM) biases E_ν reconstruction

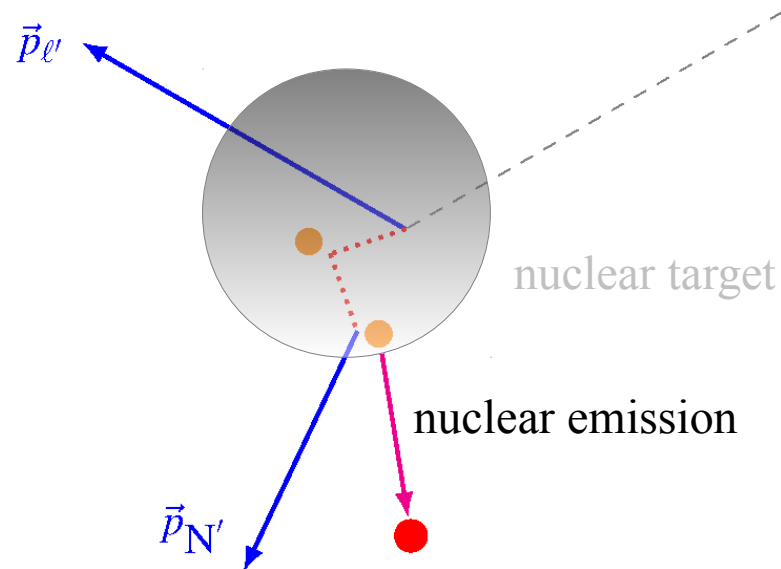
Multinucleon correlations:

cross section unknown, strong bias to *all* final-state kinematics

QE-like: π absorbed in nucleus \leftarrow final-state interaction (FSI)

FSI \rightarrow energy-momentum transferred in nucleus, possible nuclear emission

charged current (CC) $\nu \rightarrow l'$



QE-like $N \rightarrow N'$

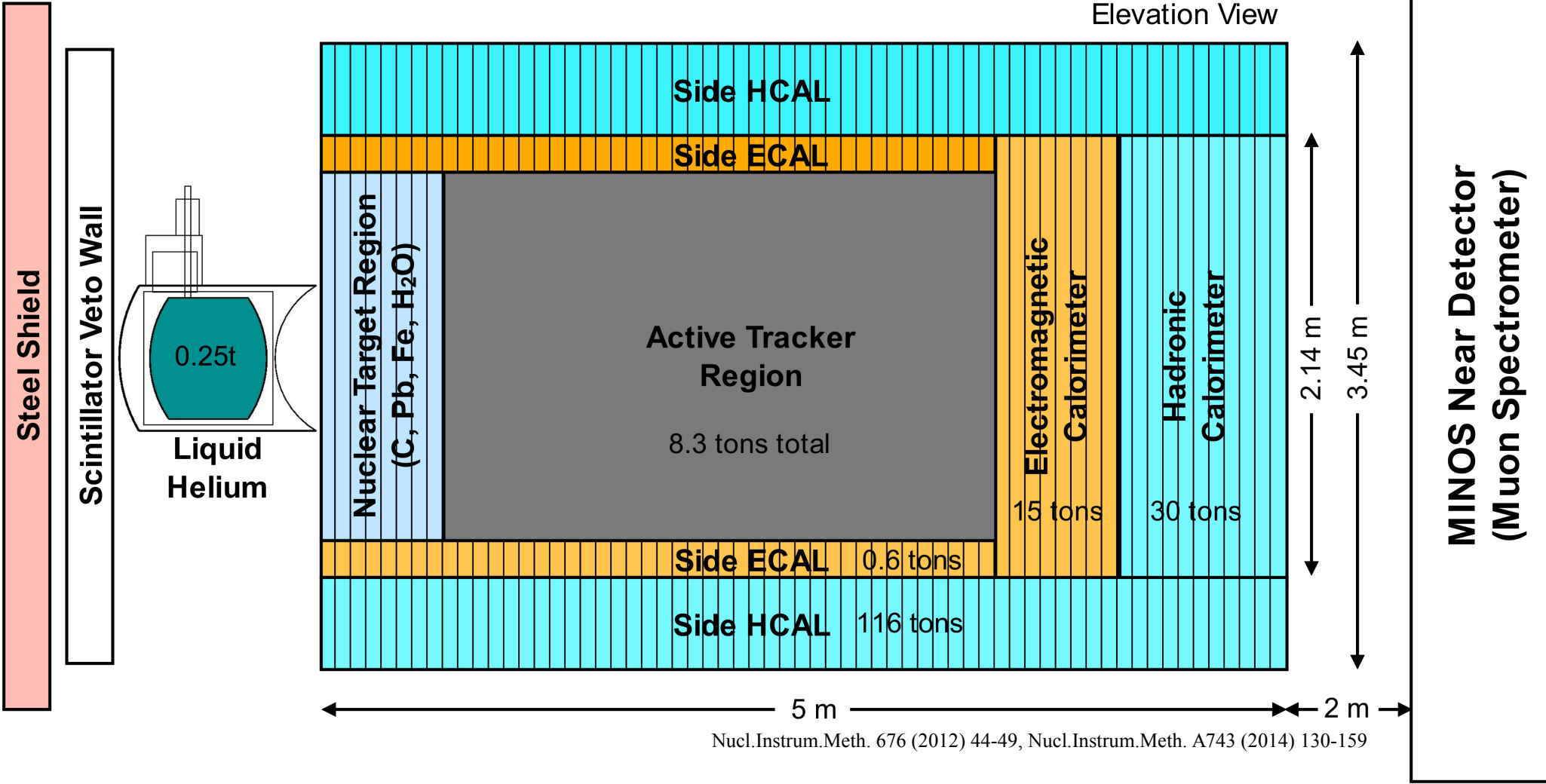
including resonance production (RES) $\Delta \rightarrow N'\pi$ followed by π absorption

MINERvA



MINERvA

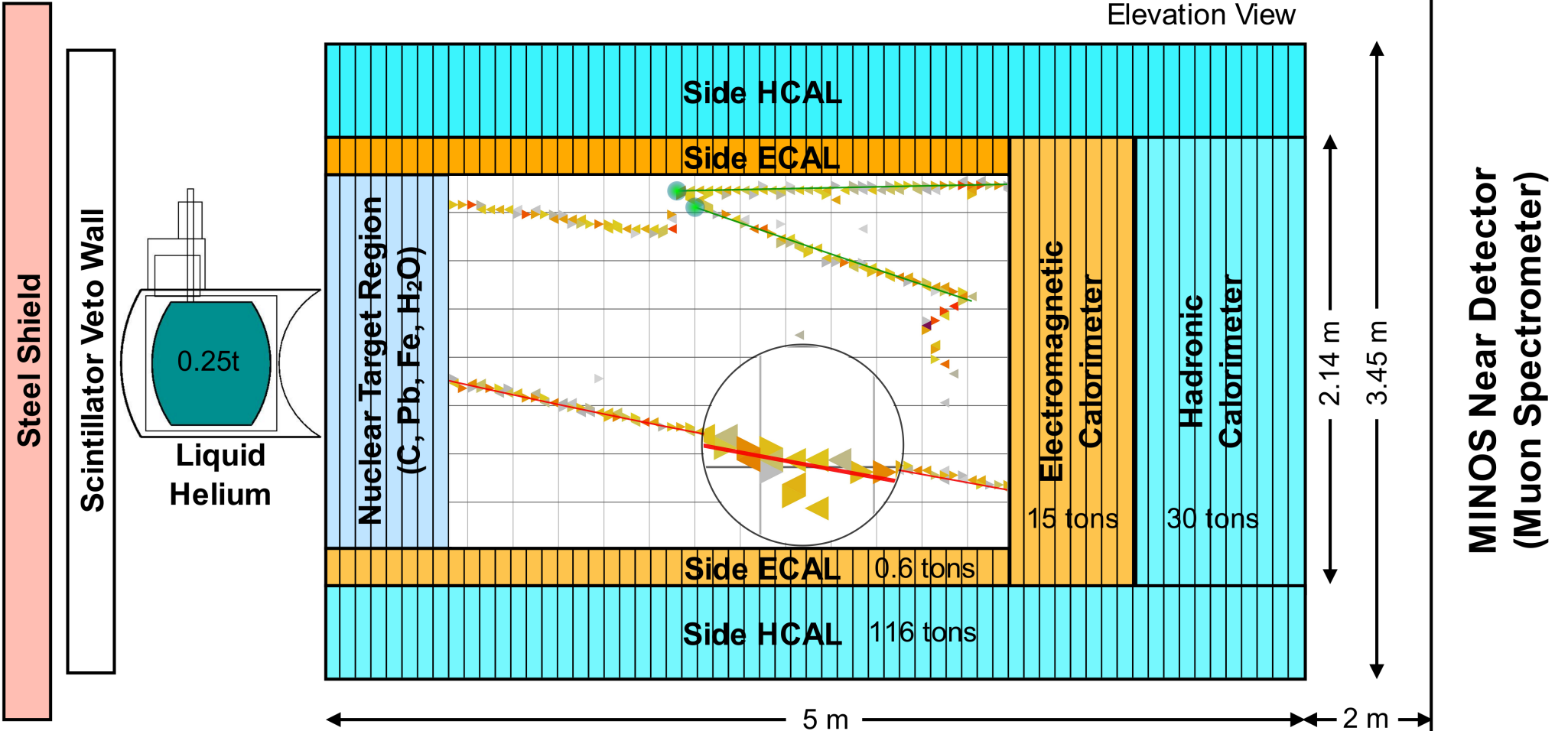
Thanks to MINOS!



Scintillator tracker

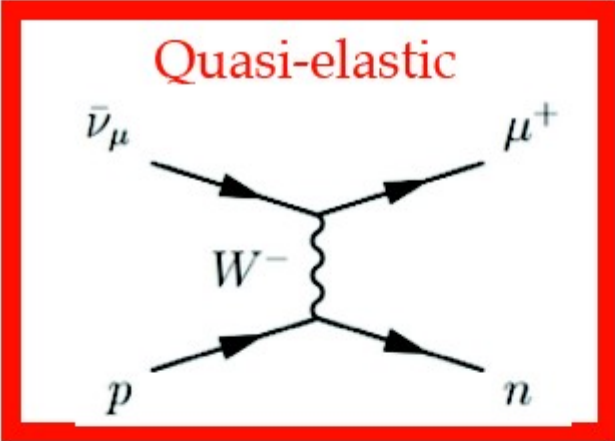
MINERvA

Thanks to MINOS!

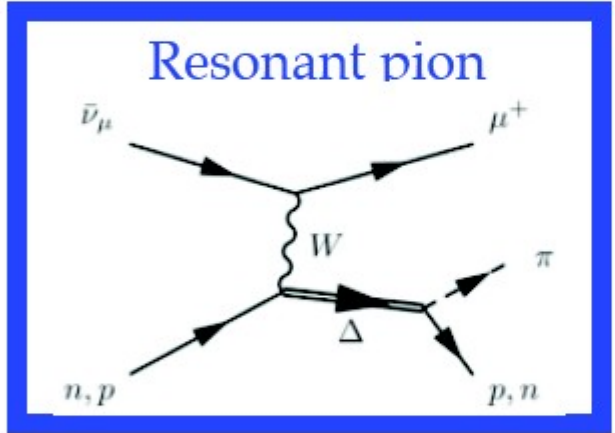


Nucl.Instrum.Meth. 676 (2012) 44-49, Nucl.Instrum.Meth. A743 (2014) 130-159

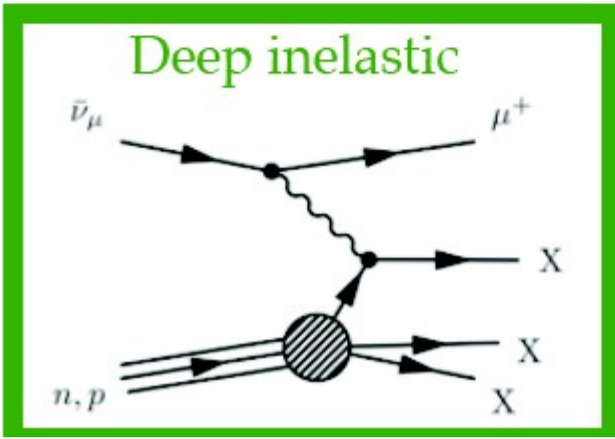
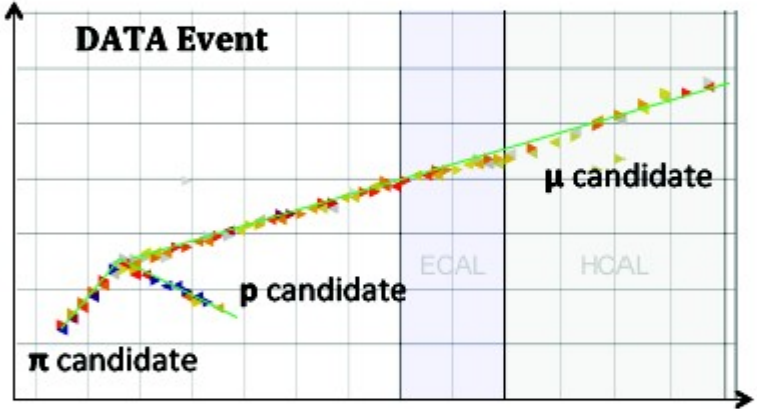
Scintillator tracker:
 Hydrocarbon (CH) target
 Homogeneous non-magnetized active tracker



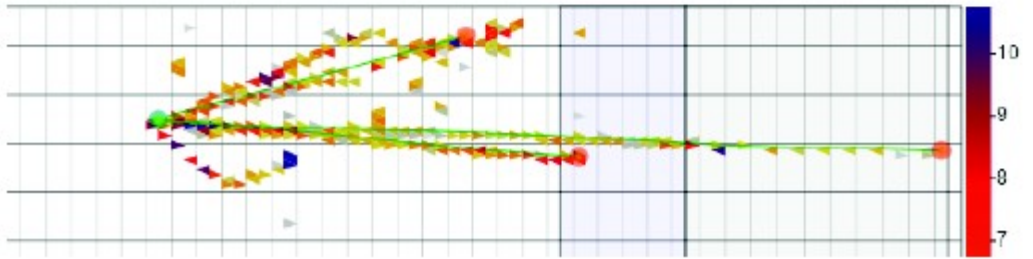
QE

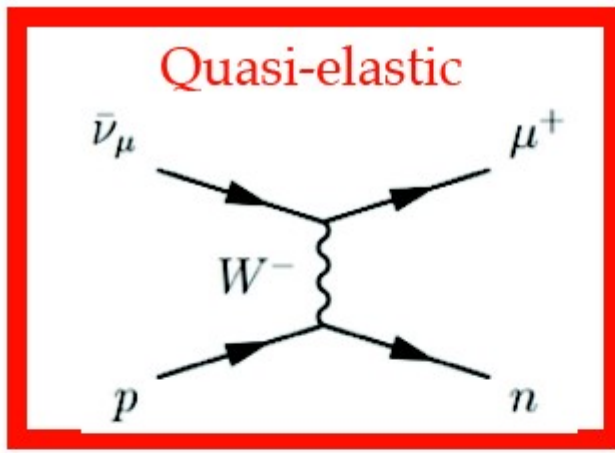


RES

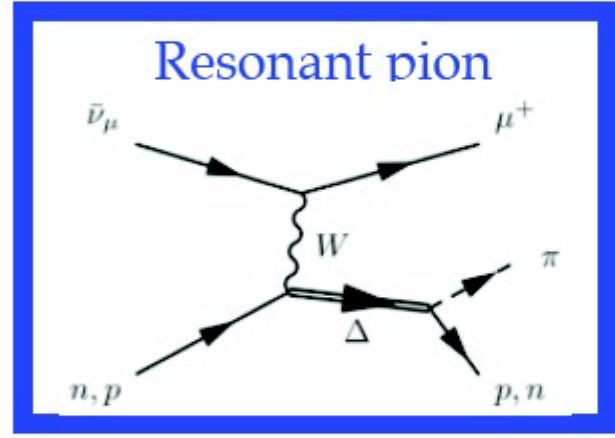


DIS

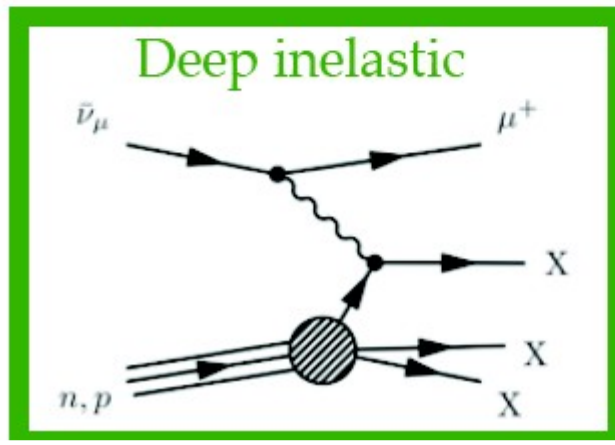




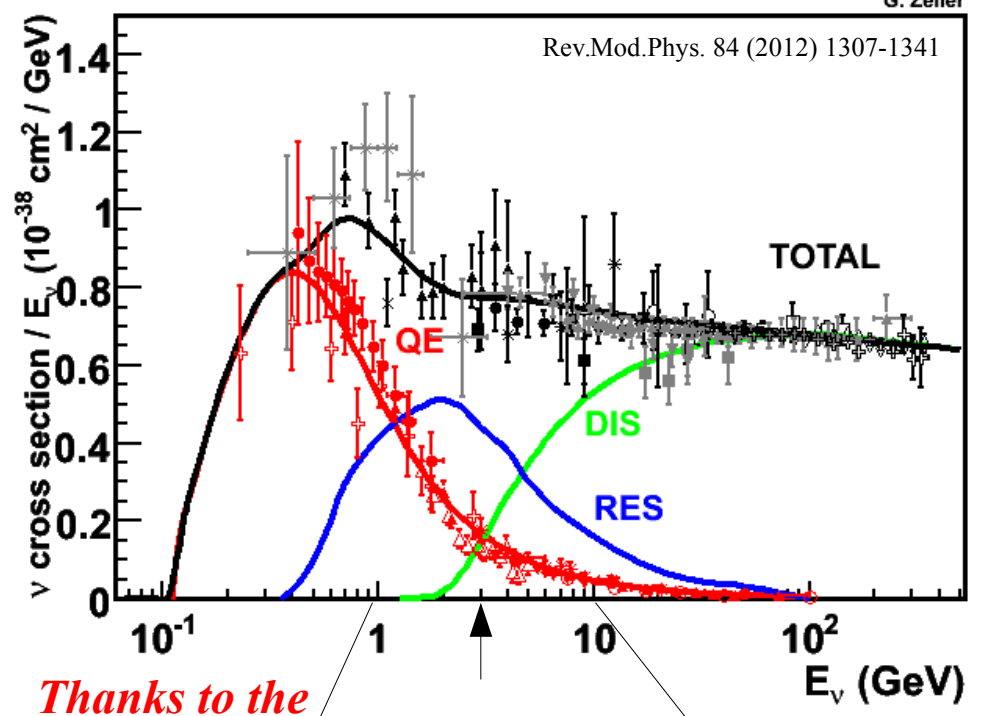
QE



RES

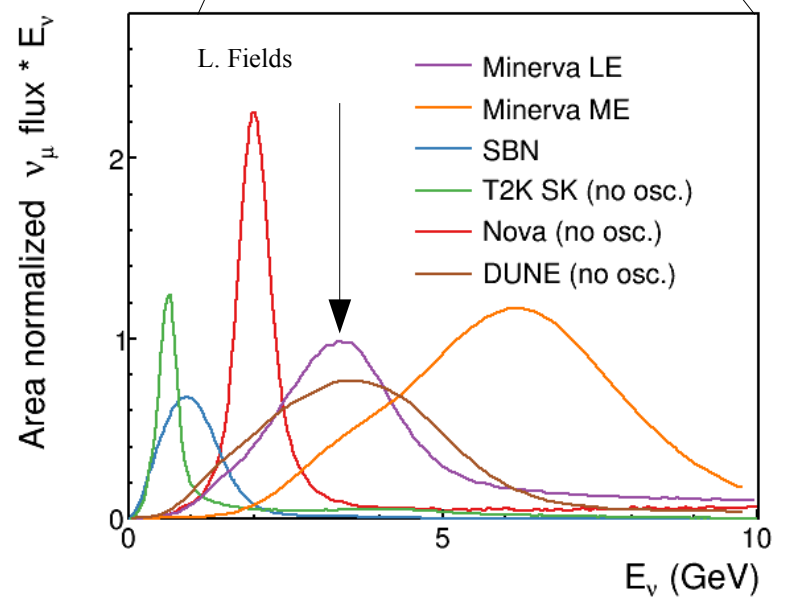


DIS



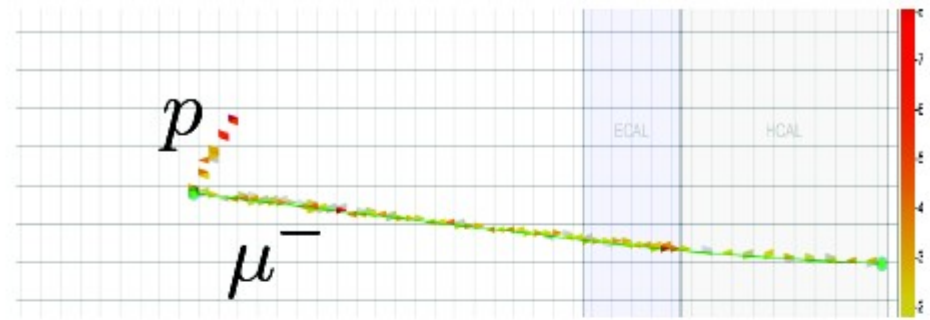
Thanks to the accelerator division!

NuMI low energy beam $\langle E_\nu \rangle \sim 3 \text{ GeV}$



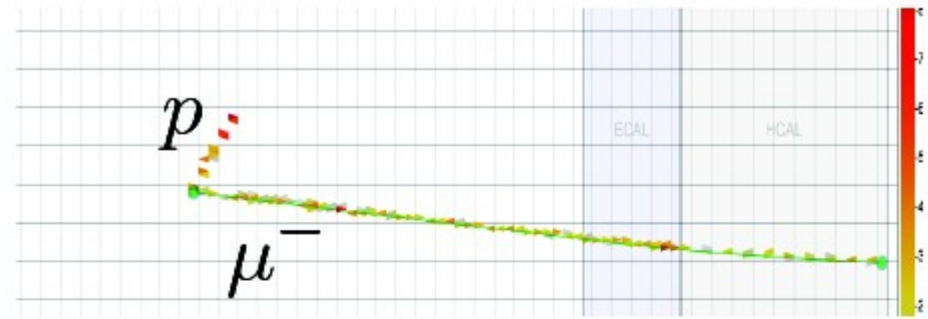
Today's topic:

μ - p mesonless production



Today's topic:

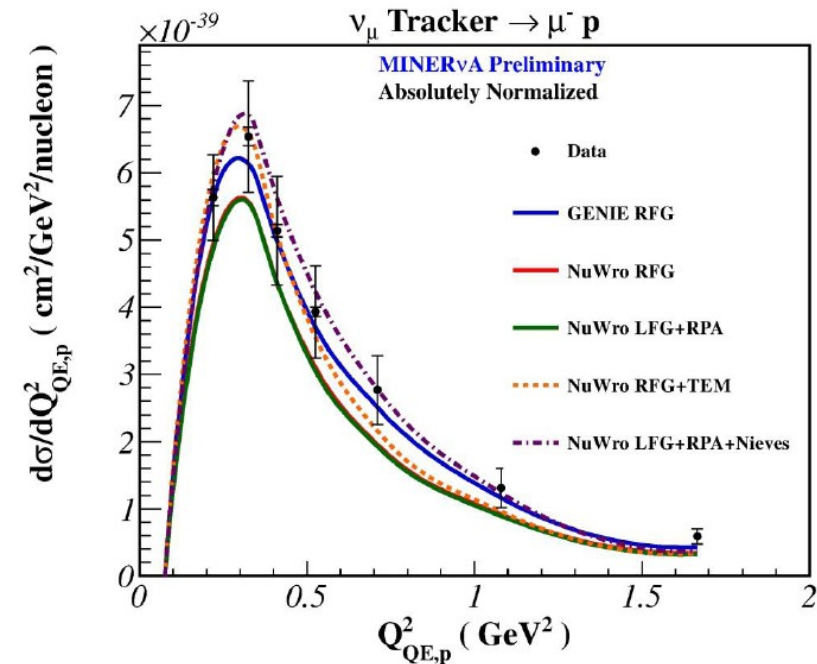
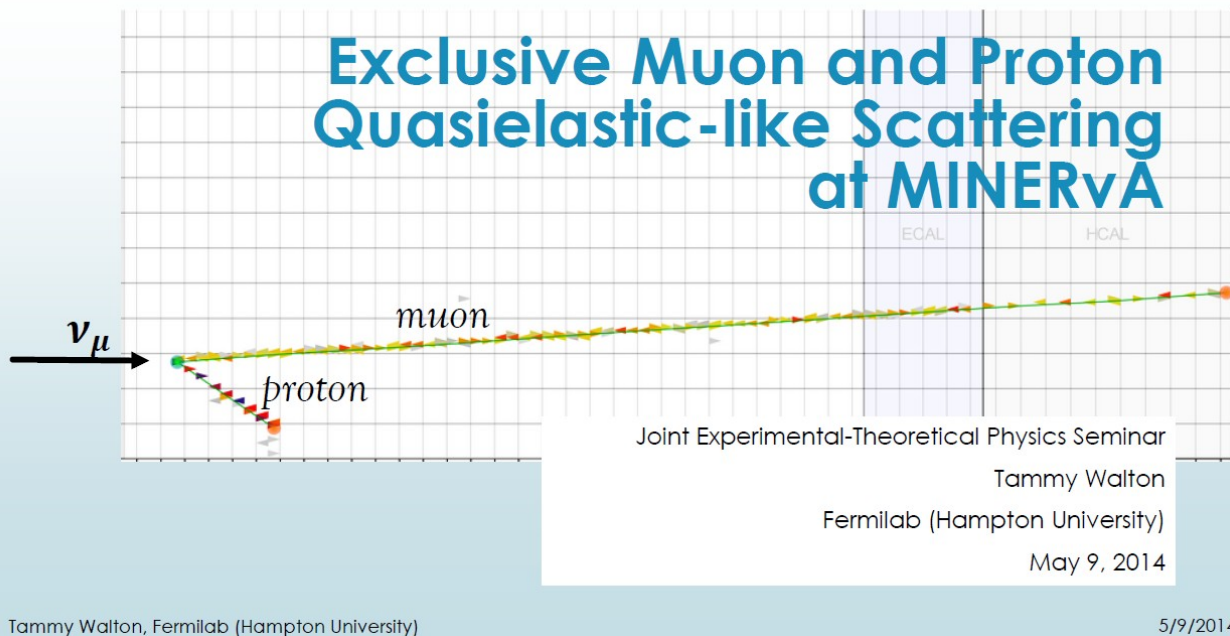
μ -p mesonless production



First presented by Tammy Walton in 2014 [Phys. Rev. D 91, 071301(R)]

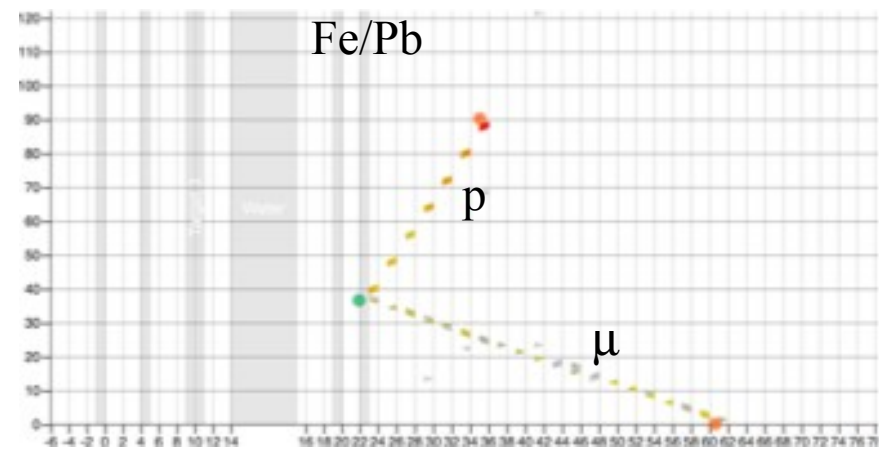
- Hydrocarbon target
- NuMI low energy (LE) neutrino beam
- Foundation of further MINERvA μ p analyses

Exclusive Muon and Proton Quasielastic-like Scattering at MINERvA



Today's topic:

μ -p mesonless production



Extension to nuclear targets presented by Minerba Betancourt in 2016 [Phys. Rev. Lett. 119, 082001]

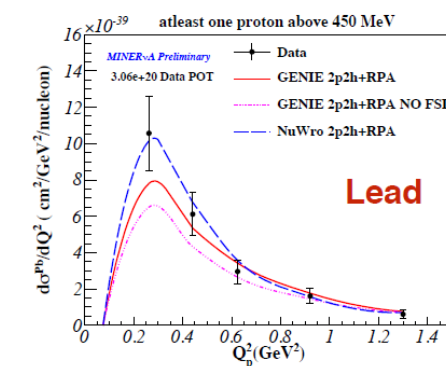
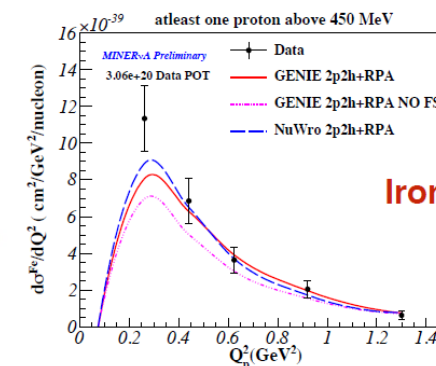
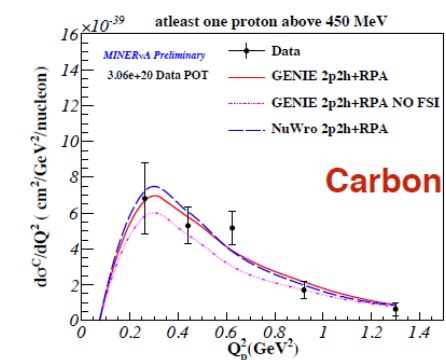
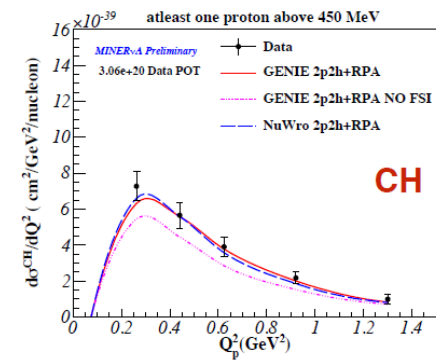
- Extended previous framework
- Challenging analysis

Managed by Fermi Research Alliance, LLC for the U.S. Department of Energy Office of Science

Nuclear Dependence of Quasi-Elastic Scattering at MINERvA

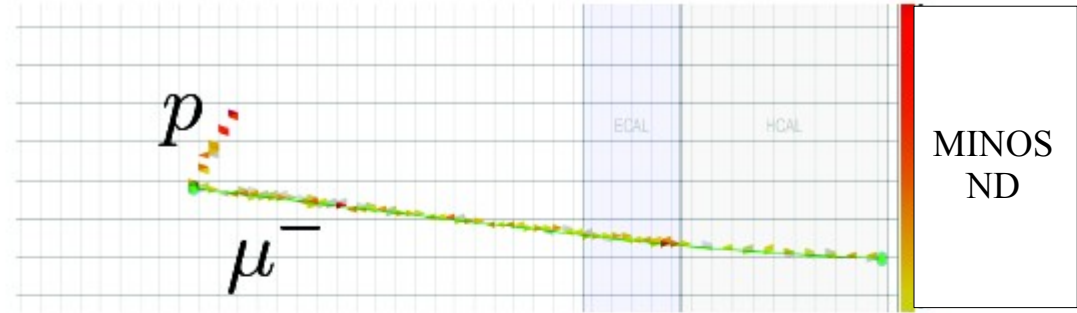
- Introduction
- The MINERvA experiment
- Event Selection
- Analysis
- Results

Minerba Betancourt, Fermilab
 Joint Experimental-Theoretical Seminar
 07 October 2016



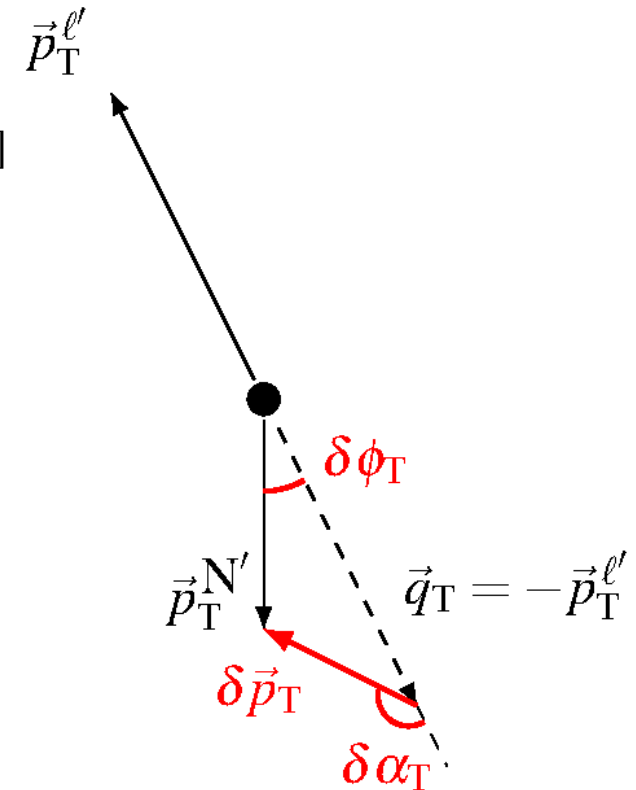
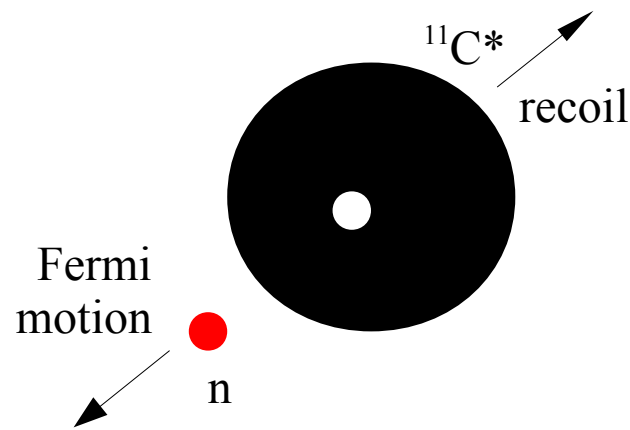
Today's topic:

μ -p mesonless production



This analysis

- Subsample of 2014 analysis
 - Muon matched to MINOS
 - **Proton kinematics** measurement significantly better
- New observables:
 - **Transverse kinematics imbalances**
[XL, L. Pickering, S. Dolan *et al.*, Phys.Rev. C94 (2016) no.1, 015503]
 - **Initial neutron momentum**
[A. Furmanski, J. Sobczyk, Phys.Rev. C95 (2017) no.6, 065501]



A brief history of Shadow Art



https://en.wikipedia.org/wiki/Cave_painting

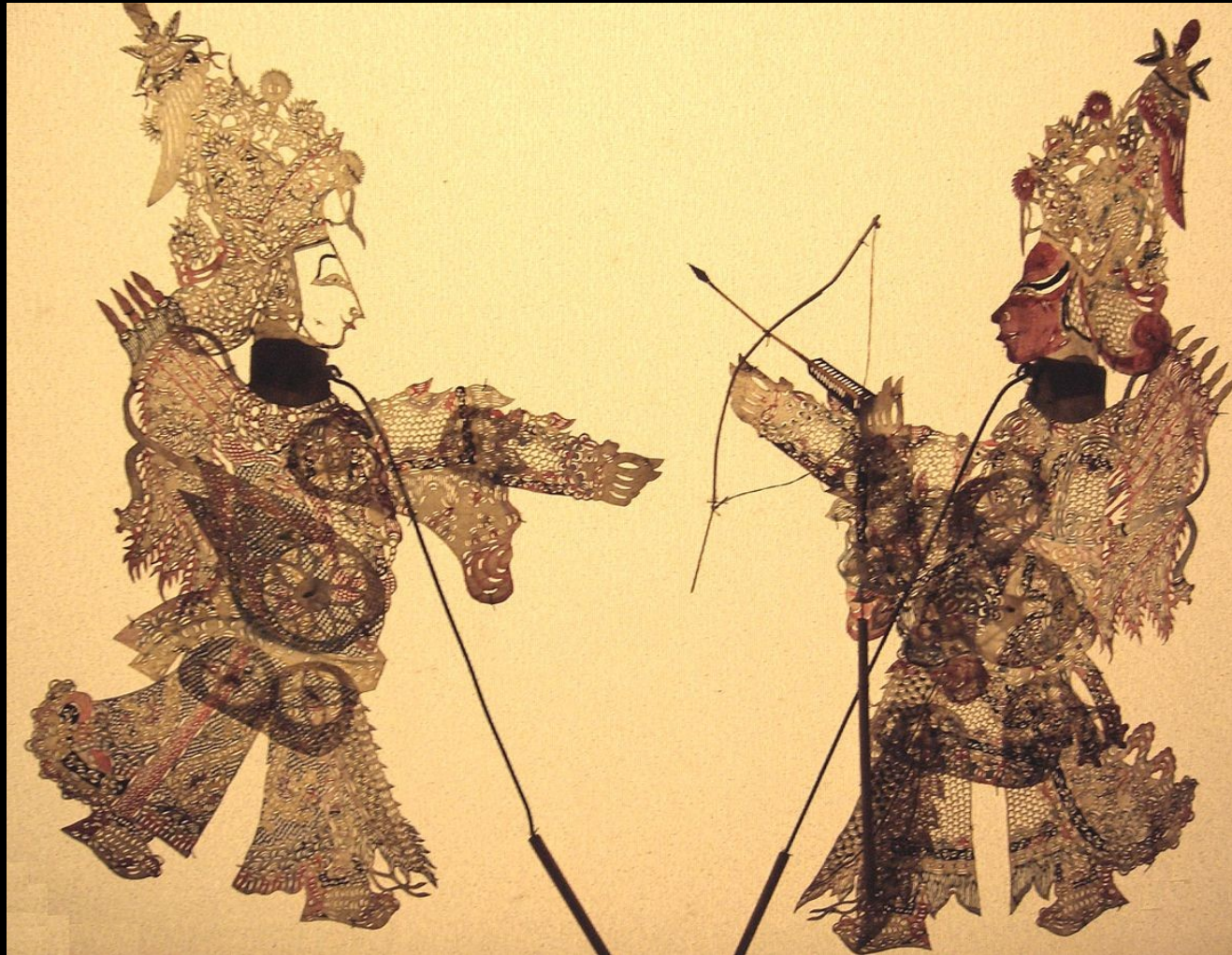
Cave of Pettakere, Bantimurung district (kecamatan), South Sulawesi,
Indonesia. Hand stencils estimated between **35,000-40,000 BP**

A brief history of Shadow Art



https://en.wikipedia.org/wiki/Cave_painting

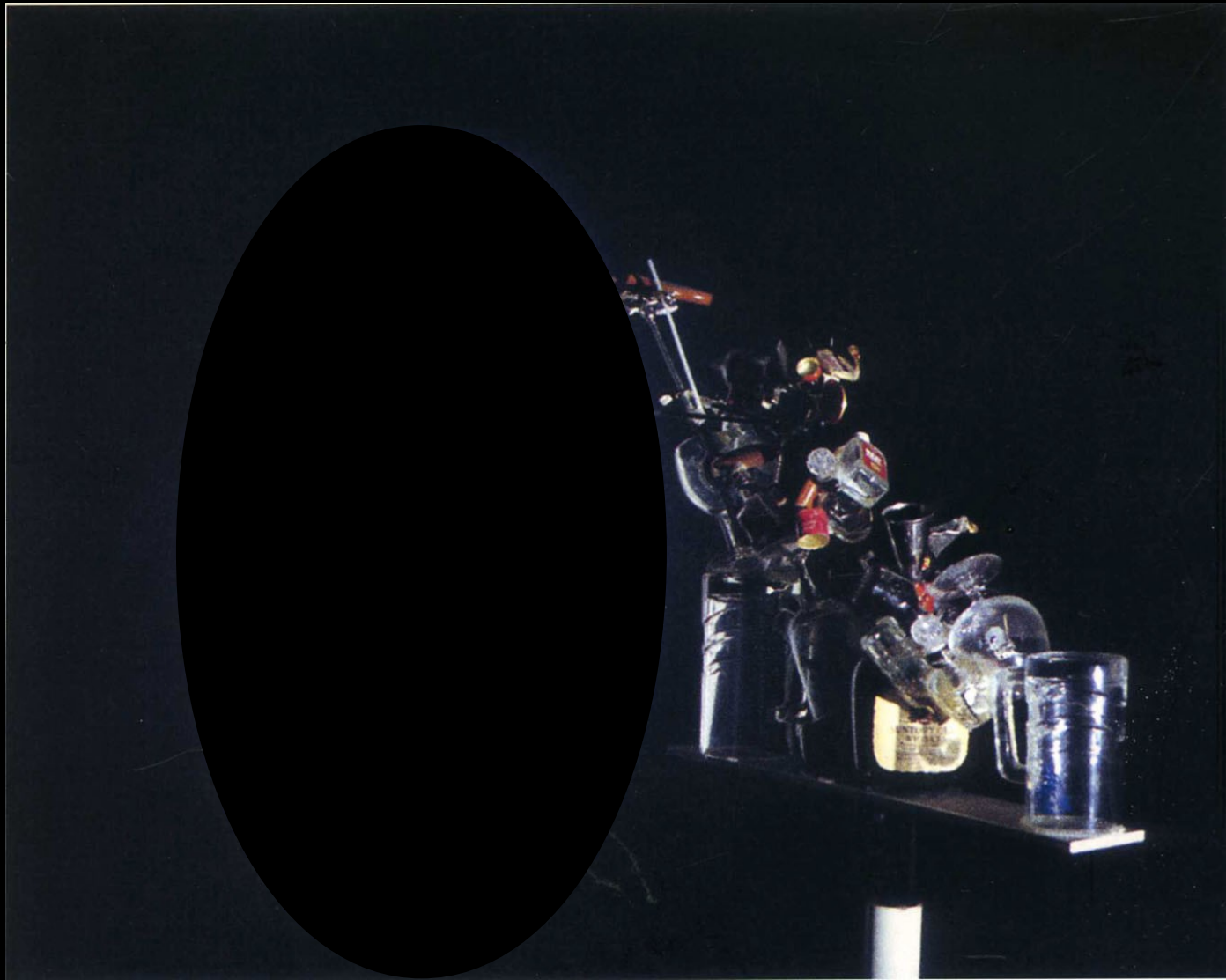
A brief history of Shadow Art



<https://zh.wikipedia.org/wiki/%E7%9A%AE%E5%BD%B1%E6%88%B2>

Traditional Chinese “movie”

A brief history of Shadow Art



<http://www.spoon-tamago.com/2015/08/03/illusionistic-shadow-art-by-shigeo-fukuda/>

Japanese modern art

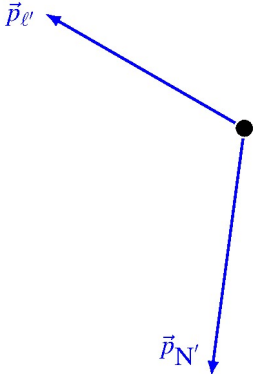
A brief history of Shadow Art



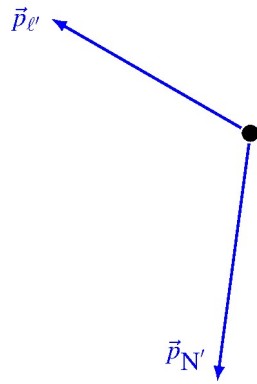
<http://www.spoon-tamago.com/2015/08/03/illusionistic-shadow-art-by-shigeo-fukuda/>

Japanese modern art

Transverse kinematic imbalances – *a neutrino shadow play*



Transverse kinematic imbalances – a neutrino shadow play



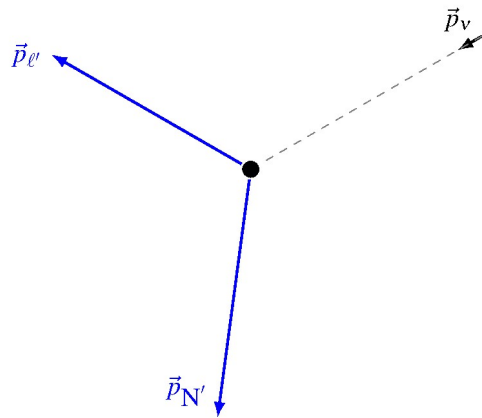
Source: <http://zhejiangpiying.sokutu.com/tupian.html>



To make *Neutrino Shadow Play*, we need

- ✓ beam of light
- ✓ screen

Transverse kinematic imbalances – a neutrino shadow play



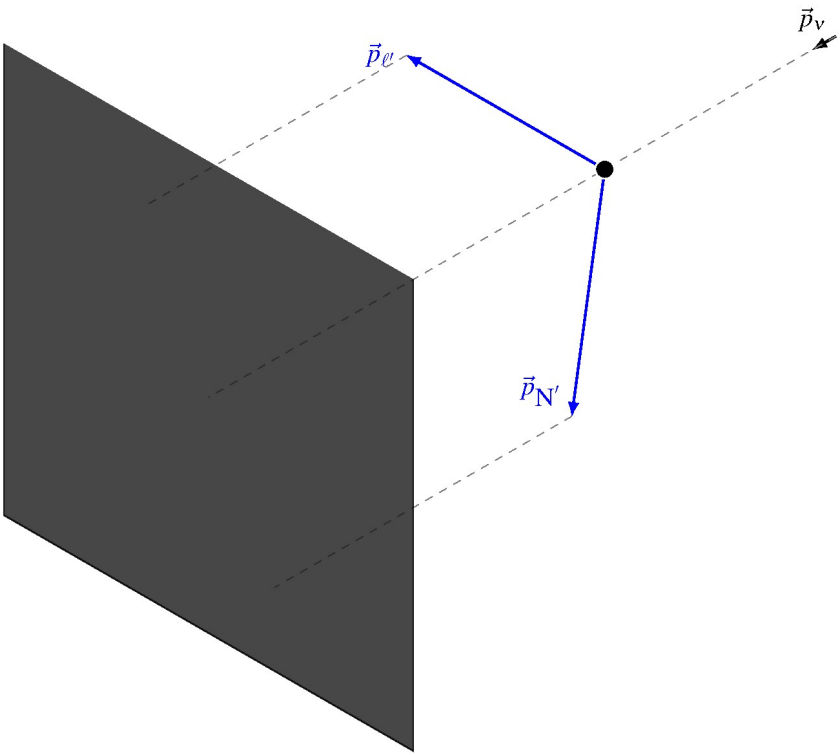
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To make *Neutrino Shadow Play*, we need

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- ✓ screen

Transverse kinematic imbalances – a neutrino shadow play

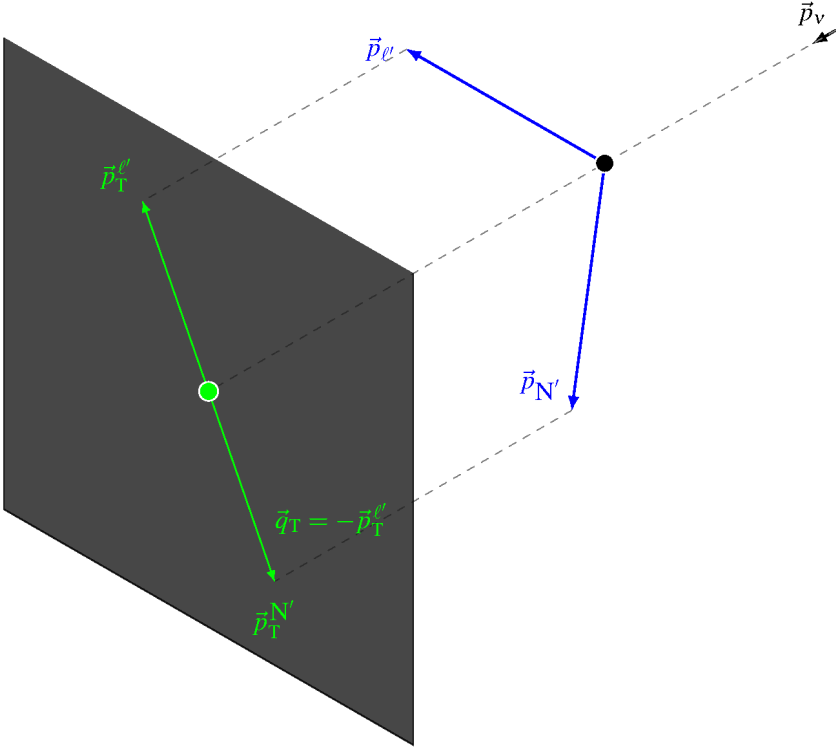


Source: <http://zhejiangpiying.sokutu.com/tupian.html>



- To make *Neutrino Shadow Play*, we need
- ✓ beam of light → accelerator
 - ✓ screen → transverse plane

Transverse kinematic imbalances – a neutrino shadow play



Static nucleon target

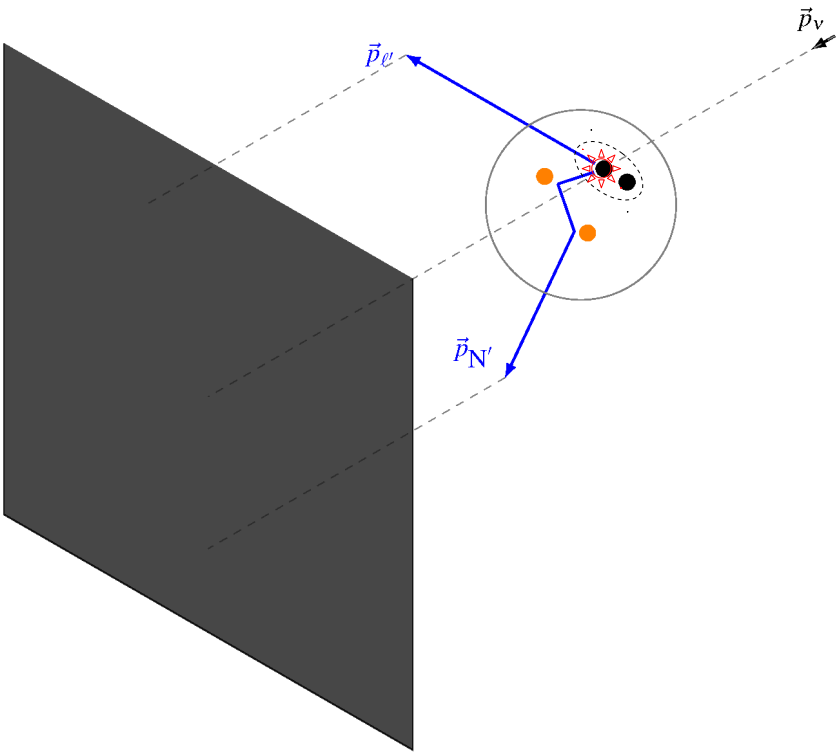


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Transverse kinematic imbalances – a neutrino shadow play



Nuclear target

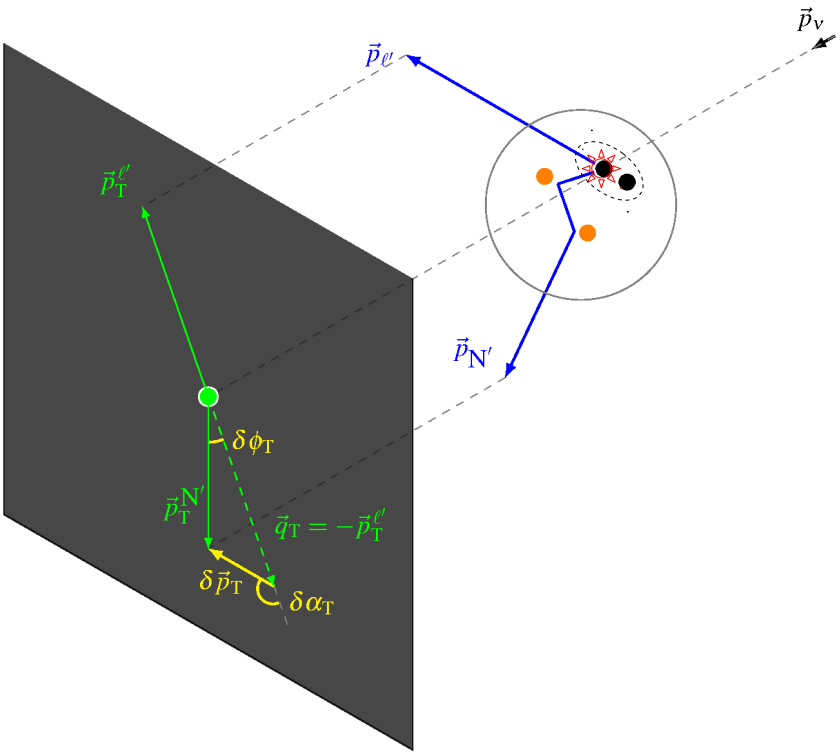


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Transverse kinematic imbalances – a neutrino shadow play



Nuclear target



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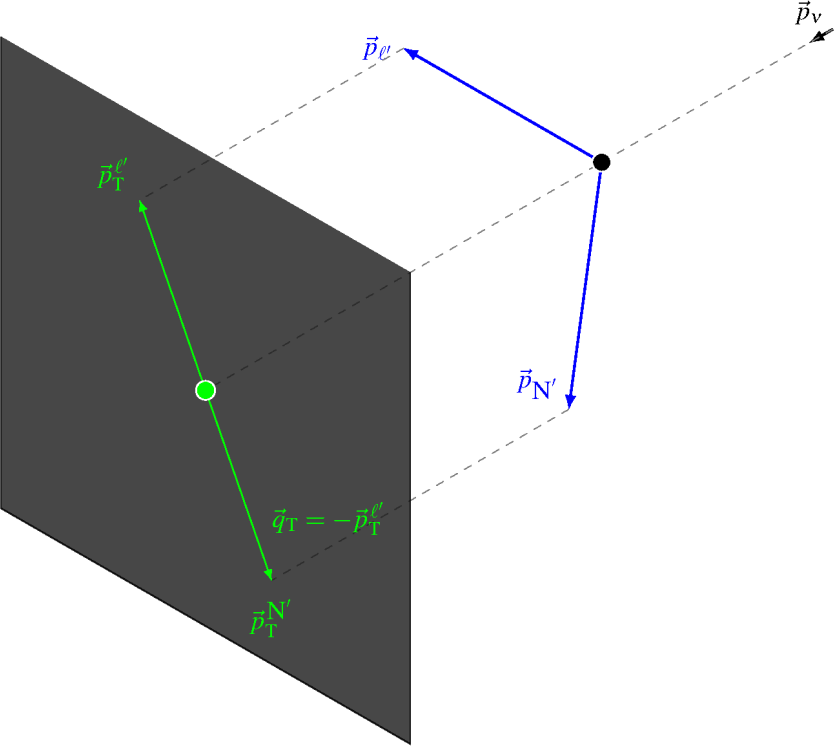


To make *Neutrino Shadow Play*, we need
 ✓ beam of light → accelerator
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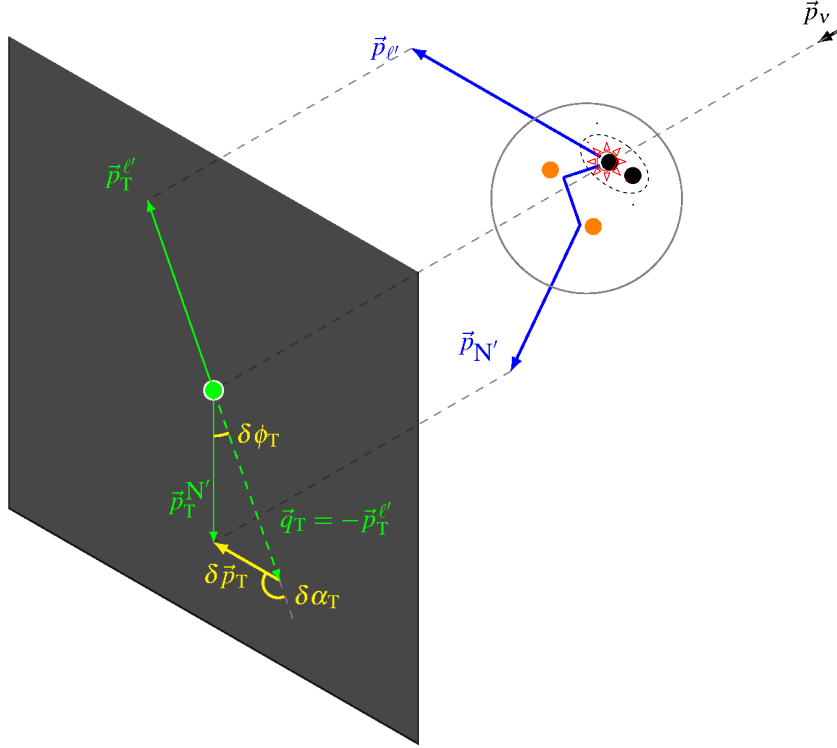
Transverse kinematic imbalances – a neutrino shadow play

$$\delta \vec{p}_T = \vec{p}_T^N - \Delta \vec{p}_T$$

Convolution of Fermi motion and intra-nuclear momentum transfer due to FSI, resonance production, 2p2h etc.



Static nucleon target



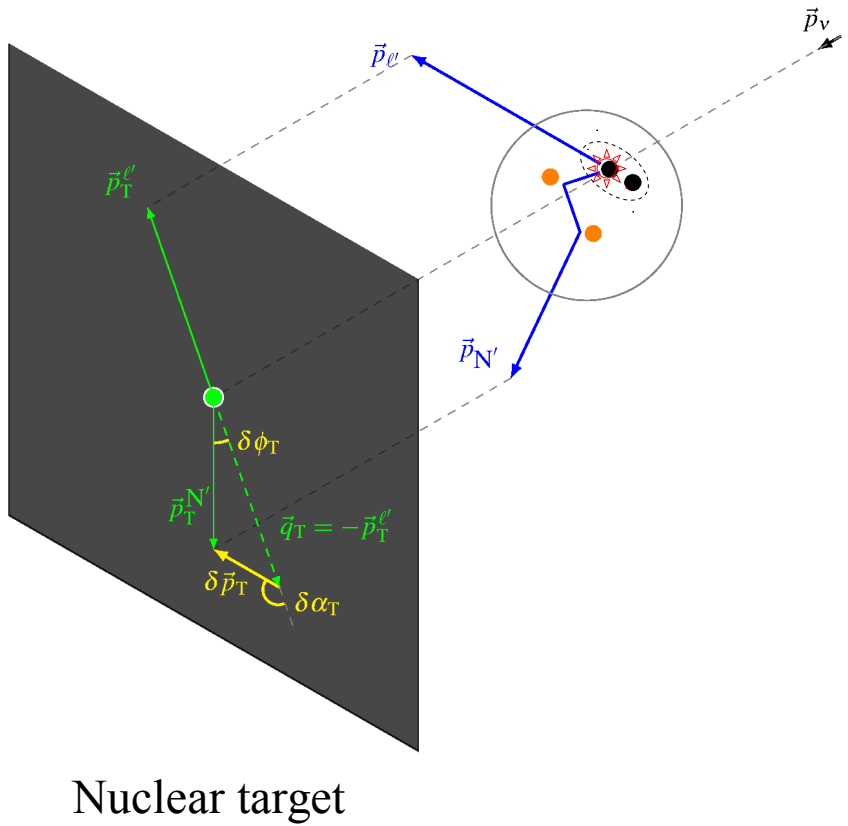
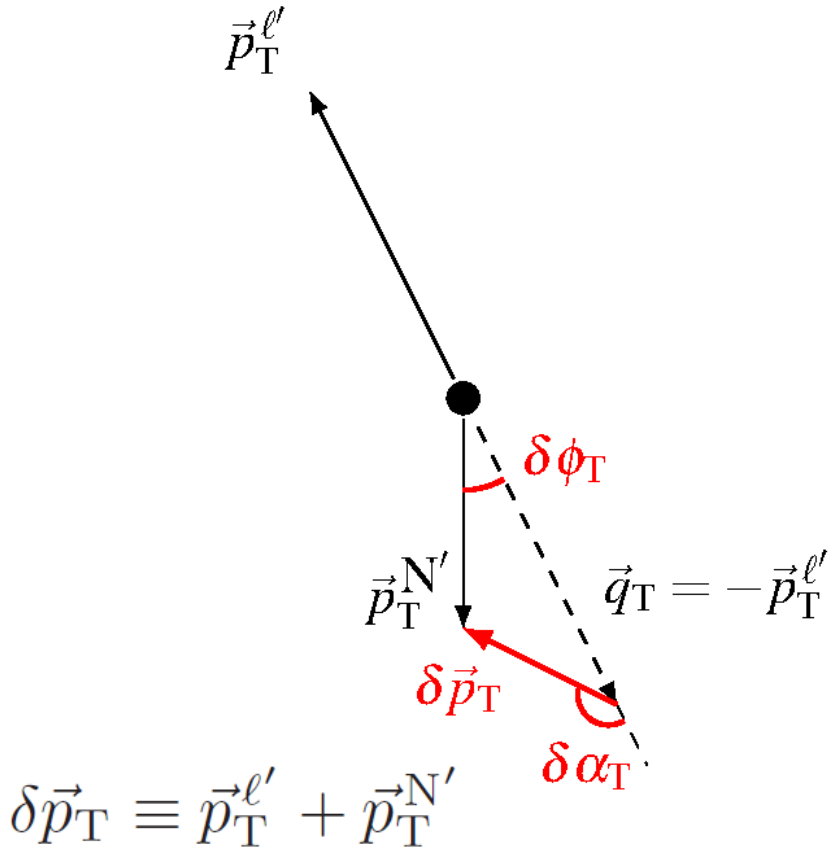
Nuclear target

XL, L. Pickering, S. Dolan *et al.*, Phys.Rev. C94 (2016) no.1, 015503

Transverse kinematic imbalances – a neutrino shadow play

$$\delta \vec{p}_T = \vec{p}_T^N - \Delta \vec{p}_T$$

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XL, L. Pickering, S. Dolan *et al.*, Phys.Rev. C94 (2016) no.1, 015503

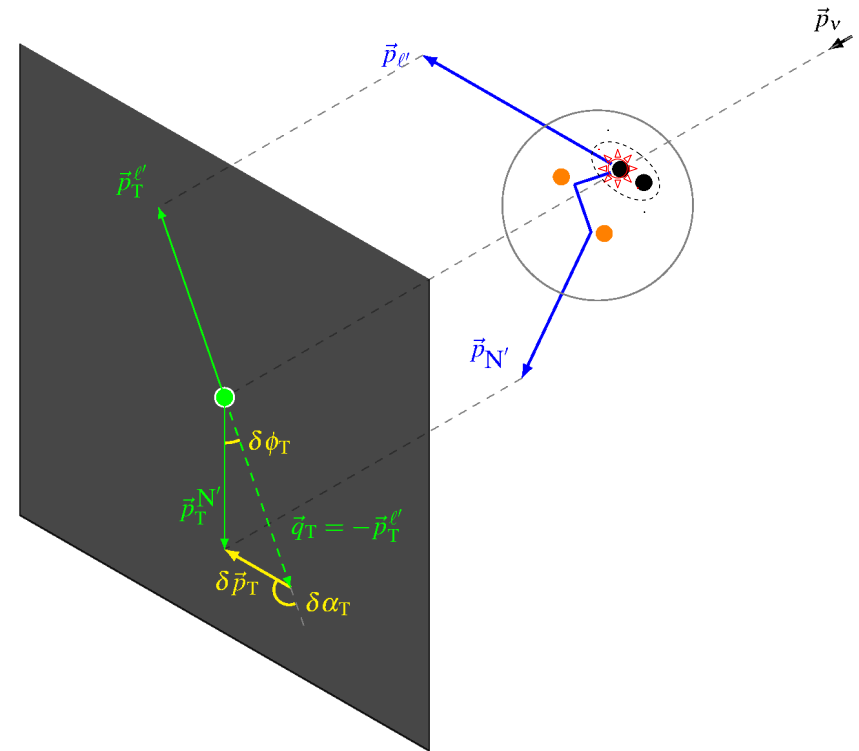
A more general analysis of kinematic imbalance

Transverse: $0 = \vec{p}_T^{\ell'} + \vec{p}_T^{N'} - \delta\vec{p}_T$

Longitudinal: $E_\nu = p_L^{\ell'} + p_L^{N'} - \delta p_L$

New variable: $p_n \equiv \sqrt{\delta p_T^2 + \delta p_L^2}$

Neutrino energy is unknown (in the first place), equations are not closed.



A more general analysis of kinematic imbalance

Transverse: $0 = \vec{p}_T^{\ell'} + \vec{p}_T^{N'} - \delta\vec{p}_T$

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Neutrino energy is unknown (in the first place), equations are not closed.

Assuming exclusive μ -p-A' final states

Use energy conservation to close the equations

$$E_\nu + m_A = E_{\ell'} + E_{N'} + E_{A'}$$

$$E_{A'} = \sqrt{m_{A'}^2 + p_n^2}$$

p_n : recoil momentum of the nuclear remnant

A more general analysis of kinematic imbalance

Transverse: $0 = \vec{p}_T^{\ell'} + \vec{p}_T^{N'} - \delta\vec{p}_T$

Longitudinal: $E_\nu = p_L^{\ell'} + p_L^{N'} - \delta p_L$

New variable: $p_n \equiv \sqrt{\delta p_T^2 + \delta p_L^2}$

Neutrino energy is unknown (in the first place), equations are not closed.

For CCQE with elastic FSI, $A' = {}^{11}\text{C}^*$
 No more unknowns
 p_n : neutron Fermi motion *weakly smeared

A. Furmanski, J. Sobczyk, Phys.Rev. C95 (2017) no.6, 065501

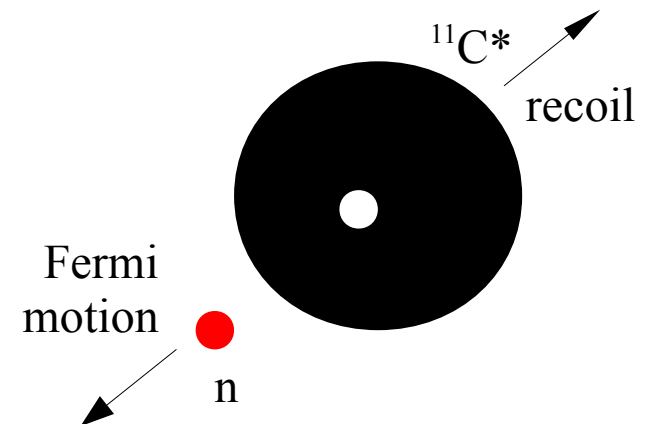
Xianguo Lu, Oxford

Assuming exclusive μ -p-A' final states
 Use energy conservation to close the equations

$$E_\nu + m_A = E_{\ell'} + E_{N'} + E_{A'}$$

$$E_{A'} = \sqrt{m_{A'}^2 + p_n^2}$$

p_n : recoil momentum of the nuclear remnant



$$p_L = \frac{1}{2}(M(A) + k'_L + p'_L - E' - E_{p'})$$

$$- \frac{p_T^2 + M^*(A-1)^2}{2(M(A) + k'_L + p'_L - E' - E_{p'})}$$

Measurement of final-state correlations in neutrino charged-current muon-proton mesonless production on hydrocarbon at $\langle E_\nu \rangle = 3 \text{ GeV}$

Signal definition:

- Charged current
- One muon and at least one proton in the restricted final-state phase space
- No mesons

$$1.5 \text{ GeV}/c < p_\mu < 10 \text{ GeV}/c, \theta_\mu < 20^\circ,$$
$$0.45 \text{ GeV}/c < p_p < 1.2 \text{ GeV}/c, \theta_p < 70^\circ$$

Measurement:

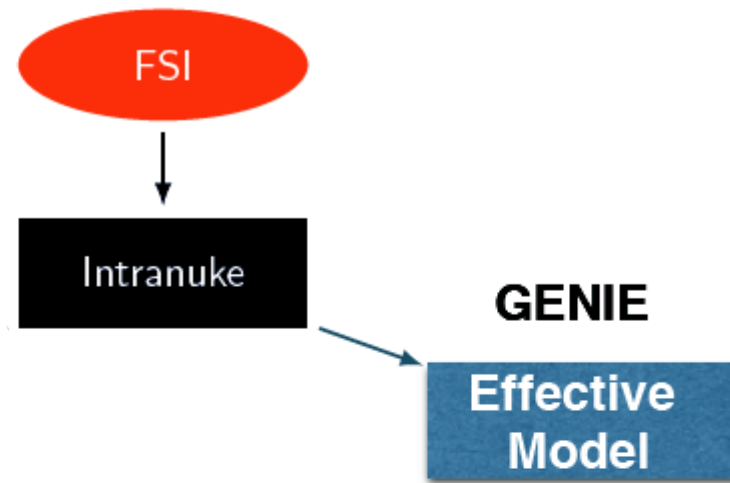
Data sample: NuMI low energy neutrino data, 3.28×10^{20} POT

Interaction target: tracker (mostly CH)

- Event selection
 - Background estimation and subtraction
 - Unfolding
 - Efficiency correction
- Flux integrated cross section as results
- Focus today

Simulation: GENIE [Nucl.Instrum.Meth. A614 (2010) 87-104]

- **Nominal:** version 2.8.4
 - ✓ global Fermi Gas (RFG) model with Bodek-Ritchie (BR) tail [Phys. Rev. D 23, 1070 (1981)]
 - ✓ hA FSI [AIP Conf.Proc. 1405 (2011) 213-218]
- **No-FSI:** Nominal without FSI

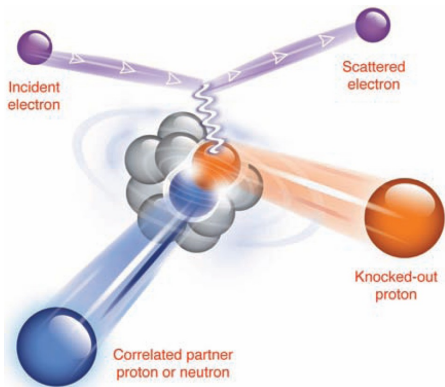


- INC-like with one “effective” interaction
- tuned do hadron-nucleus data
- easy to reweight

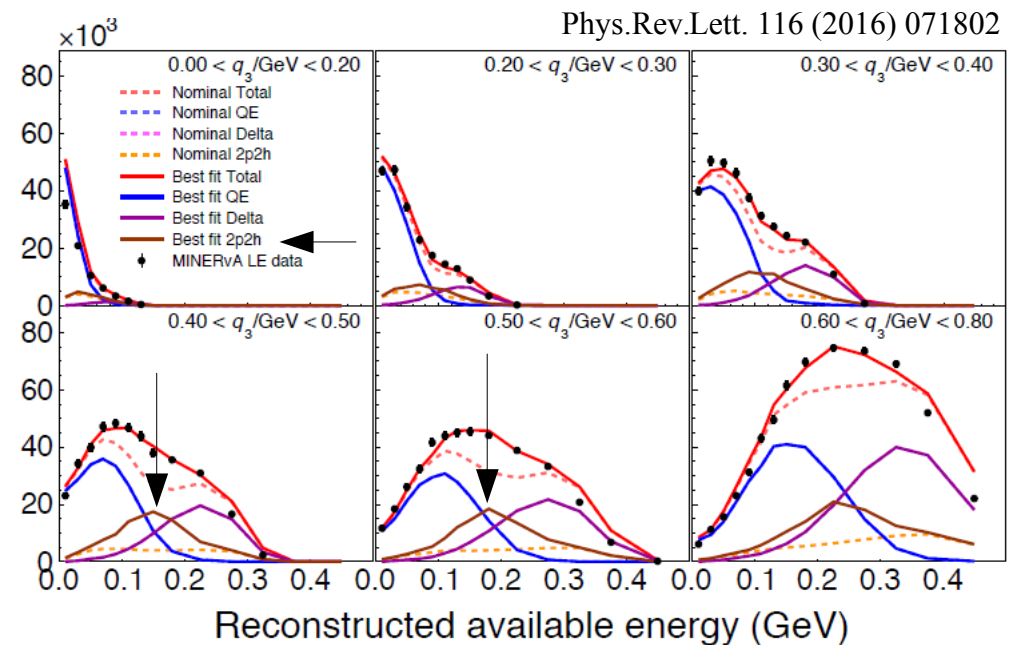
courtesy of Tomasz Golan

Simulation: GENIE [Nucl.Instrum.Meth. A614 (2010) 87-104]

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- **No-FSI:** Nominal without FSI
- **MnvGENIE-v1: GENIE MINERvA Tune (v1) [only 2p2h relevant for this analysis]**
 - Added Random Phase Approximation (RPA) [Phys.Rev. C70 (2004) 055503]
 - Non-resonance pion production scaled down by 75% [Phys.Rev. D90 (2014) no.11, 112017]
 - **Valencia 2p2h** [Nieves *et al.*, Phys.Lett. B707 (2012) 72-75, Phys. Rev. C 86, 015504 (2012), Phys.Rev. D88 (2013) no.11, 113007, arXiv:1601.02038]
 - ✓ **tuned to MINERvA inclusive data** → significant enhancement in small 4-momentum transfer region [Phys.Rev.Lett. 116 (2016) 071802]



Science 320 (2008) 1476-1478



→ representing energy transfer from the neutrino to the target

Simulation: GENIE [Nucl.Instrum.Meth. A614 (2010) 87-104]

- **Nominal:** version 2.8.4
 - ✓ global Fermi Gas (RFG) model with Bodek-Ritchie (BR) tail [Phys. Rev. D 23, 1070 (1981)]
 - ✓ hA FSI [AIP Conf.Proc. 1405 (2011) 213-218]
- **No-FSI:** Nominal without FSI
- **MnvGENIE-v1: GENIE MINERvA Tune (v1) [only 2p2h relevant for this analysis]**
 - Added Random Phase Approximation (RPA) [Phys.Rev. C70 (2004) 055503]
 - Non-resonance pion production scaled down by 75% [Phys.Rev. D90 (2014) no.11, 112017]
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 - ✓ tuned to MINERvA inclusive data → significant enhancement in small 4-momentum transfer region [Phys.Rev.Lett. 116 (2016) 071802]

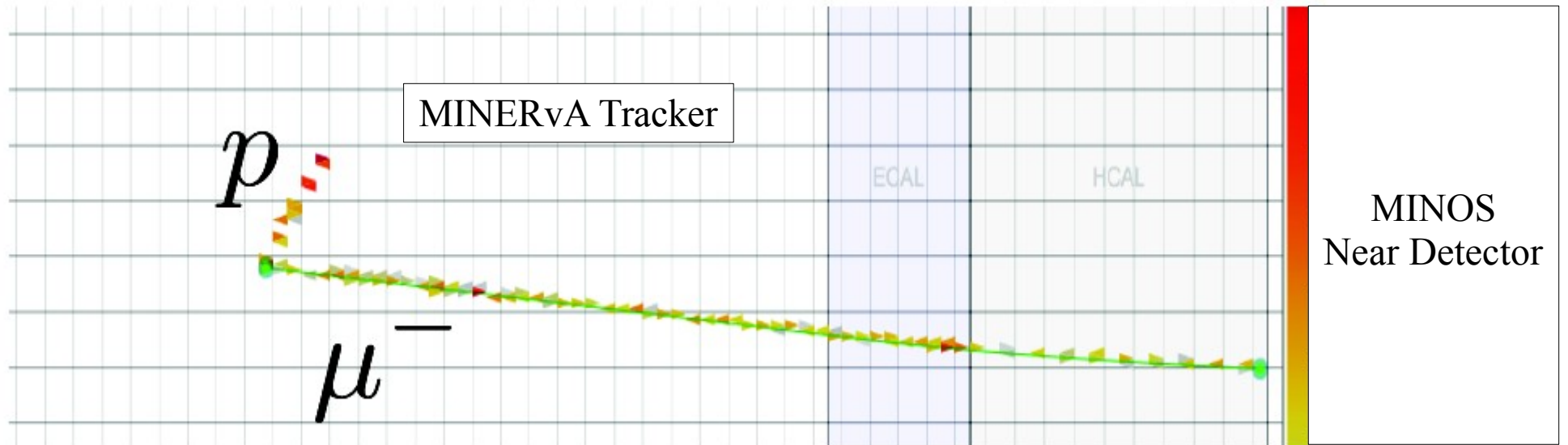
Detector simulation: GEANT4 (4.9.2)

GENIE used in other experiments (e.g. NOvA, T2K, μ BooNE, DUNE)

This analysis:
GENIE MINERvA Tune (v1) used in cross section extraction

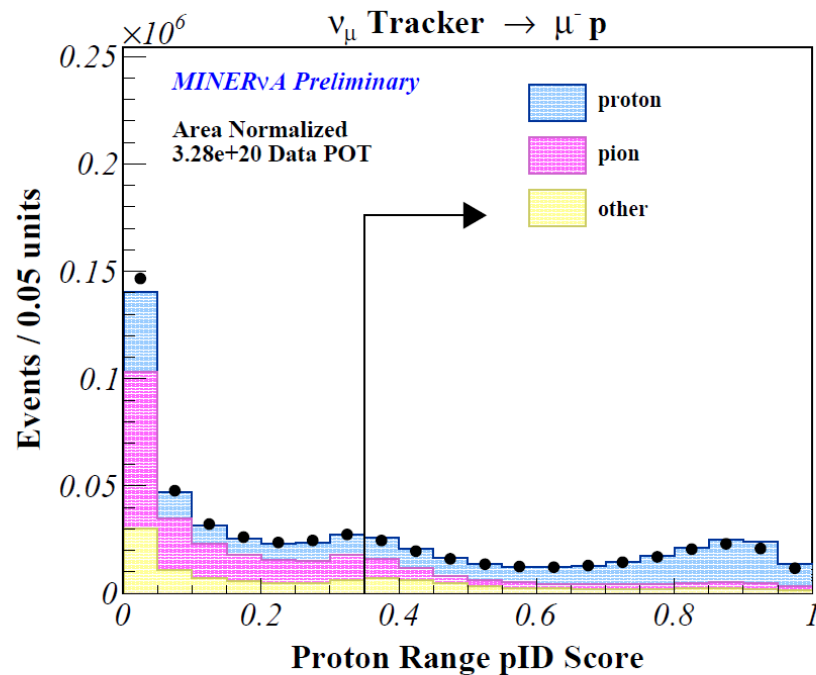
Event Selection

- One muon candidate track matched to a MINOS track
- At least one proton candidate (particle identification using dE/dx along the track)
 - Elastically Scattered Contained (ESC) proton selection
- Vertex in tracker
- Michel electron (from pion-muon-electron decay chain) tag to remove pion production
- Cut on energy far from vertex (unattached visible energy) to remove events with untracked pions



Event Selection

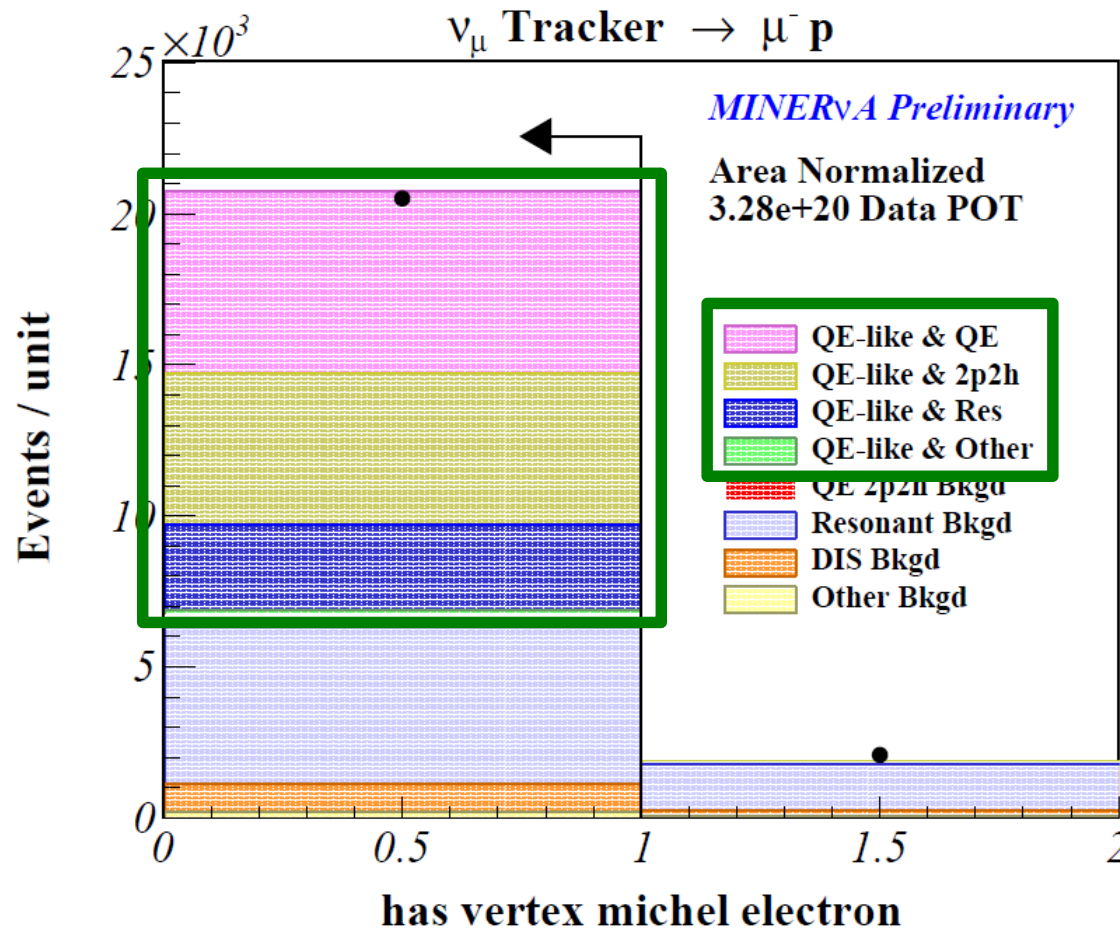
- One muon candidate track matched to a MINOS track
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- Vertex in tracker
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- Cut on energy far from vertex (unattached visible energy) to remove events with untracked pions



Based on dE/dx profile along the track

Event Selection

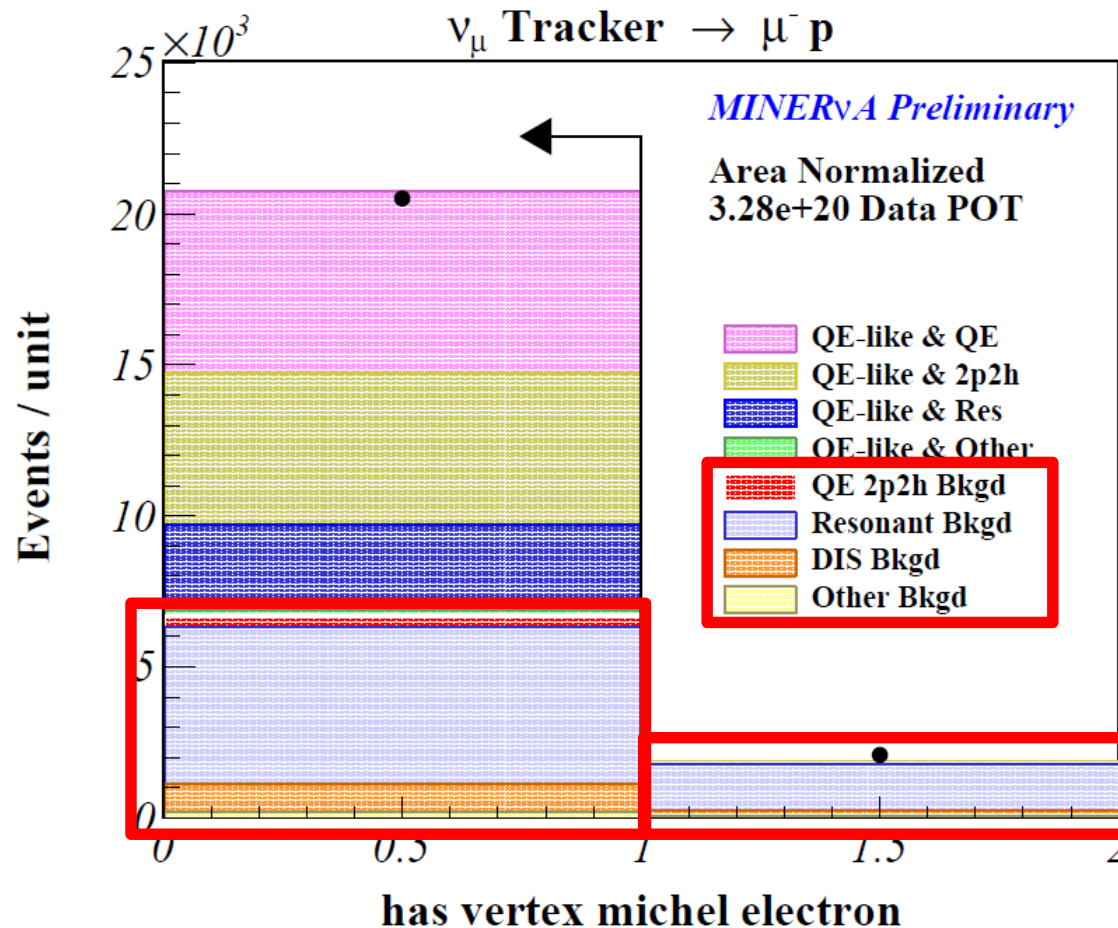
- One muon candidate track matched to a MINOS track
- At least one proton candidate (particle identification using dE/dx along the track)
 - Elastically Scattered Contained (ESC) proton selection
- Vertex in tracker
- Michel electron (from pion-muon-electron decay chain) tag to remove pion production
- Cut on energy far from vertex (unattached visible energy) to remove events with untracked pions



Signal:
QE-like events

Event Selection

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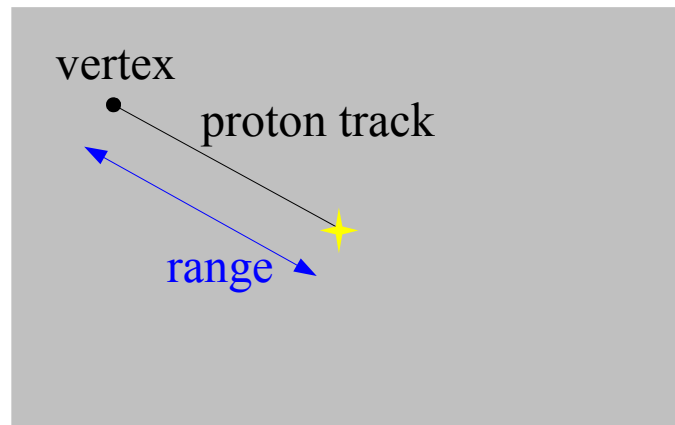


Background

Event Selection

- One muon candidate track matched to a MINOS track
- At least one proton candidate (particle identification using dE/dx along the track)
 - Elastically Scattered Contained (ESC) proton selection → **new development**
- Vertex in tracker
- Michel electron (from pion-muon-electron decay chain) tag to remove pion production
- Cut on energy far from vertex (unattached visible energy) to remove events with untracked pions

Homogeneous non-magnetized tracker
Momentum by range

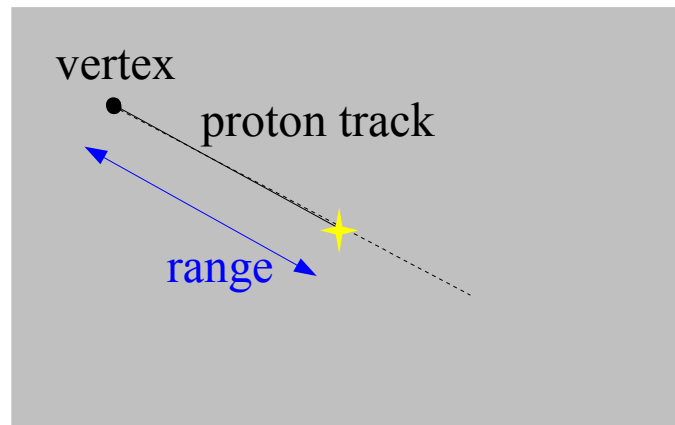


Momentum-range correlation best known when the track has “peaceful” end: stopped elastically

Event Selection

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Homogeneous non-magnetized tracker
Momentum by range



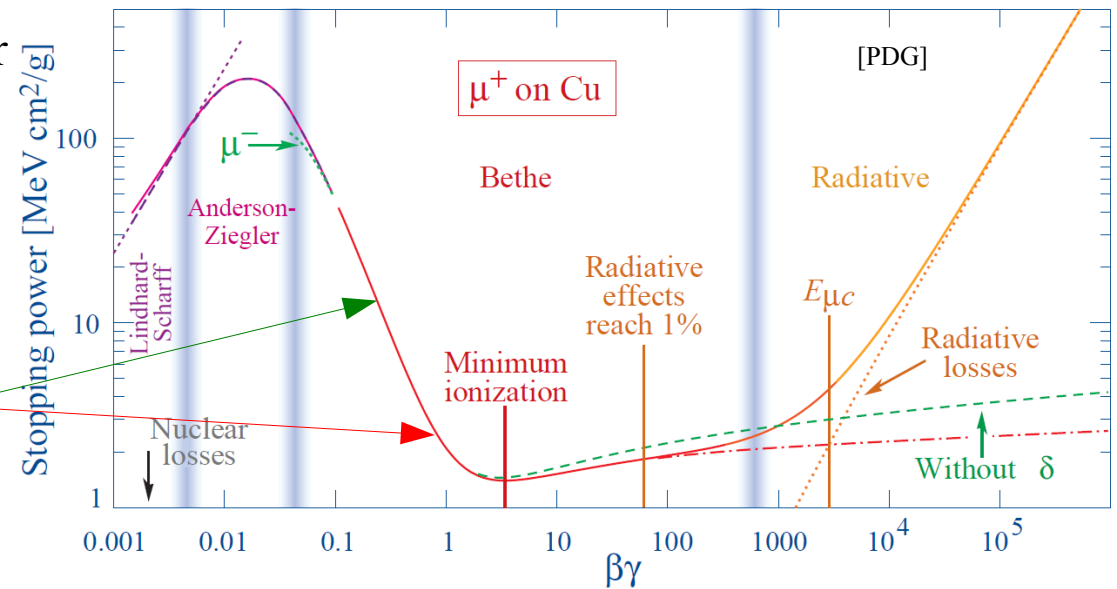
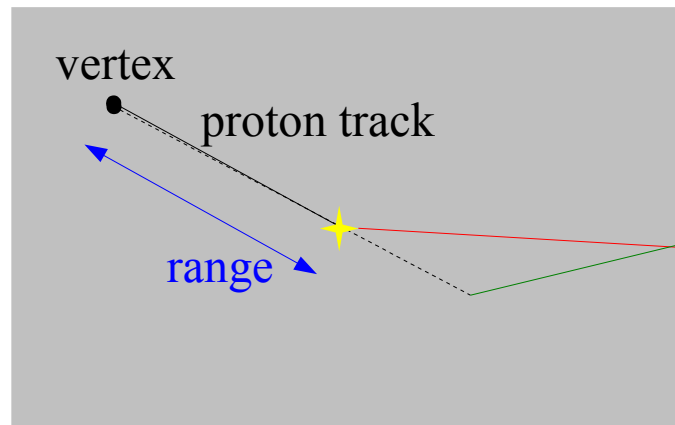
Momentum-range correlation best known when the track has “peaceful” end: stopped elastically

If track ends on the fly due to inelastic interaction in detector (e.g. $p A \rightarrow n A'$)
Range can only be measured prematurely → large bias in momentum estimation

Event Selection

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 - Elastically Scattered Contained (ESC) proton selection → **new development**
- Vertex in tracker
- Michel electron (from pion-muon-electron decay chain) tag to remove pion production
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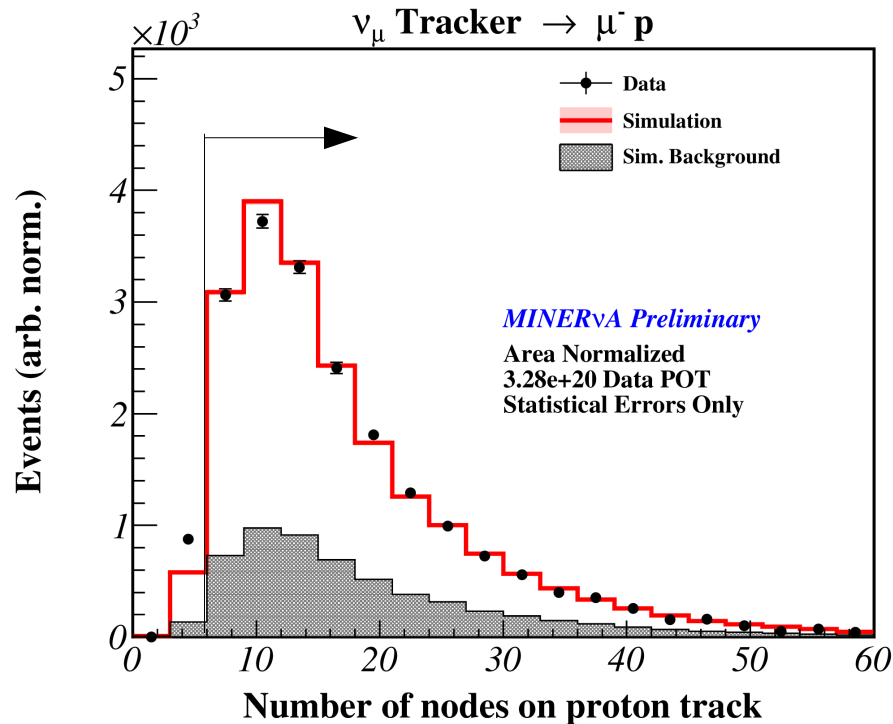
Homogeneous non-magnetized tracker
Momentum by range



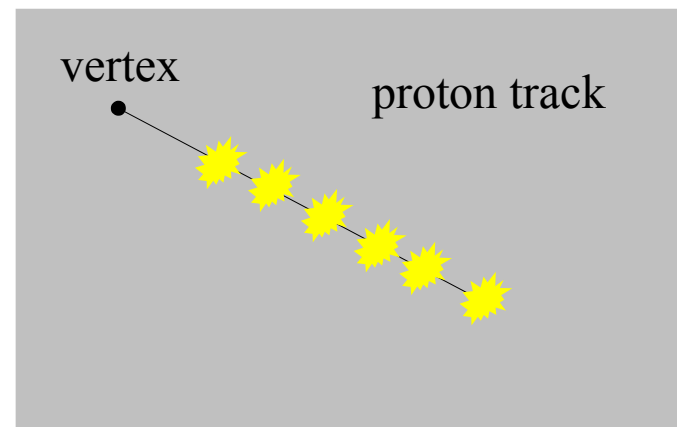
Proton stopped on the fly have smaller dE/dx
→ Cut on dE/dx from track end point

Event Selection

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- At least one proton candidate (particle identification using dE/dx along the track)
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- Vertex in tracker
- Michel electron (from pion-muon-electron decay chain) tag to remove pion production
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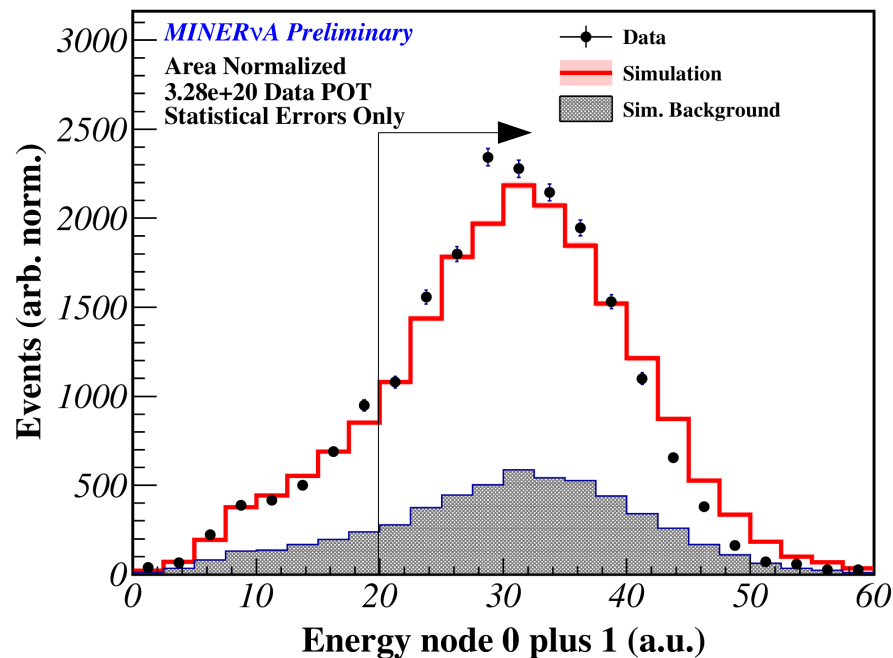


- Require at least 6 dE/dx nodes from track end point

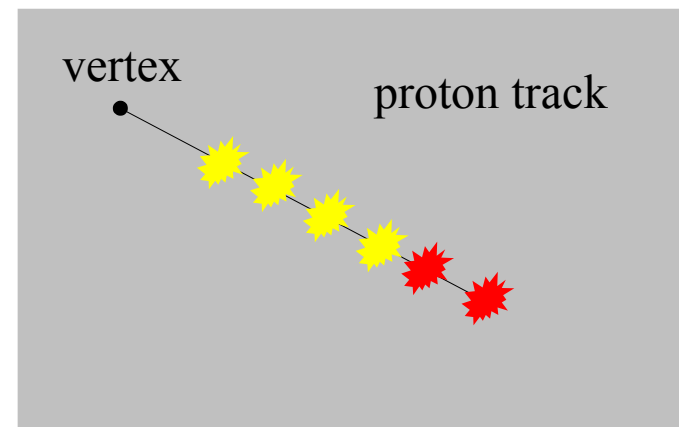


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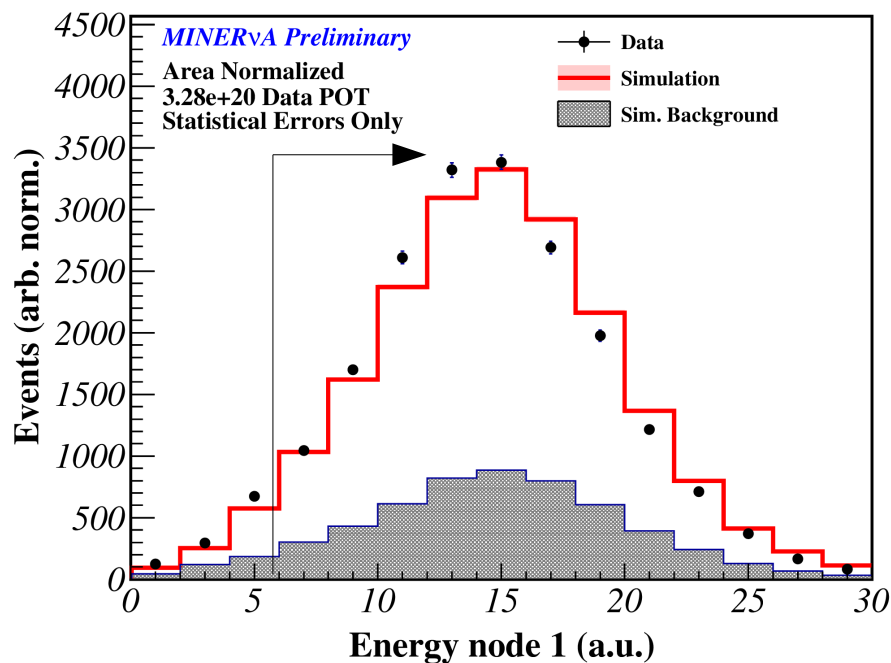


- Require at least 6 dE/dx nodes from track end point
- Cut on summed dE/dx of last two nodes (node 0 and 1 correlated)

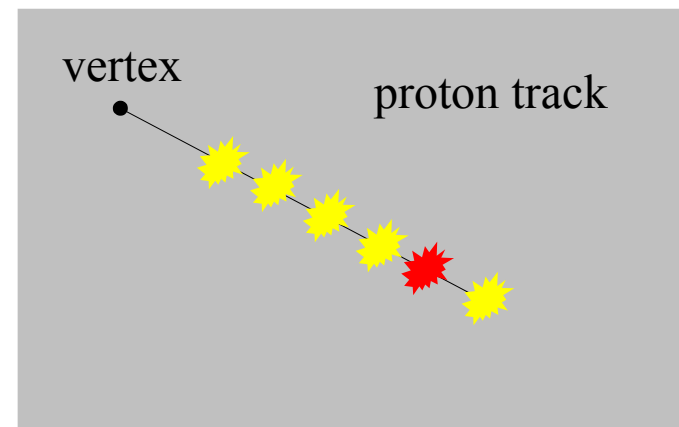


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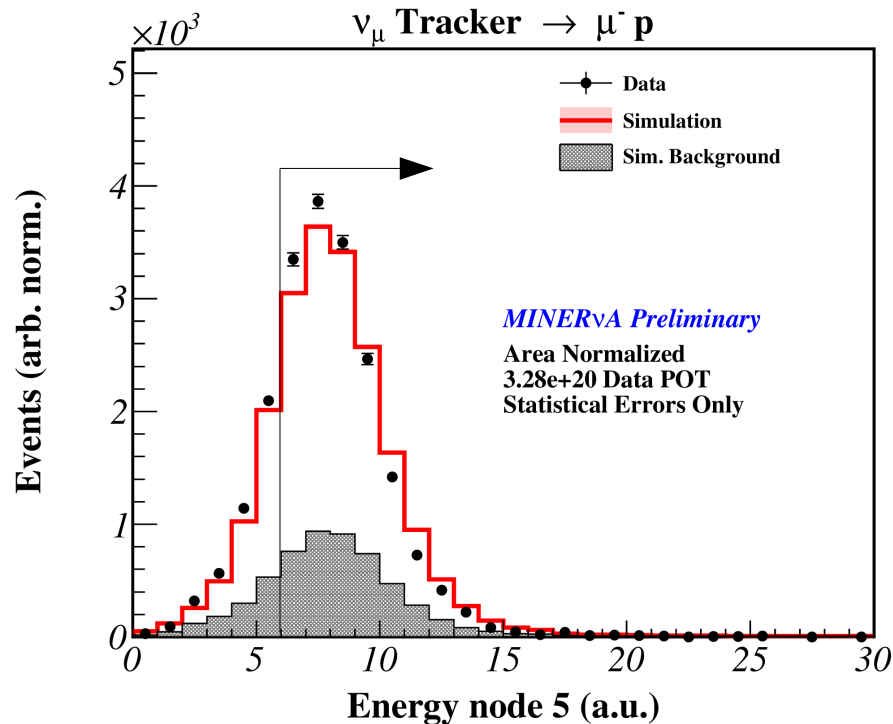


- Require at least 6 dE/dx nodes from track end point
- Cut on summed dE/dx of last two nodes (node 0 and 1 correlated)
- Cut on individual nodes (node 1-5)

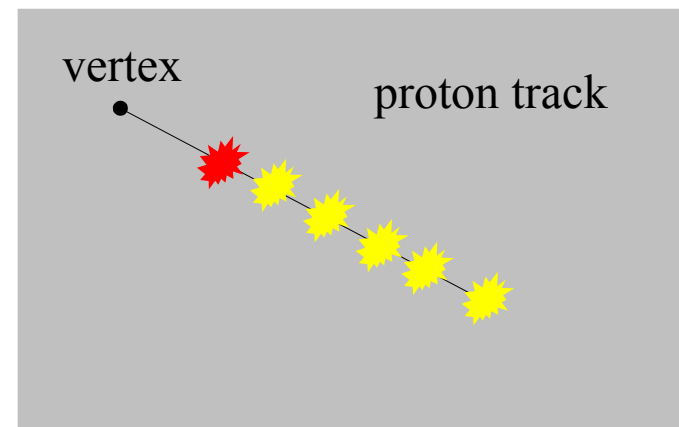


Event Selection

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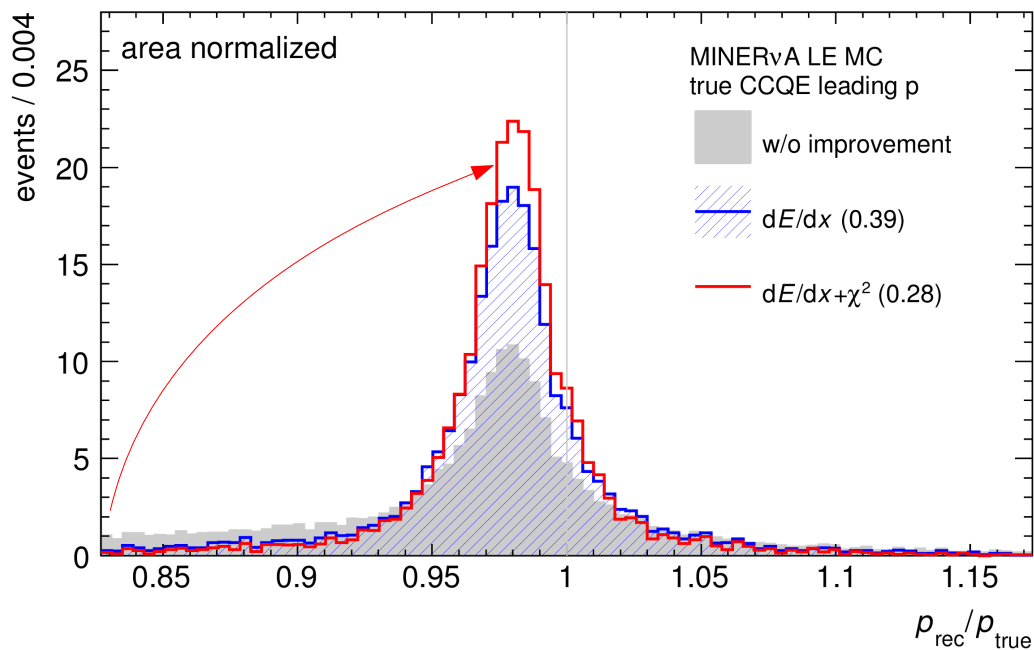


- Require at least 6 dE/dx nodes from track end point
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Event Selection

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* dE/dx +cleanup cut efficiency 30-40-20% @ 0.6-0.75-1 GeV/c

ESC proton selection:

- Cut efficiency $\sim 40\%$
- Reconstructed momentum spread much reduced @ 0.7 – 1.1 GeV, resolution 3% \sim 2%
- 5-10% uncertainty in efficiency

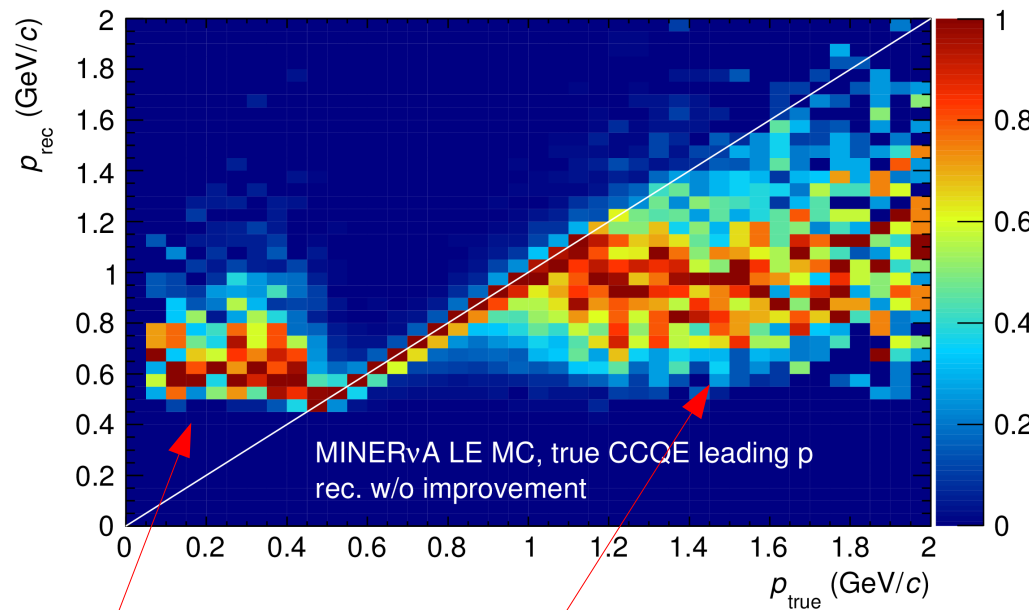
Clean-up cuts to improve proton and muon momentum resolution:

- proton dE/dx profile χ^2
- number of MINOS track nodes

Also need to correct p_T scales of both muon and protons.

Event Selection

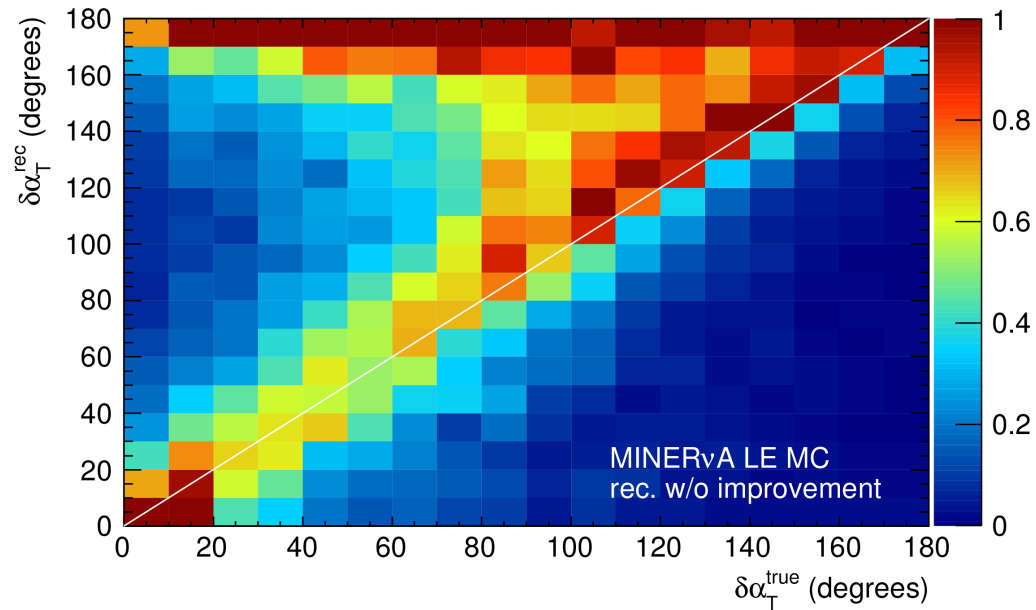
- One muon candidate track matched to a MINOS track
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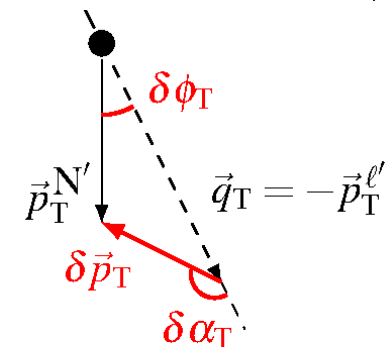
Short track
around vertex

Inelastic
scattered
protons

Xianguo Lu, Oxford

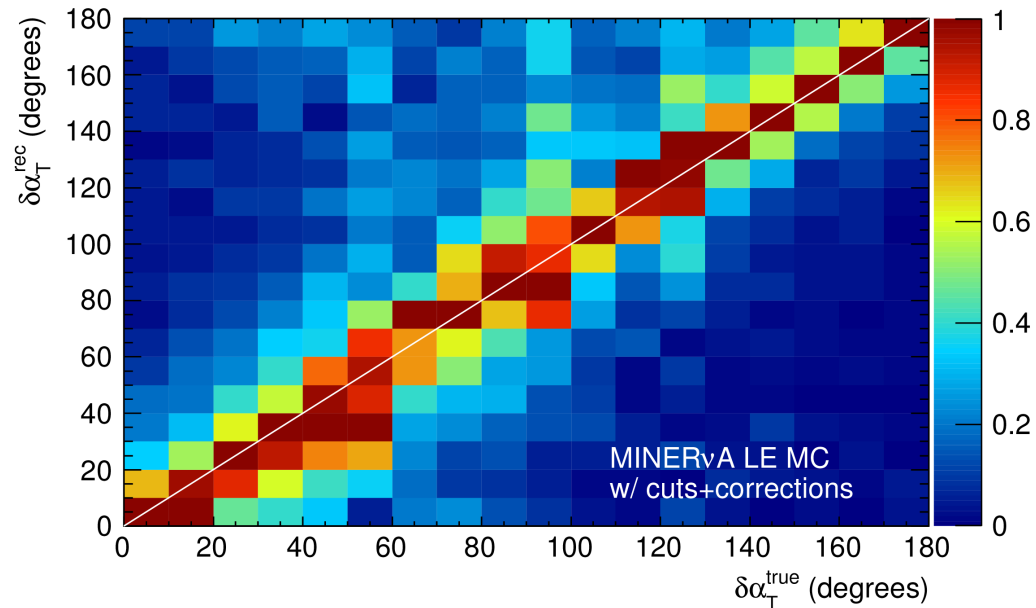
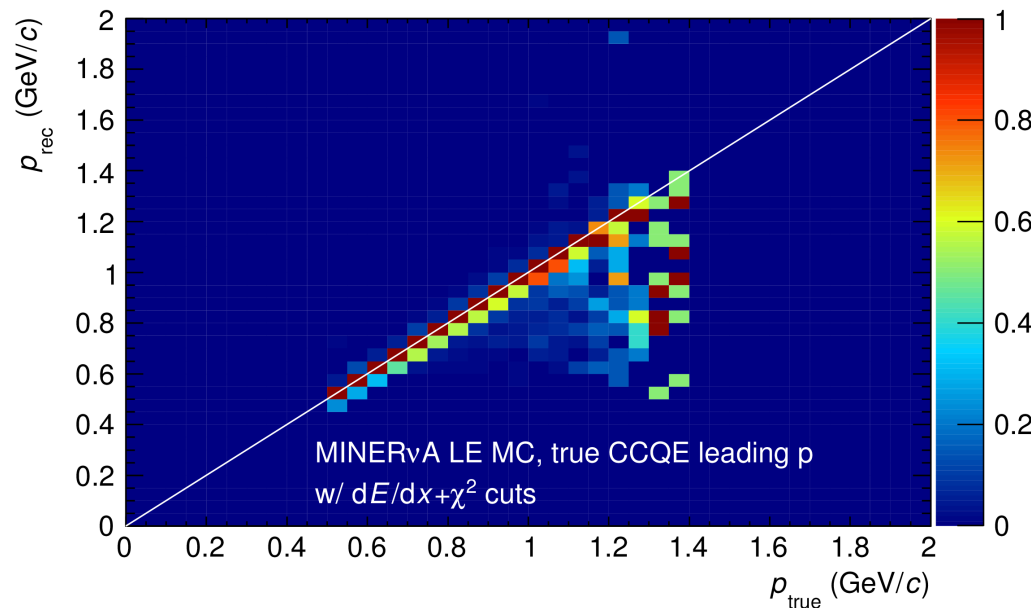


Reconstructed as a
function of true quantity
before ESC proton
selection for proton
momentum and $\delta\alpha_T$

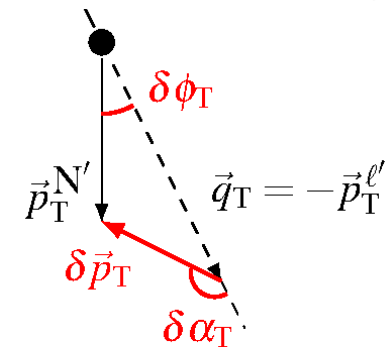


Event Selection

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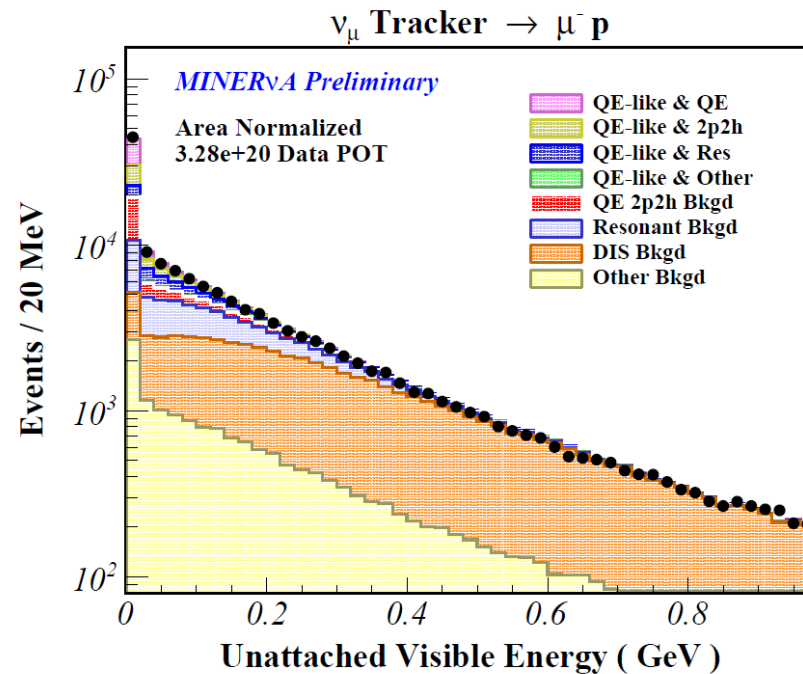


Reconstructed as a function of true quantity **after** ESC proton selection and p_T -scale correction (only applied to derived variables) for proton momentum and $\delta\alpha_T$



Event Selection

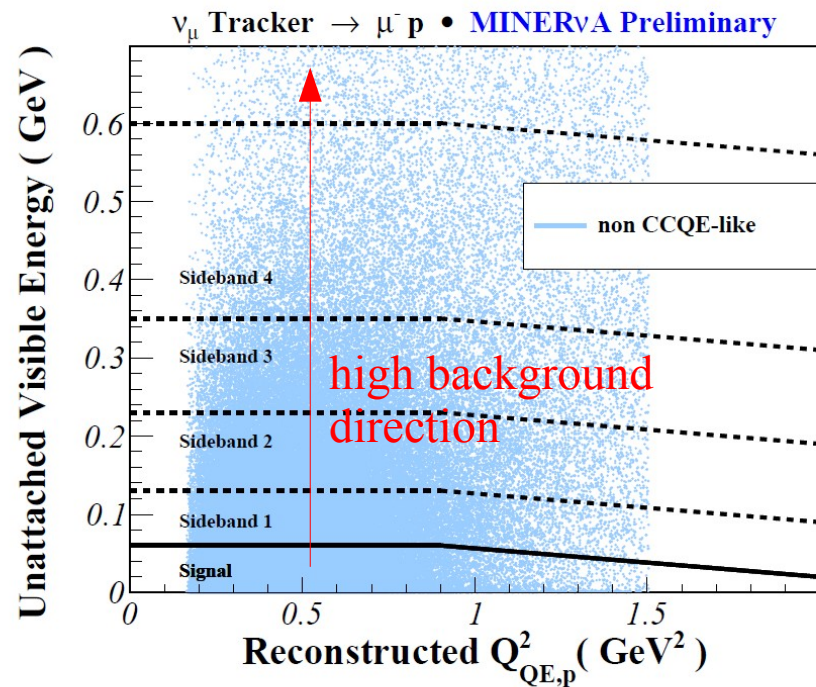
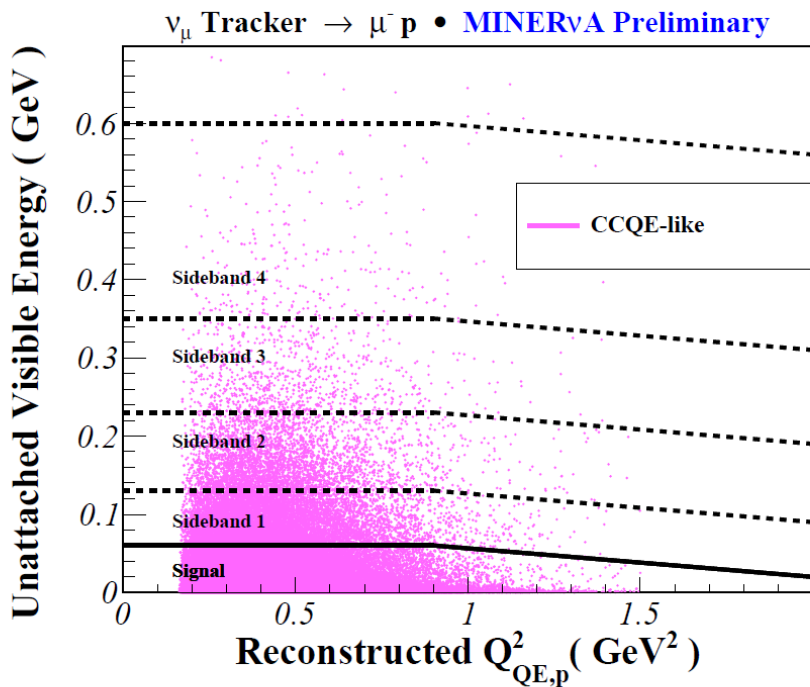
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Large unattached visible energy dominated by background.

Event Selection

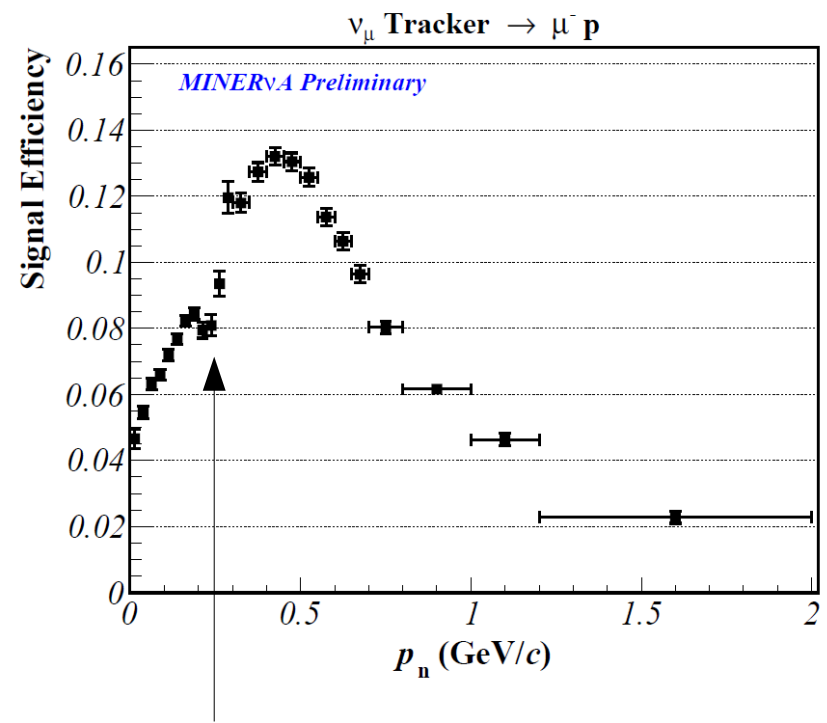
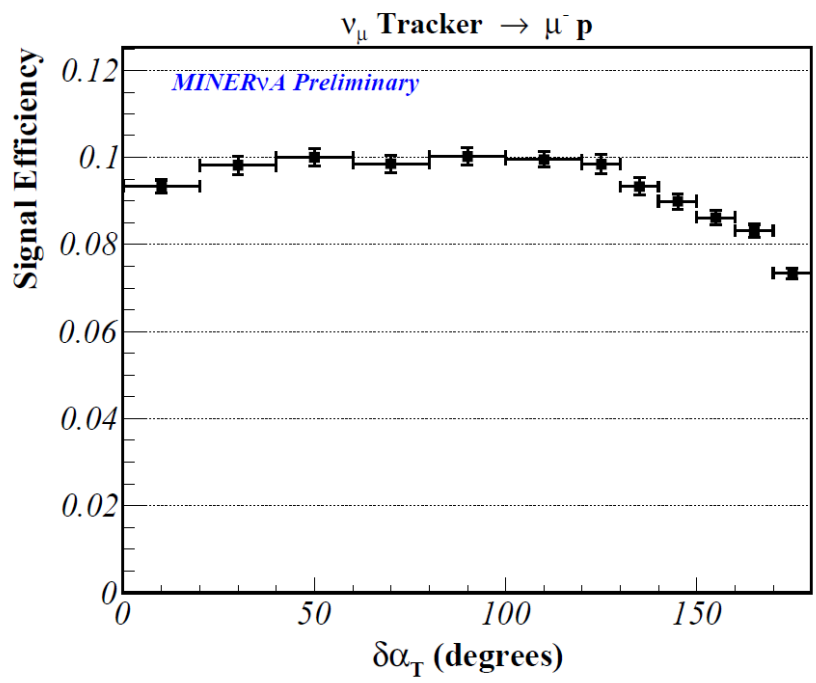
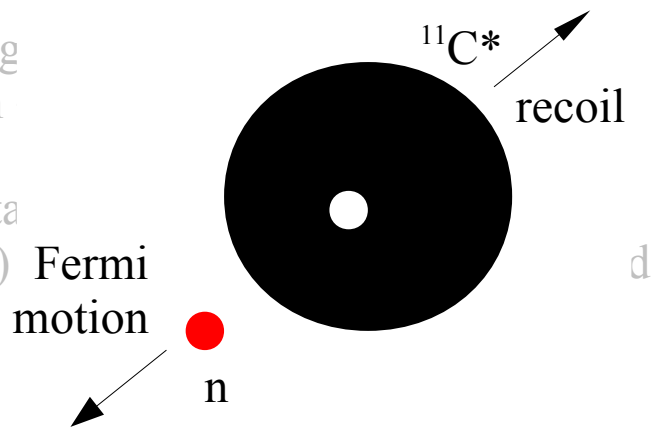
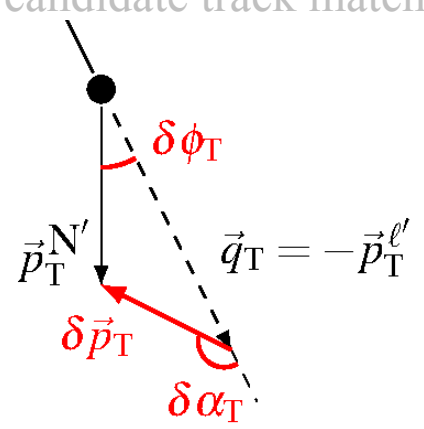
- One muon candidate track matched to a MINOS track
- At least one proton candidate (particle identification using dE/dx along the track)
 - Elastically Scattered Contained (ESC) proton selection \rightarrow new development
- Vertex in tracker
- Michel electron (from pion-muon-electron decay chain) tag to remove pion production
- Cut on energy far from vertex (unattached visible energy) to remove events with untracked pions



Signal and sidebands are defined

Event Selection

- One muon candidate track matched to a MINOS track
- At least one proton identification using (ESC) proton selection
- Vertex in the detector
- Michel electron (from muon decay chain) tagged
- Cut on energy of attached visible energy) Fermi motion
- Cut on energy of attached visible energy) Fermi motion

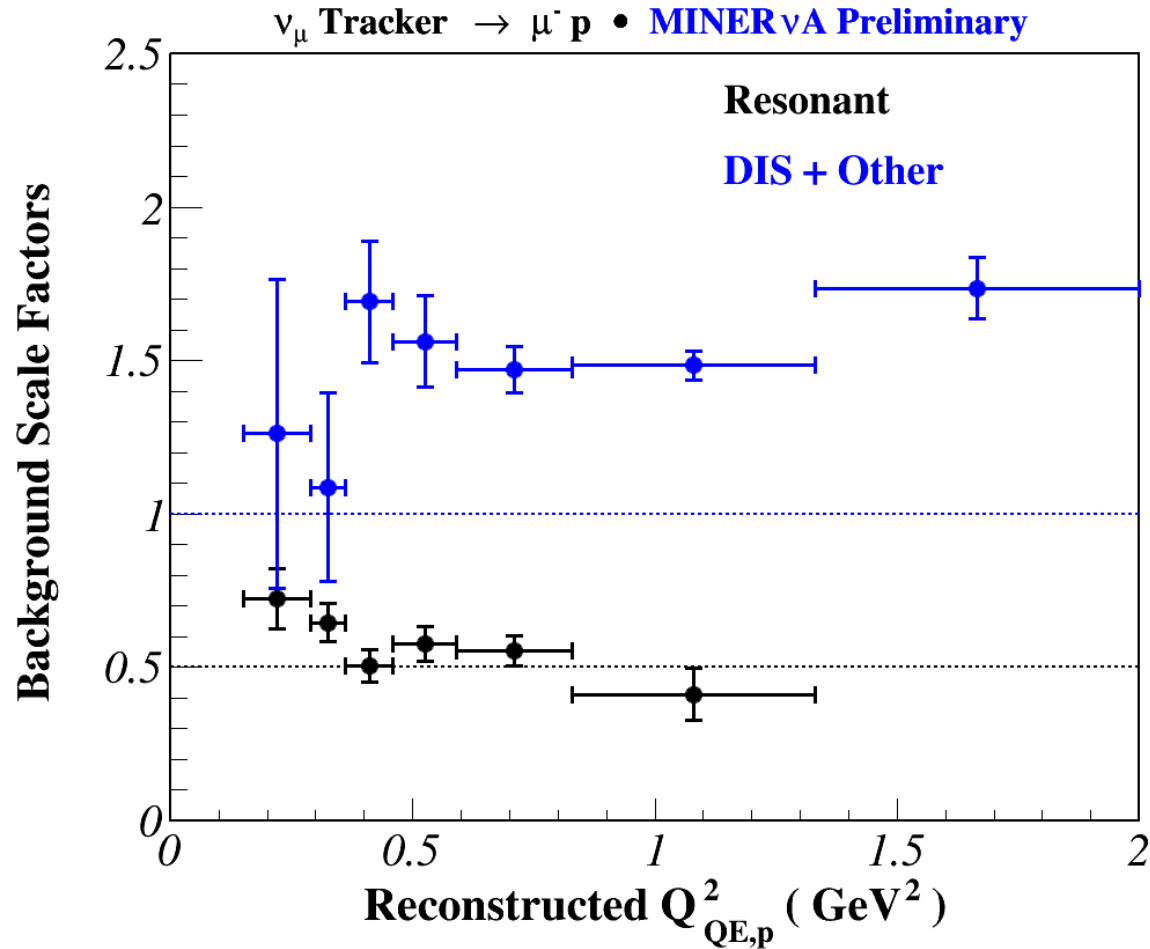


Overall efficiency: $\sim 9\%$

Peak region (see later slides)

Background Estimation

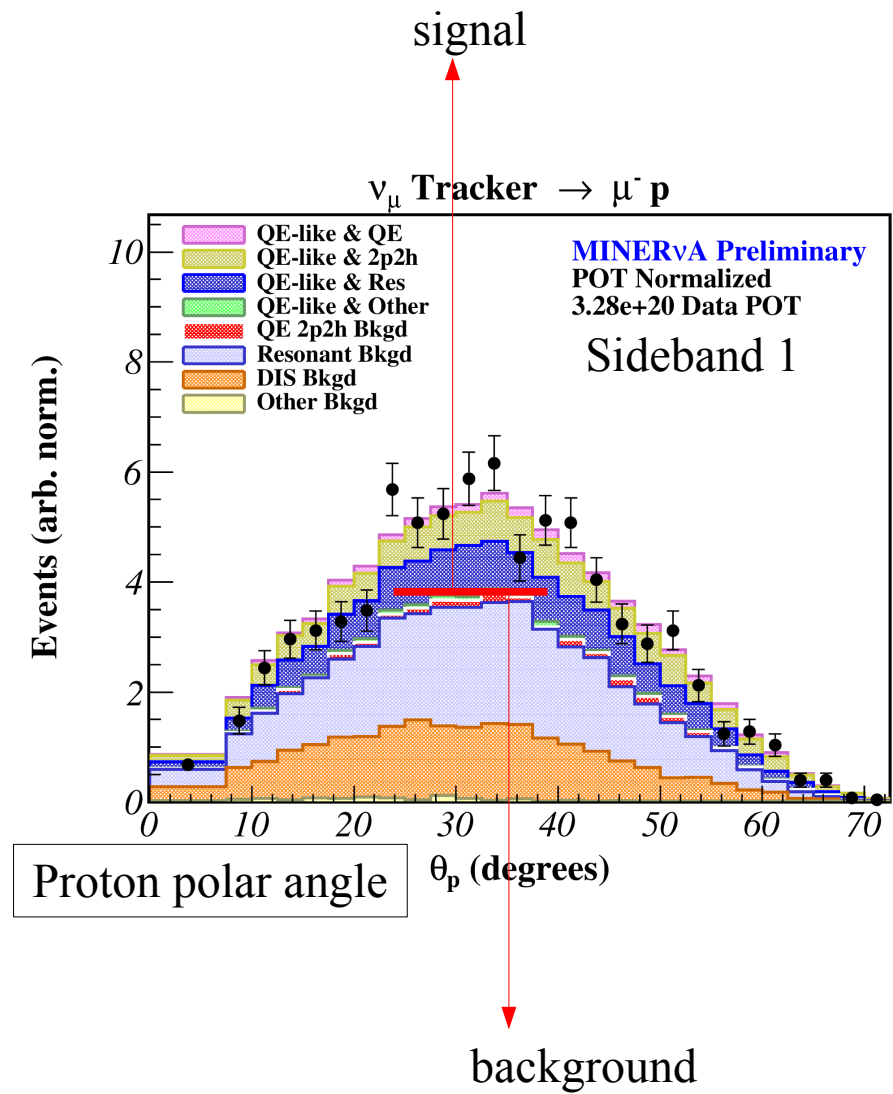
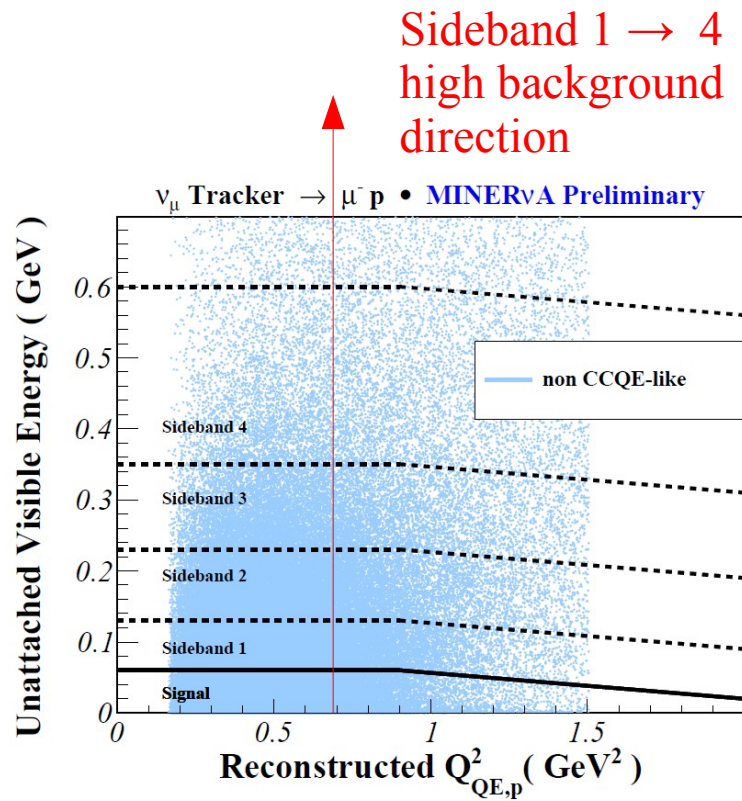
- Data driven sideband fit



Background in MC are rescaled according to data driven sideband fit

Background Estimation

- Data-MC comparison at reconstructed level after sideband fit



Selected Sample

- Data-MC comparison at reconstructed level after sideband fit

Selected sample

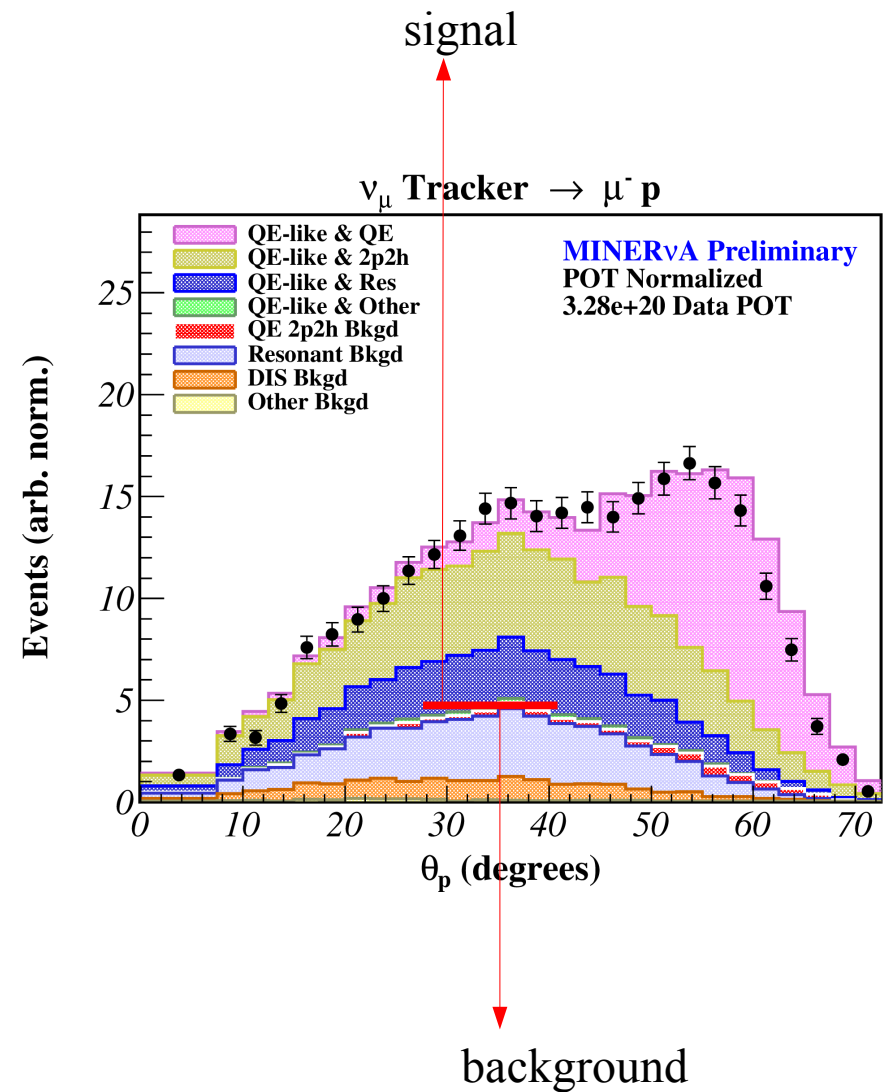
Overall purity: 78%

Overall QE-2p2h bkgd fraction: 3%

Overall RES bkgd fraction: 13%

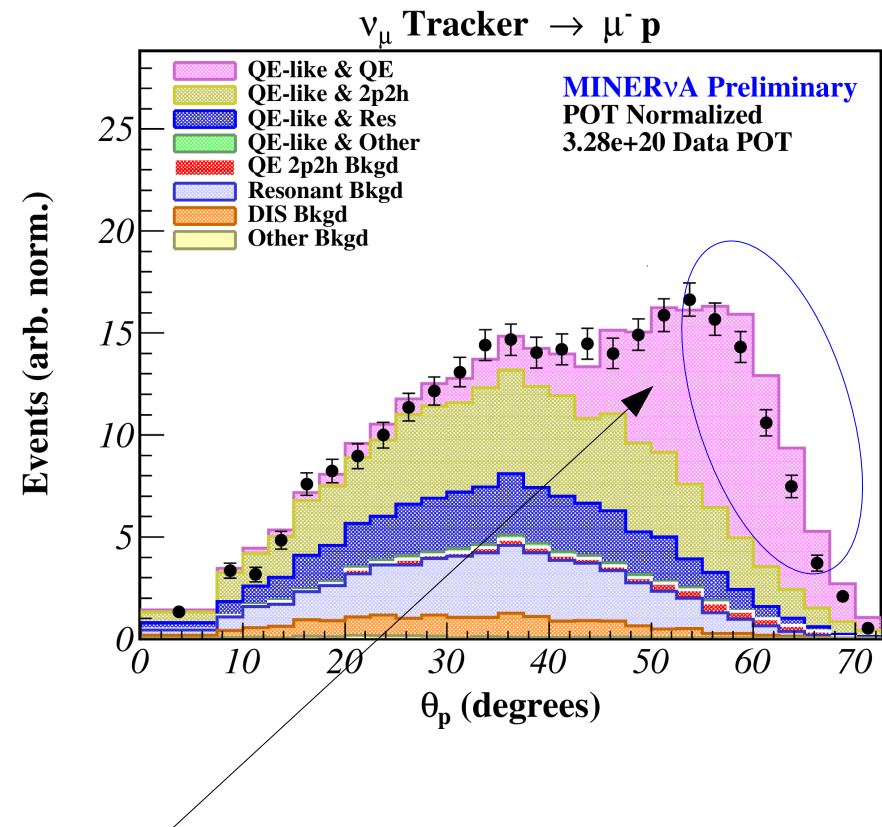
Overall DIS bkgd fraction: 5%

Overall Other bkgd fraction < 1%



Selected Sample

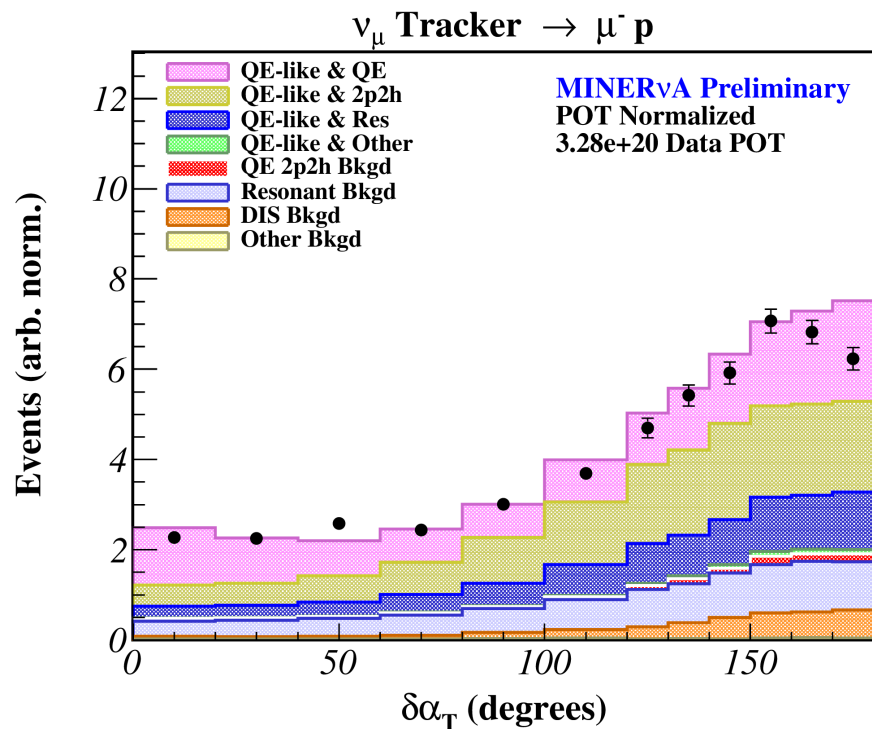
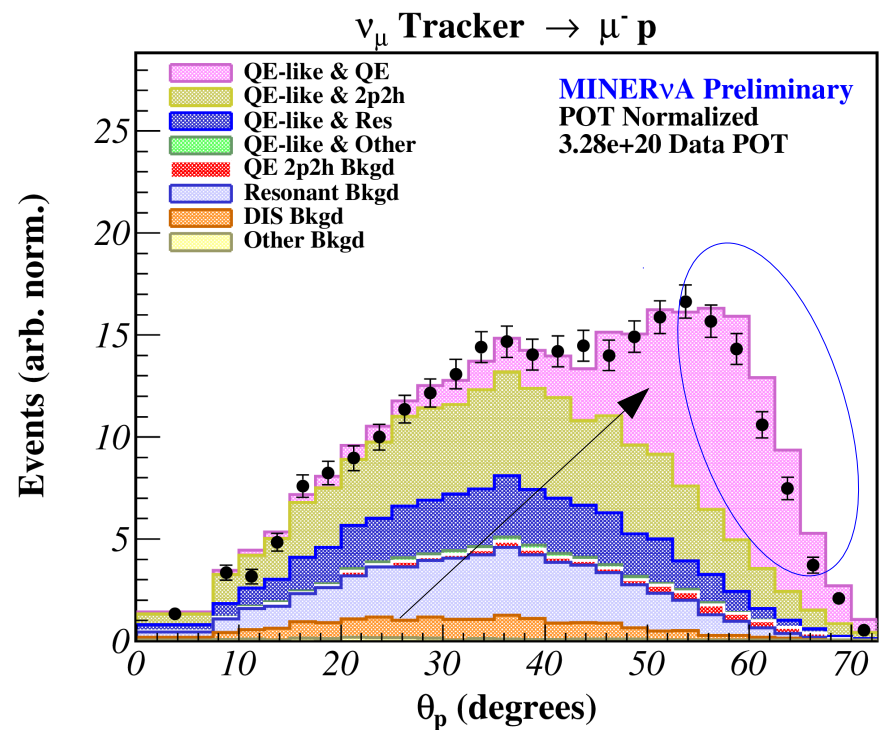
- Data-MC comparison at reconstructed level after sideband fit



GENIE MINERvA Tune (v1) describes data well (to first order)
Large concentration of pure QE at high angle
GENIE excess above data beyond 60 deg (see discussion later slides)

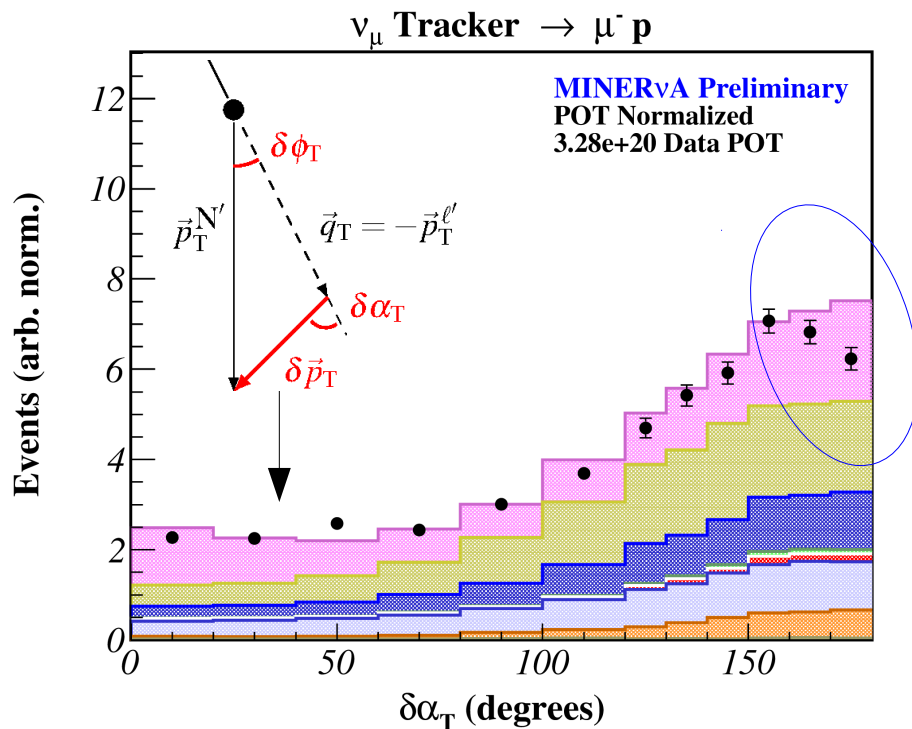
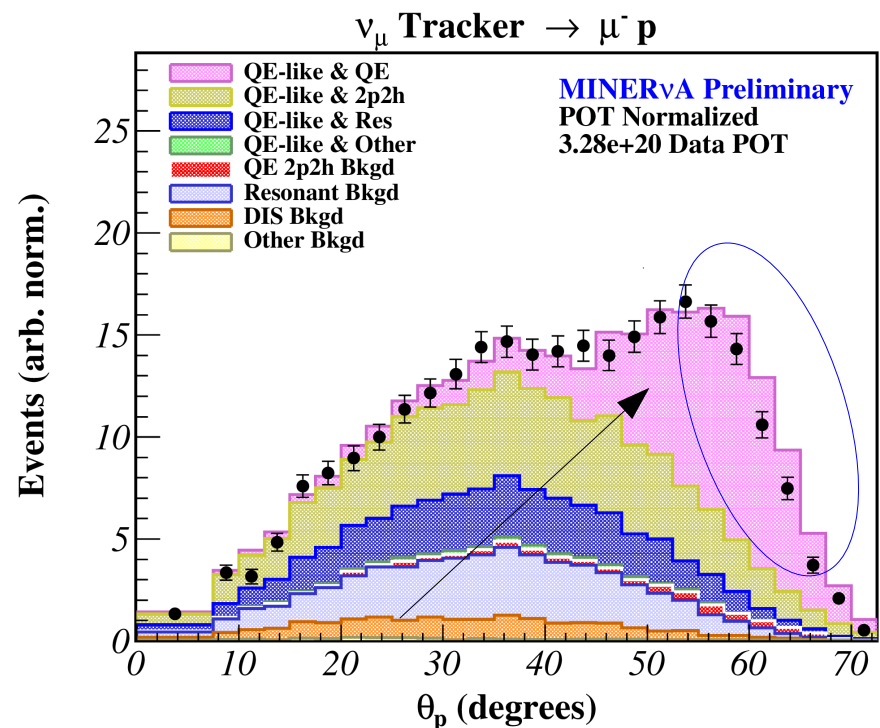
Selected Sample

- Data-MC comparison at reconstructed level after sideband fit



Selected Sample

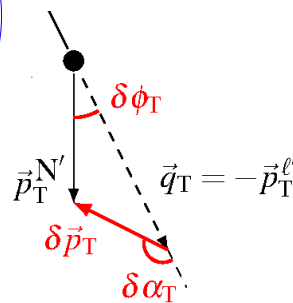
- Data-MC comparison at reconstructed level after sideband fit



GENIE MINERvA Tune (v1) describes data well (to first order)

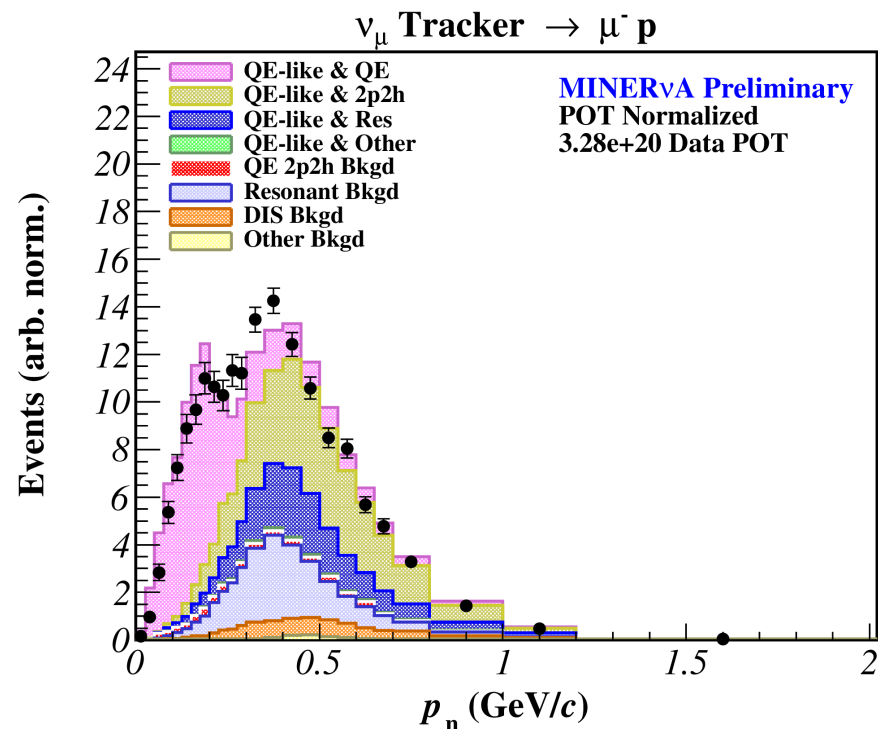
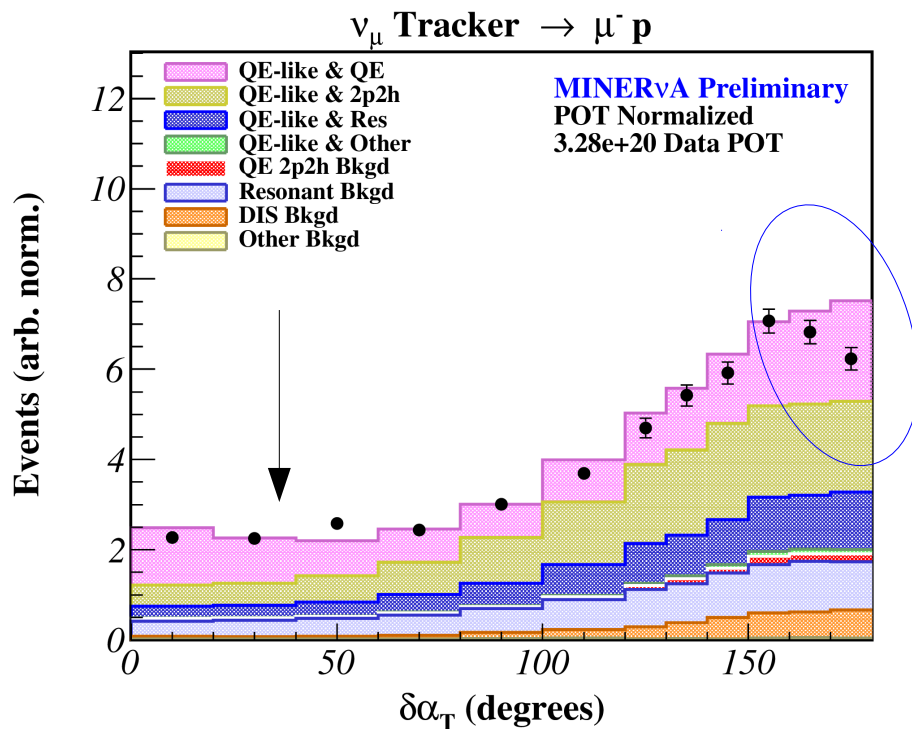
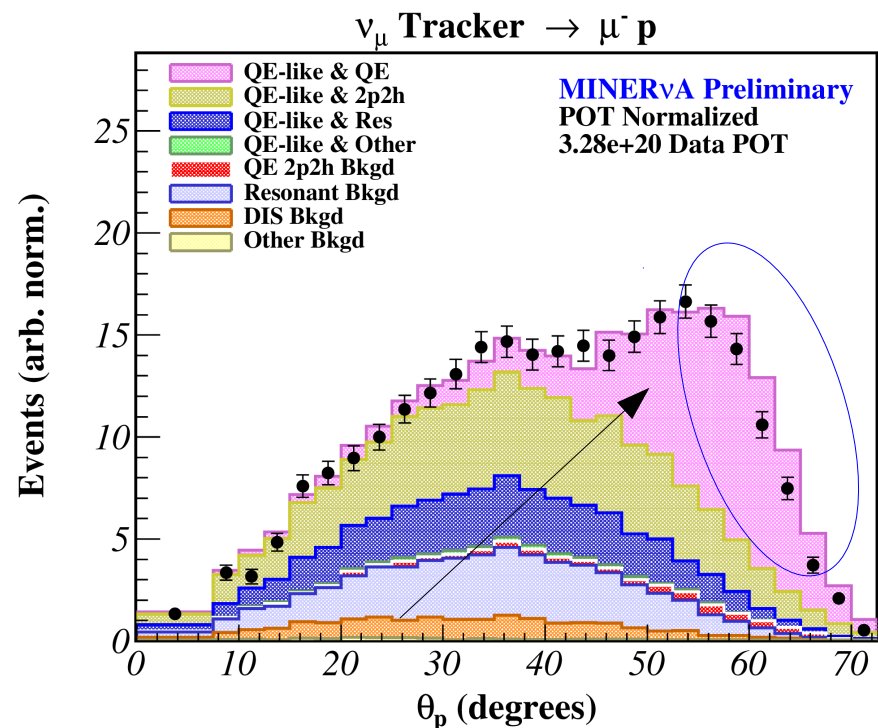
Depletion at small $\delta\alpha_T$

GENIE excess at $\delta\alpha_T \rightarrow 180$ deg.



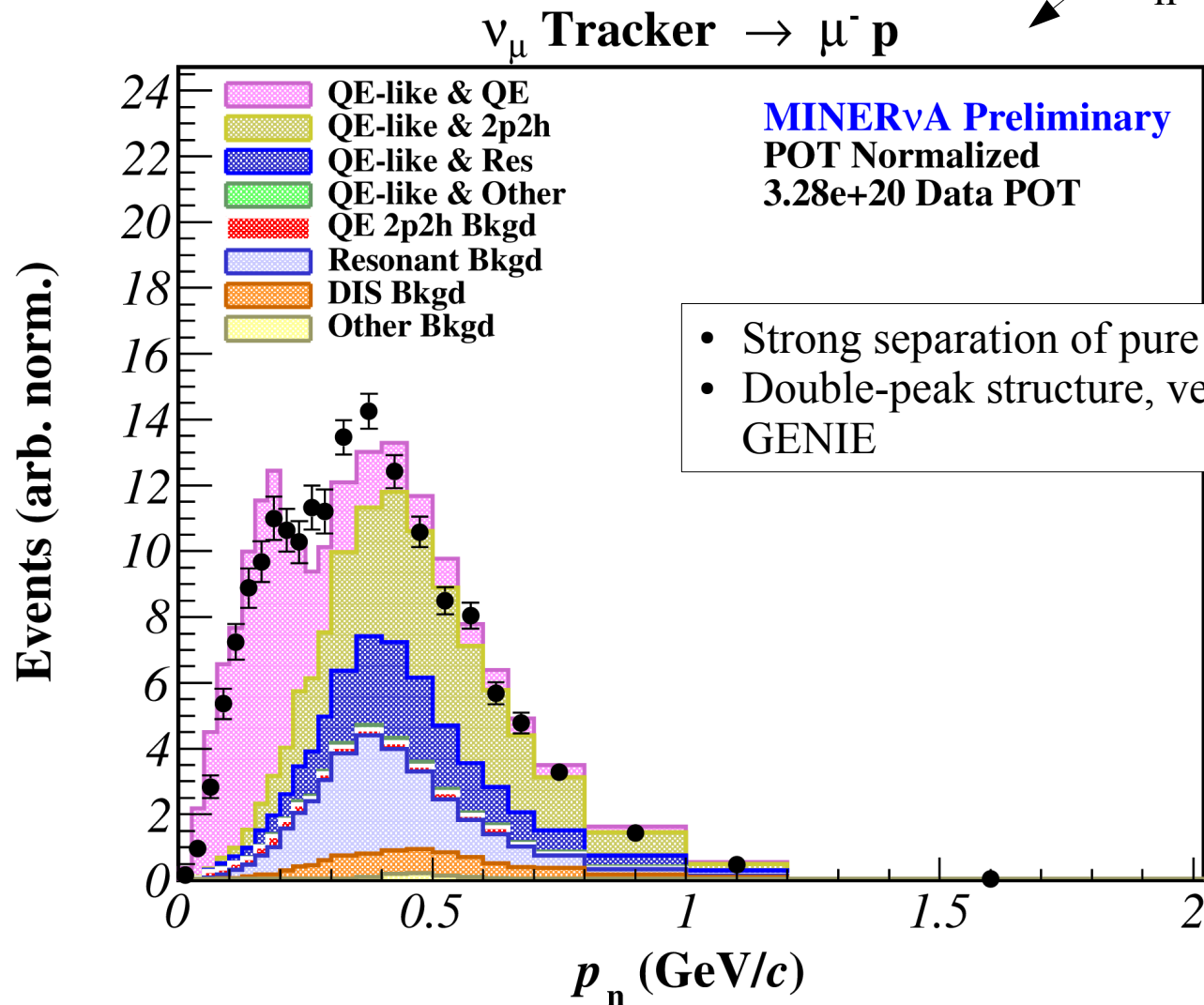
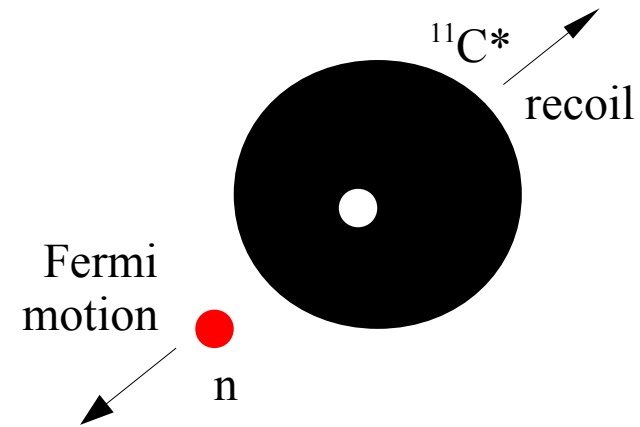
Selected Sample

- Data-MC comparison at reconstructed level after sideband fit

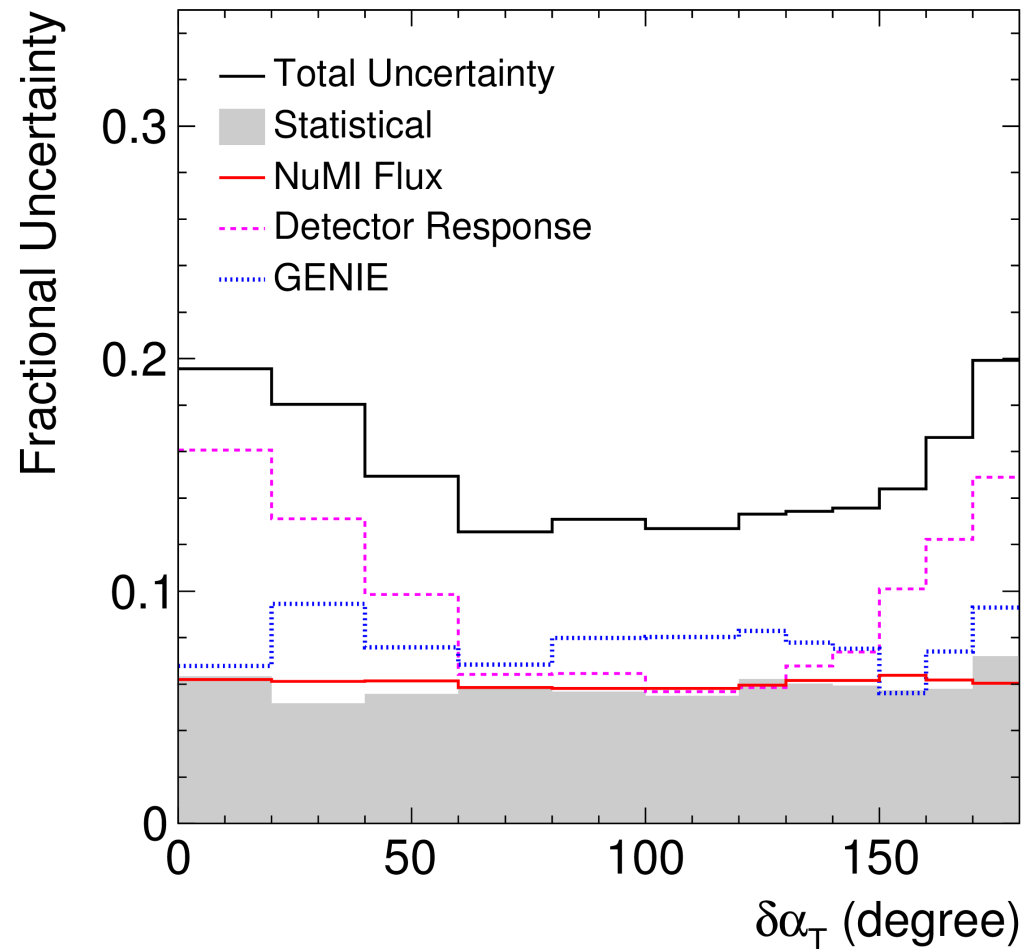


Selected Sample

- Data-MC comparison at reconstructed level after sideband fit



Systematic Uncertainties



Total uncertainty = 12-20%

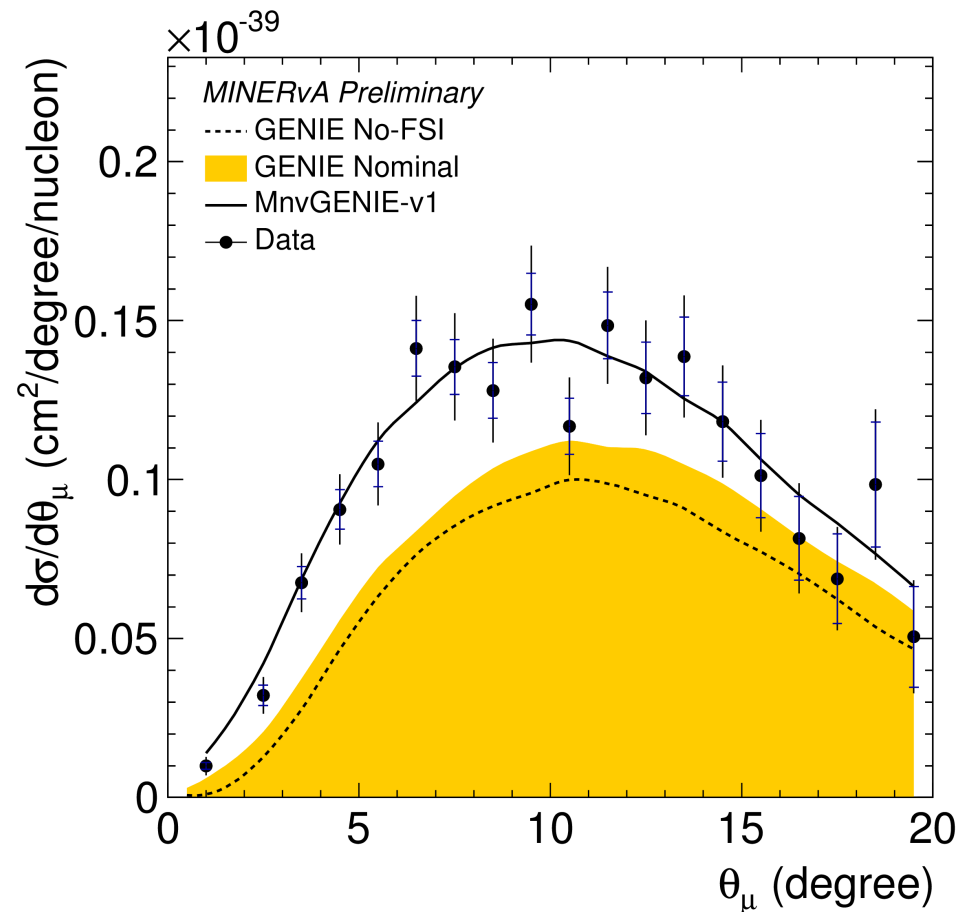
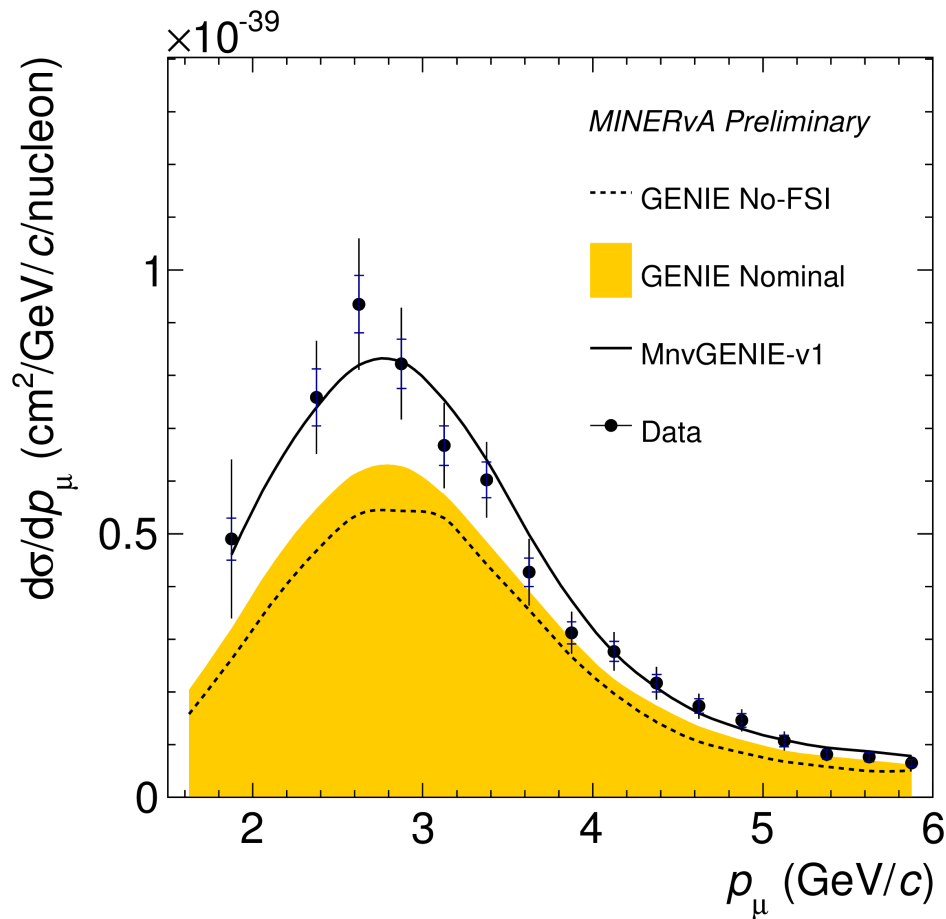
~ 6% (stat) + 6% (flux, mostly normalization) + 8% (GENIE) + 6-16% (detector)

Detector systematics dominated by transverse projection and ESC proton selection uncertainties.
GENIE systematics dominated by 2p2h model uncertainties.

RESULTS

Single-Particle Kinematics

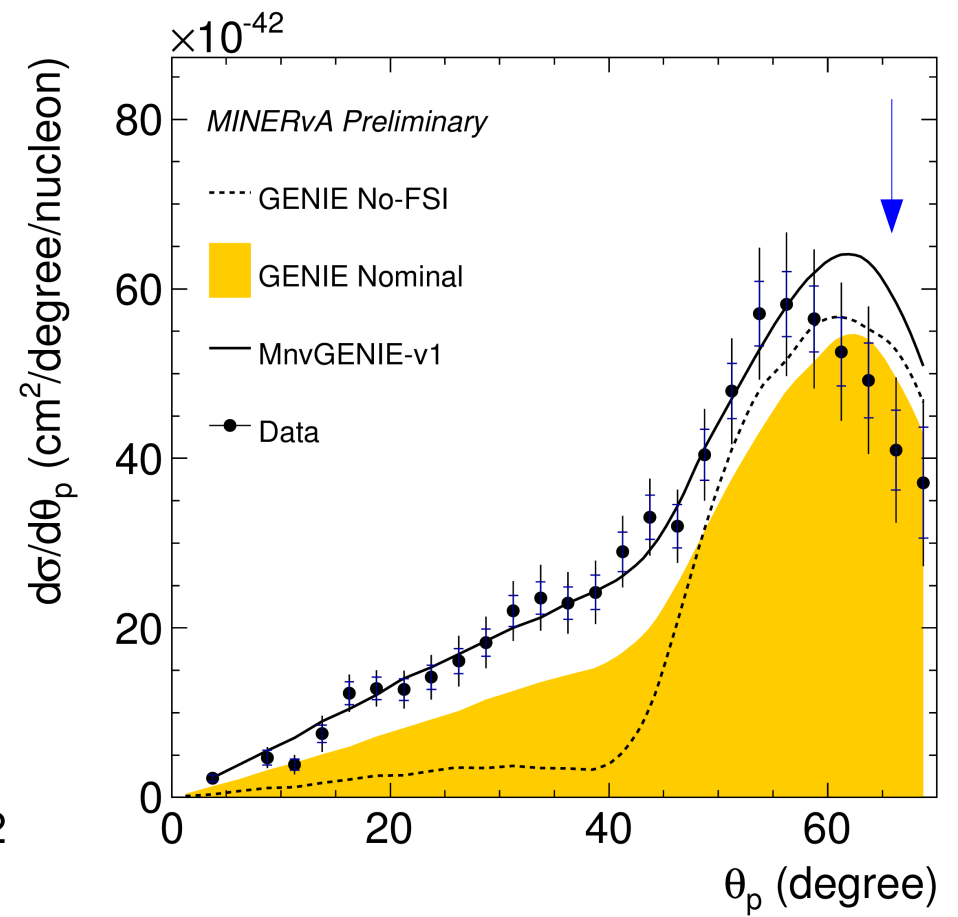
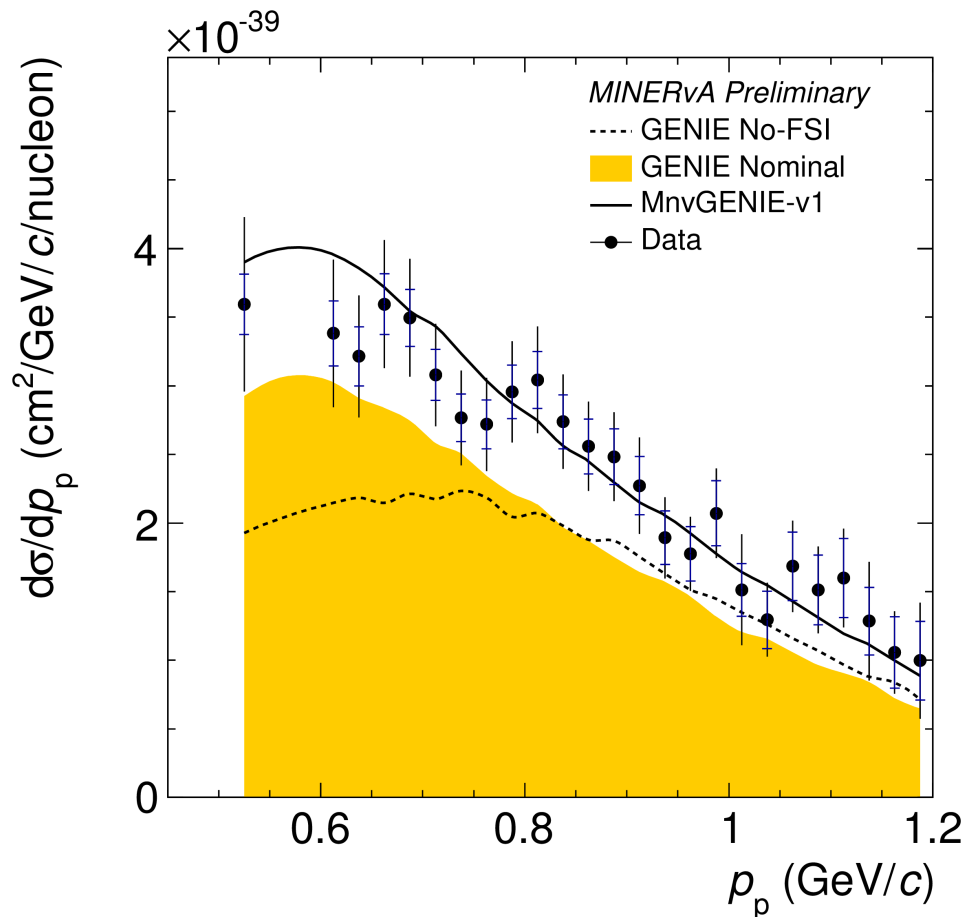
- Muon momentum, angle



- Good description by GENIE MINERvA Tune (v1)
- All predictions have same shape

Single-Particle Kinematics

- Muon momentum, angle
- Proton momentum, angle



- GENIE Nominal and No-FSI have different shape
- GENIE MINERvA Tune (v1) excess at high angle

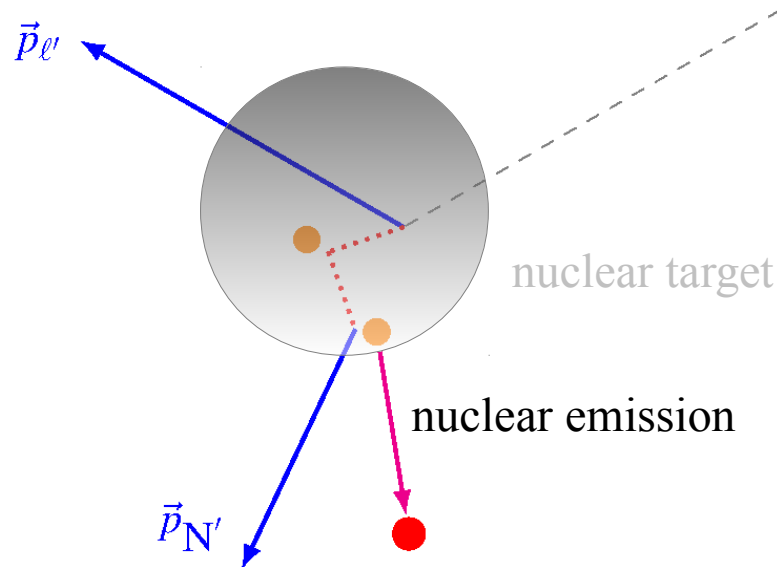
FSI decomposition in mesonless proton production:

Proton FSI:

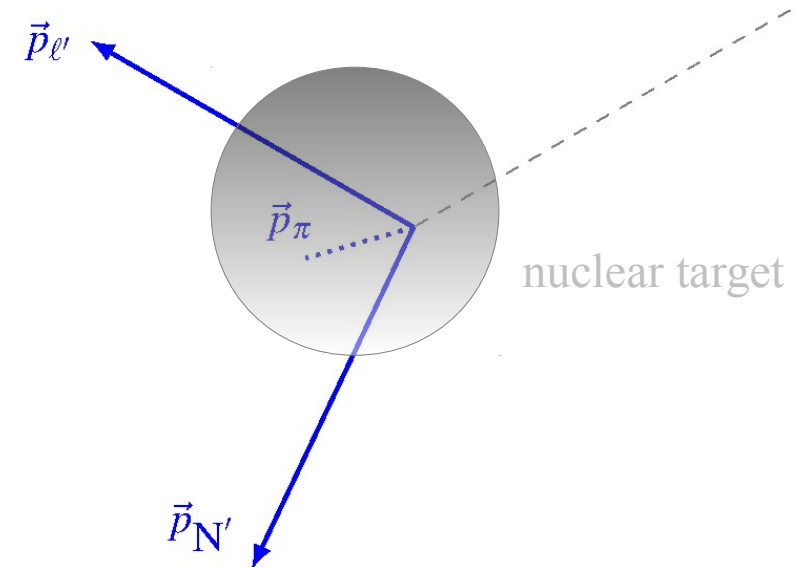
- Non-interacting (no change of energy and direction of the proton)
- **Acceleration**: energy of proton increased after FSI
- **Deceleration**: energy of proton decreased after FSI

Pion FSI: pion absorption

charged current (CC) $\nu \rightarrow l'$

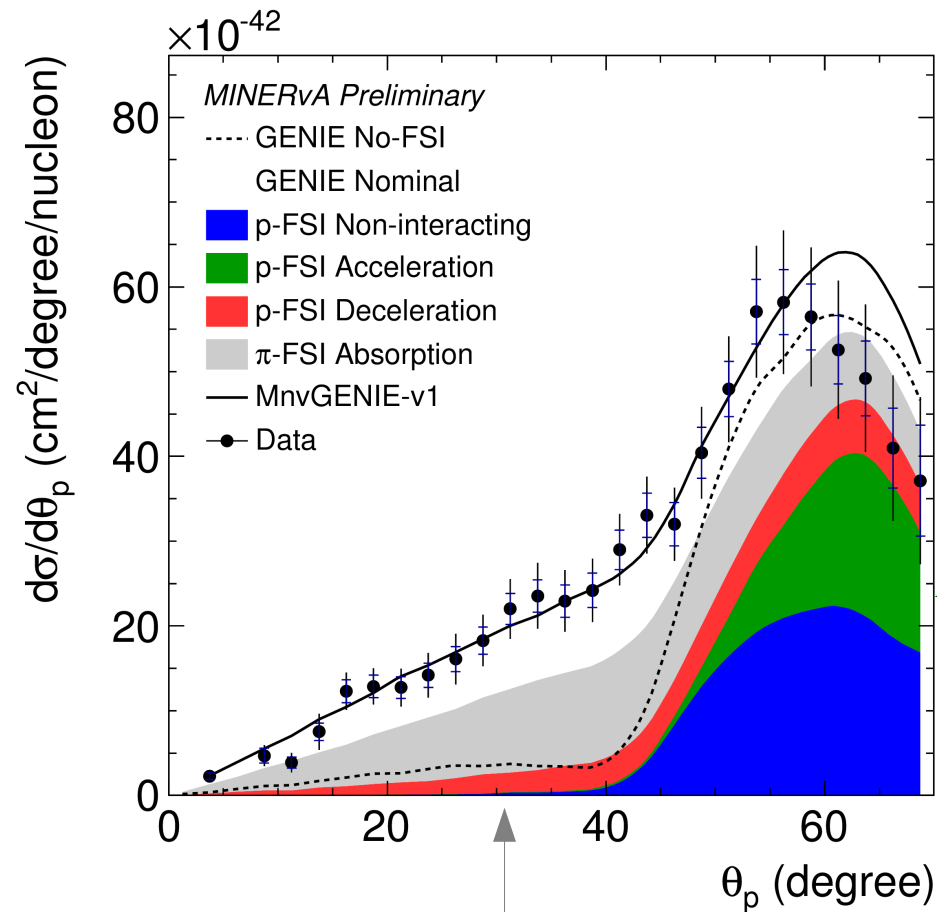


charged current (CC) $\nu \rightarrow l'$



Single-Particle Kinematics

- Muon momentum, angle
- Proton momentum, angle



Proton FSI acceleration
localized at high angle

Pionless resonant production dominates low angle

NUCLEAR EFFECT DIAGNOSTICS

A more general analysis of kinematic imbalance

Transverse: $0 = \vec{p}_T^{\ell'} + \vec{p}_T^{N'} - \delta\vec{p}_T$

Longitudinal: $E_\nu = p_L^{\ell'} + p_L^{N'} - \delta p_L$

New variable: $p_n \equiv \sqrt{\delta p_T^2 + \delta p_L^2}$

Neutrino energy is unknown (in the first place), equations are not closed.

For CCQE with elastic FSI, $A' = {}^{11}\text{C}^*$
 No more unknowns
 p_n : neutron Fermi motion *weakly smeared

A. Furmanski, J. Sobczyk, Phys.Rev. C95 (2017) no.6, 065501

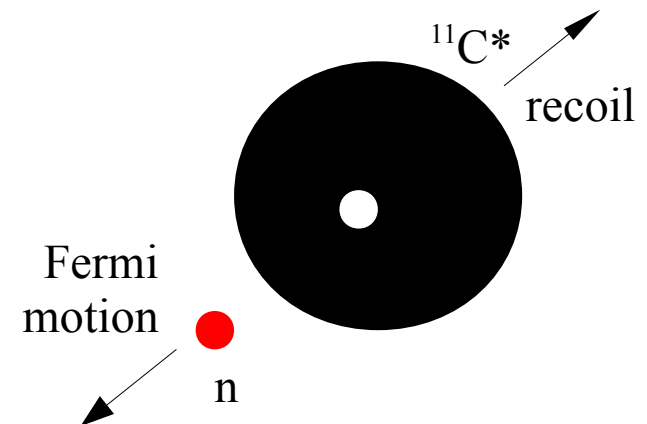
Xianguo Lu, Oxford

Assuming exclusive μ -p-A' final states
 Use energy conservation to close the equations

$$E_\nu + m_A = E_{\ell'} + E_{N'} + E_{A'}$$

$$E_{A'} = \sqrt{m_{A'}^2 + p_n^2}$$

p_n : recoil momentum of the nuclear remnant

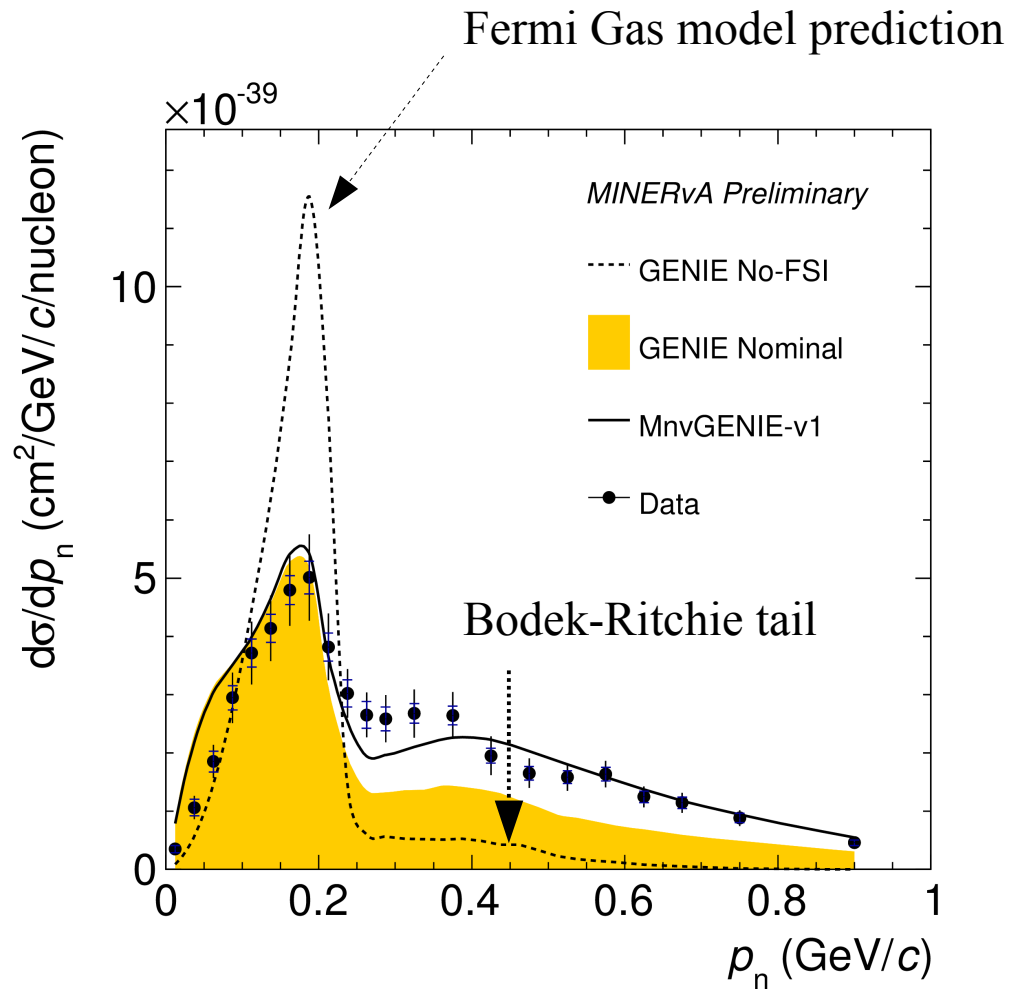


$$p_L = \frac{1}{2}(M(A) + k'_L + p'_L - E' - E_{p'})$$

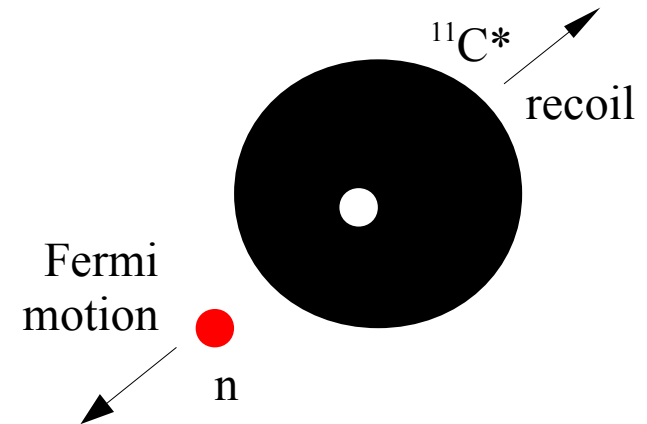
$$- \frac{p_T^2 + M^*(A-1)^2}{2(M(A) + k'_L + p'_L - E' - E_{p'})}$$

Nuclear Effect Diagnostics

- CCQE with Fermi motion

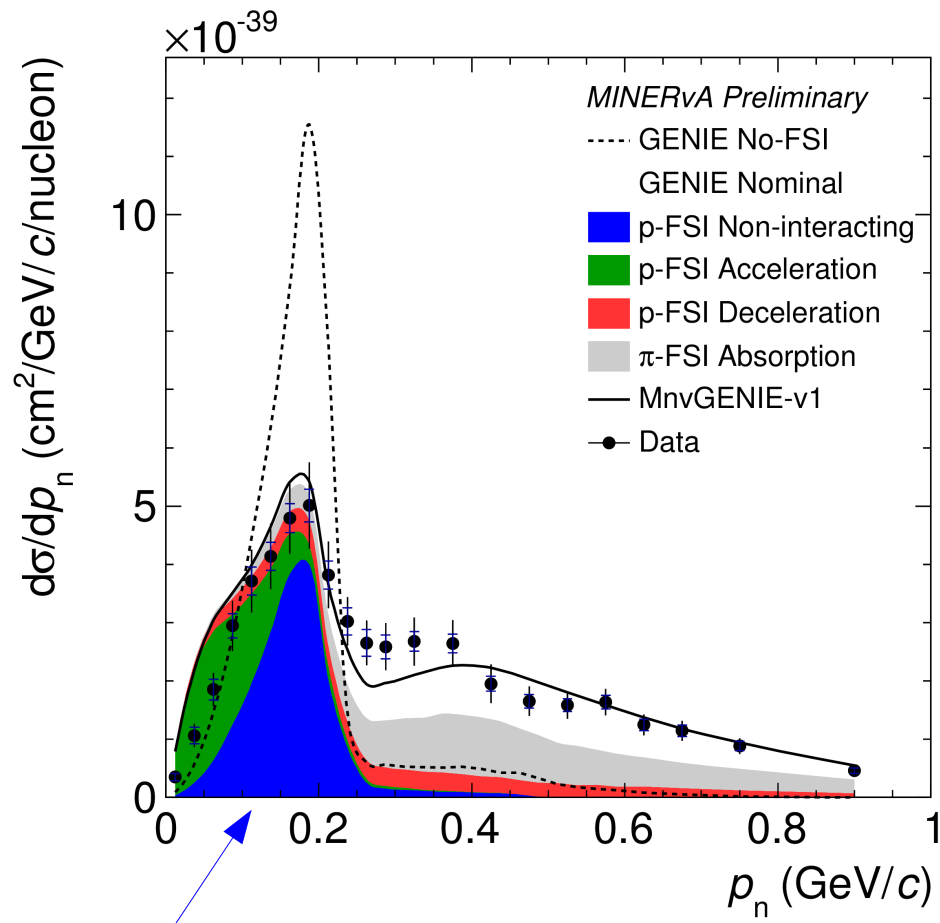


p_n is Fermi motion magnitude \rightarrow “QE peak”
– GENIE No-FSI



Nuclear Effect Diagnostics

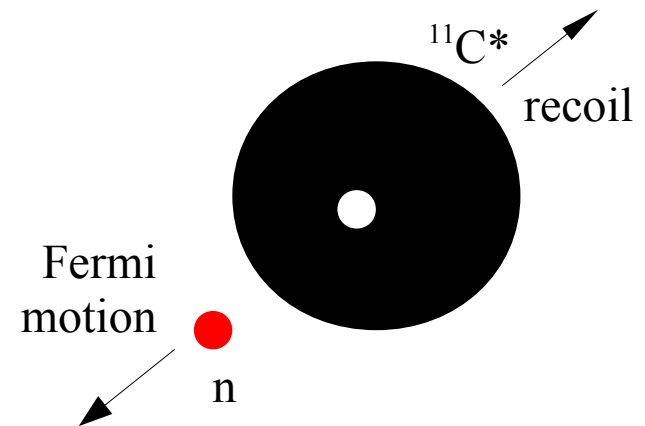
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p_n is Fermi motion magnitude \rightarrow “QE peak”

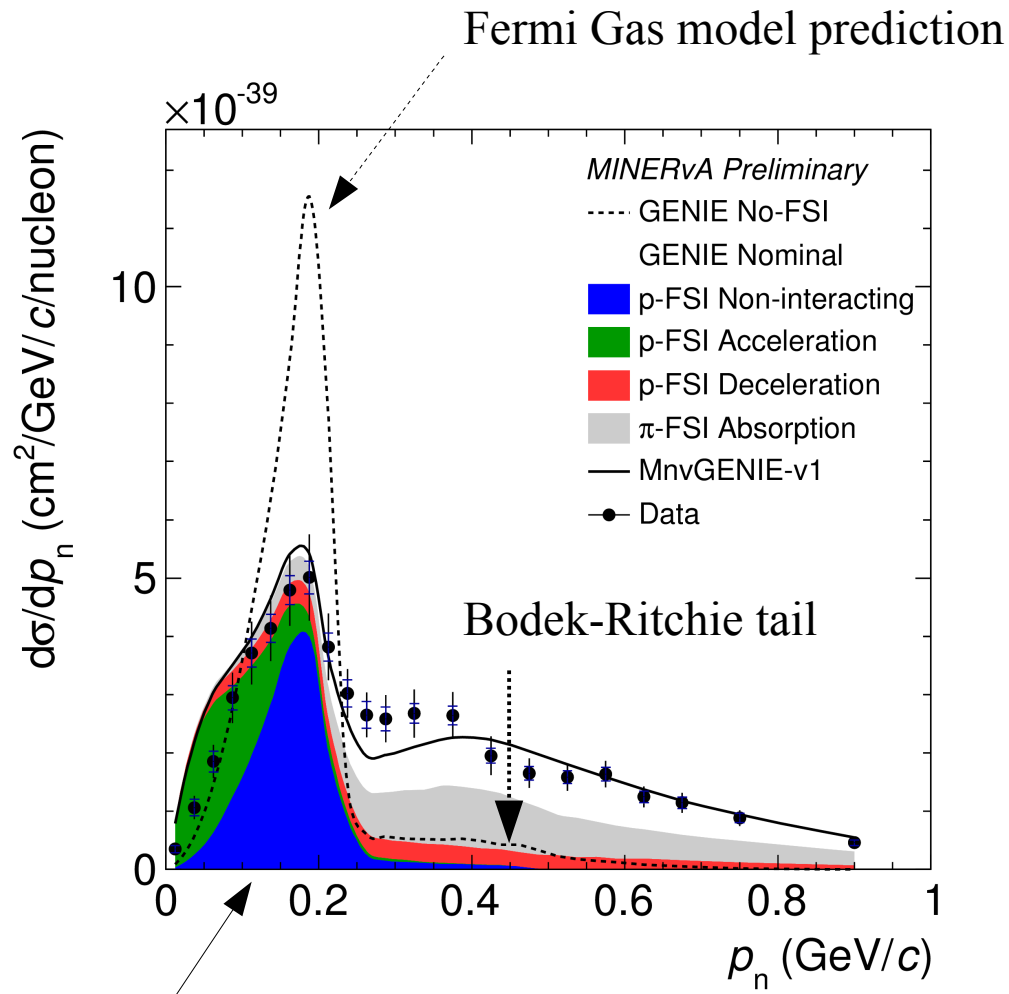
— GENIE No-FSI

— p-FSI Non-interacting



Nuclear Effect Diagnostics

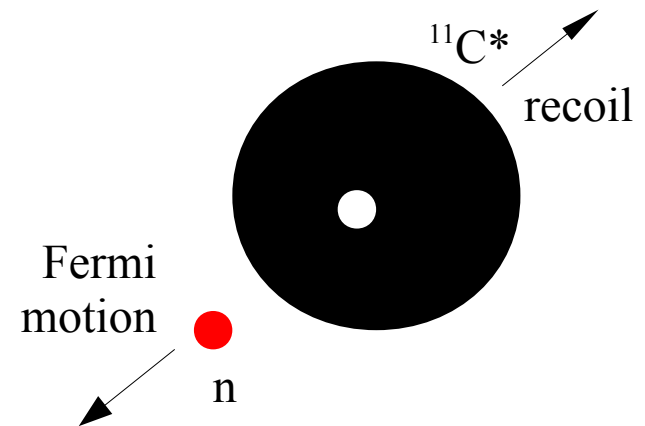
- CCQE with Fermi motion

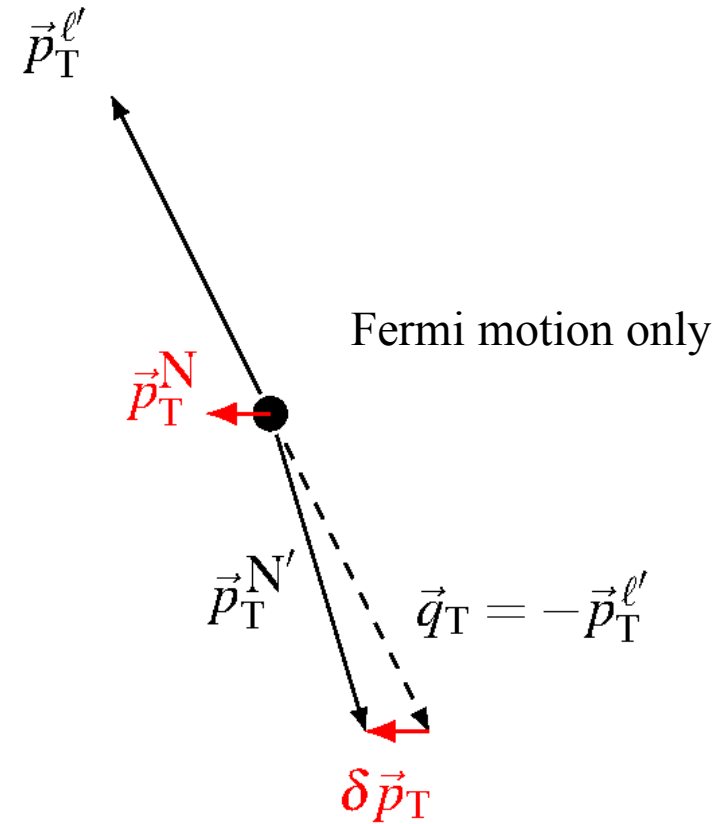


p_n is Fermi motion magnitude → “QE peak”

- GENIE No-FSI
- p-FSI Non-interacting

QE peak dominated by CCQE without FSI
Direct constraint of Fermi motion





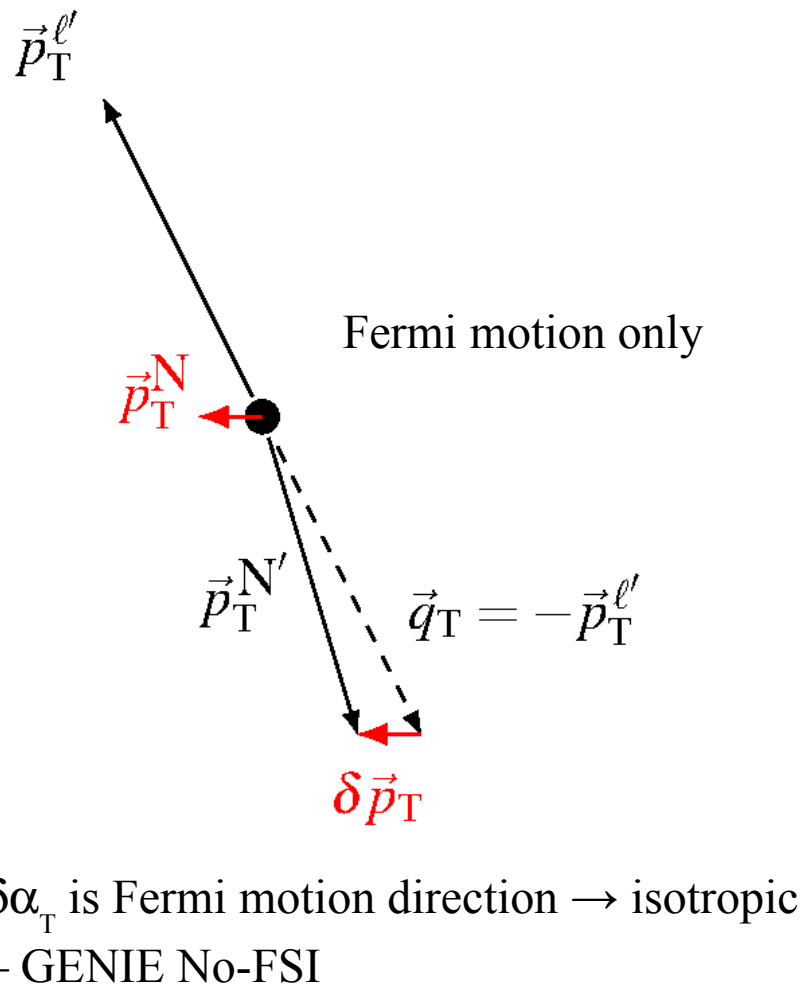
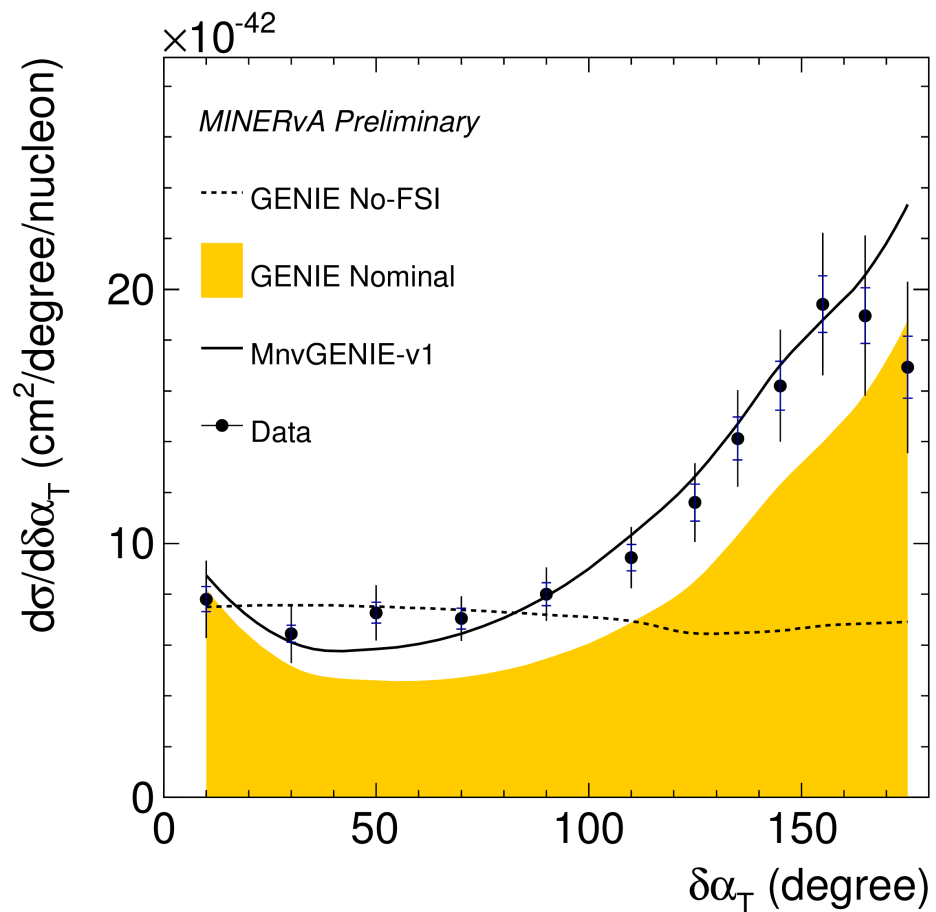
For CCQE with elastic FSI, $A' = {}^{11}\text{C}^*$
 No more unknowns
 \vec{p}_n : neutron Fermi motion *weakly smeared

$$\delta \vec{p}_T = \vec{p}_T^N$$

$\delta \alpha_T$ is Fermi motion direction \rightarrow isotropic

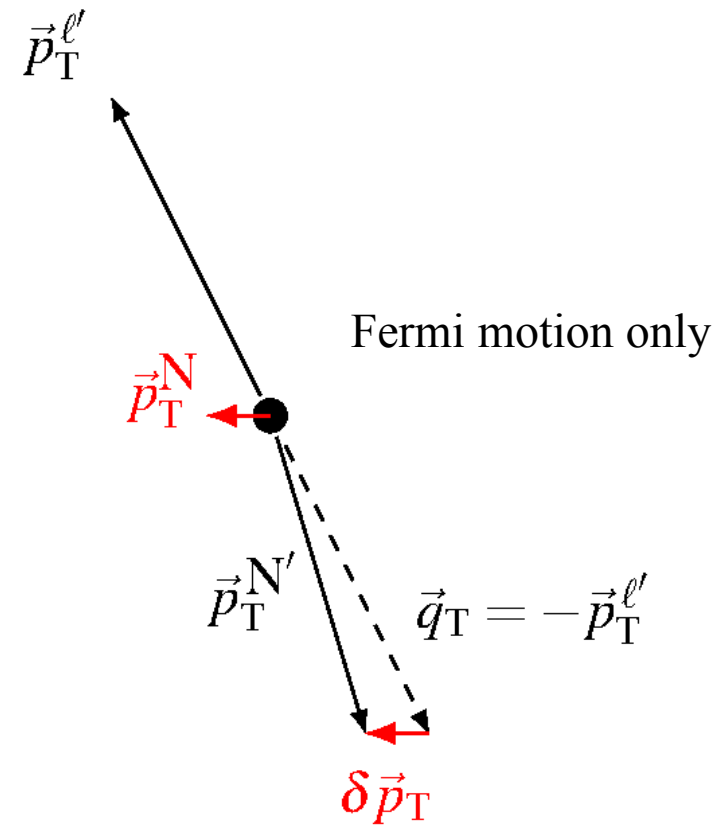
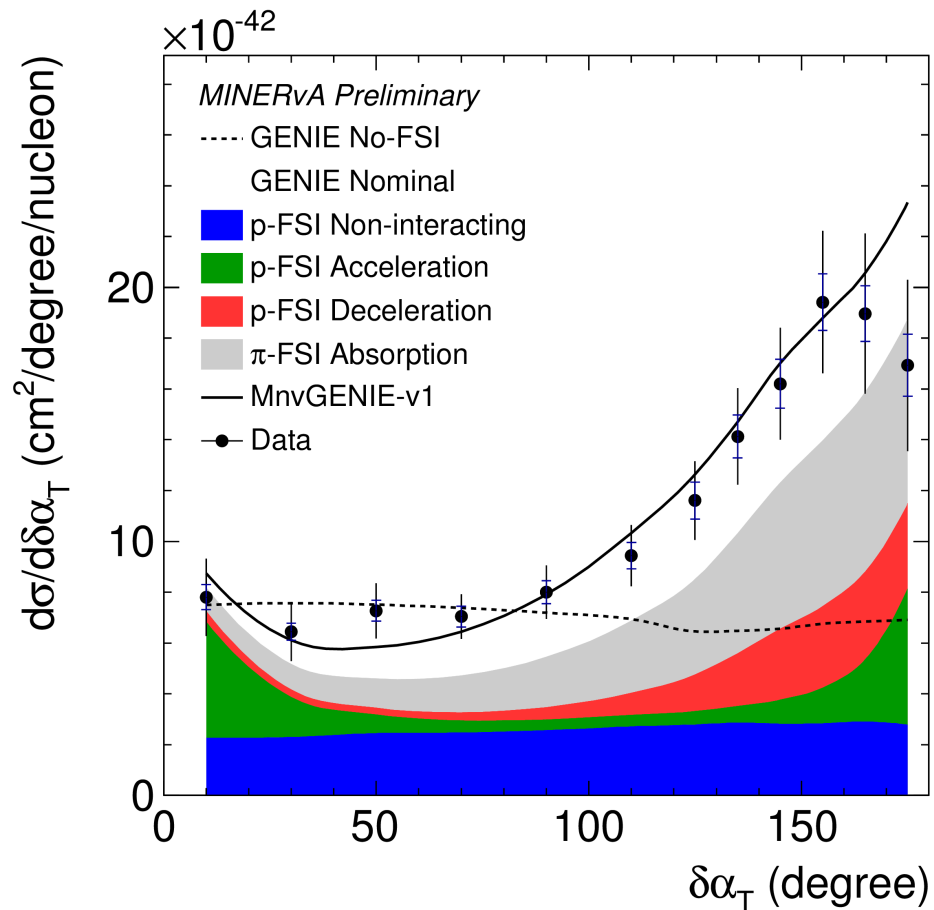
Nuclear Effect Diagnostics

- CCQE with Fermi motion



Nuclear Effect Diagnostics

- CCQE with Fermi motion



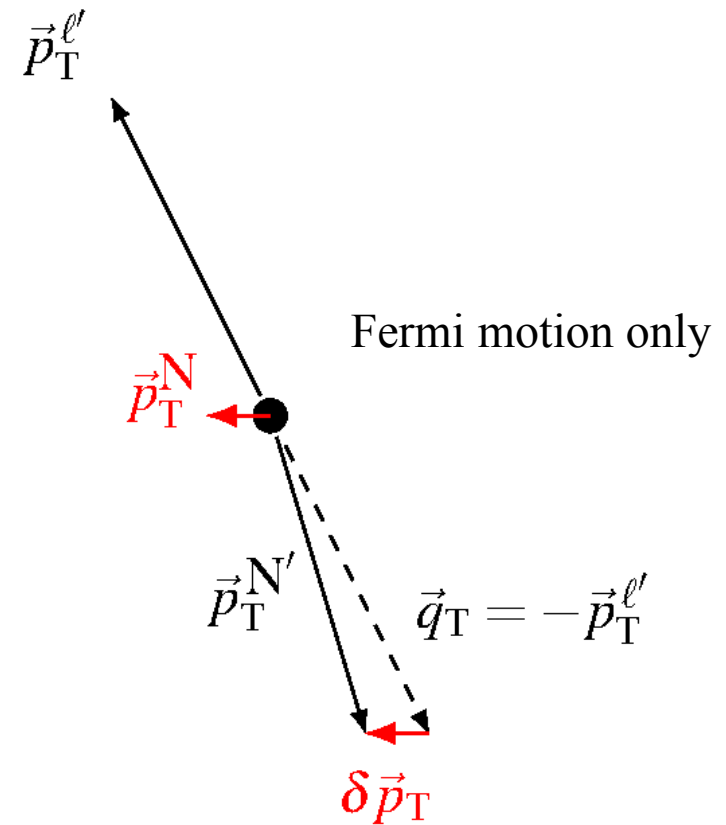
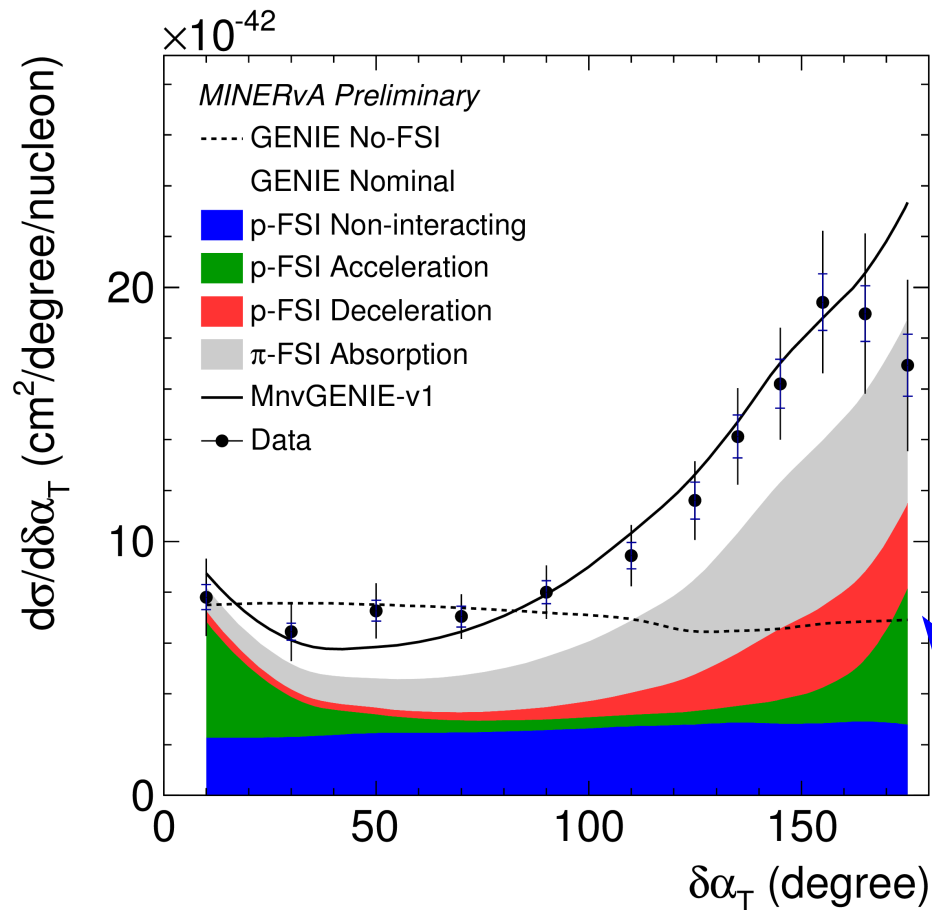
$\delta\alpha_T$ is Fermi motion direction \rightarrow isotropic

– GENIE No-FSI

– p-FSI Non-interacting

Nuclear Effect Diagnostics

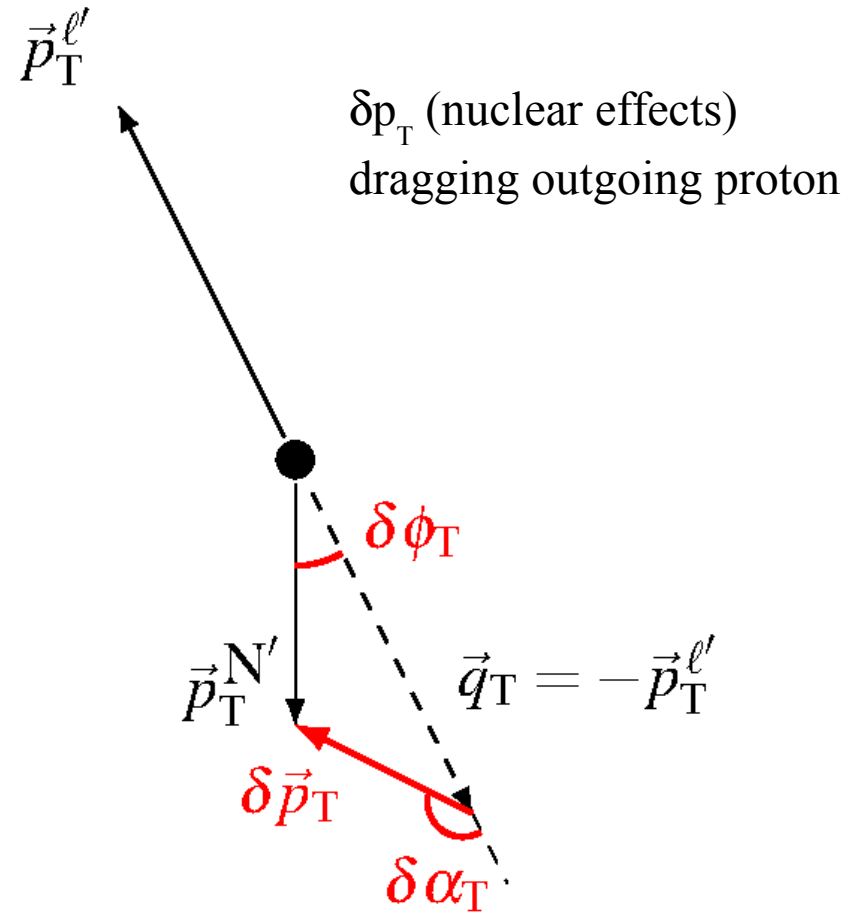
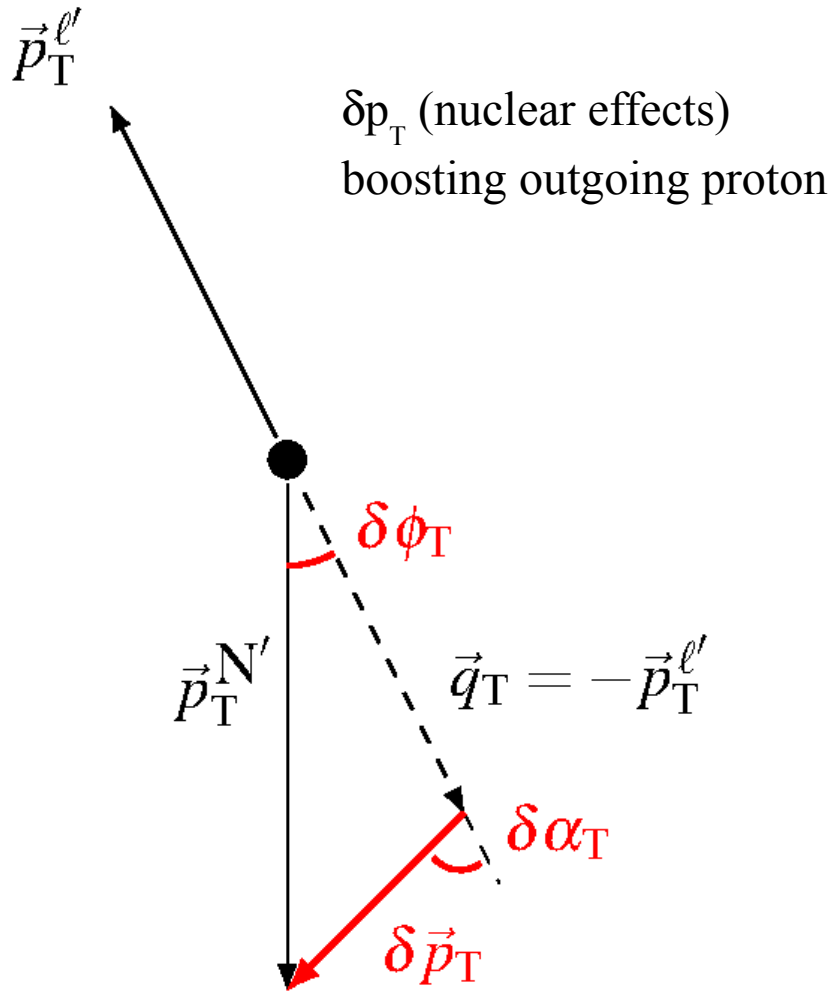
- CCQE with Fermi motion



$\delta\alpha_T$ is Fermi motion direction → isotropic

- GENIE No-FSI
- p-FSI Non-interacting

Baseline for all non-Fermi motion effects
 Factor out Fermi motion uncertainty
 Complementary to p_n



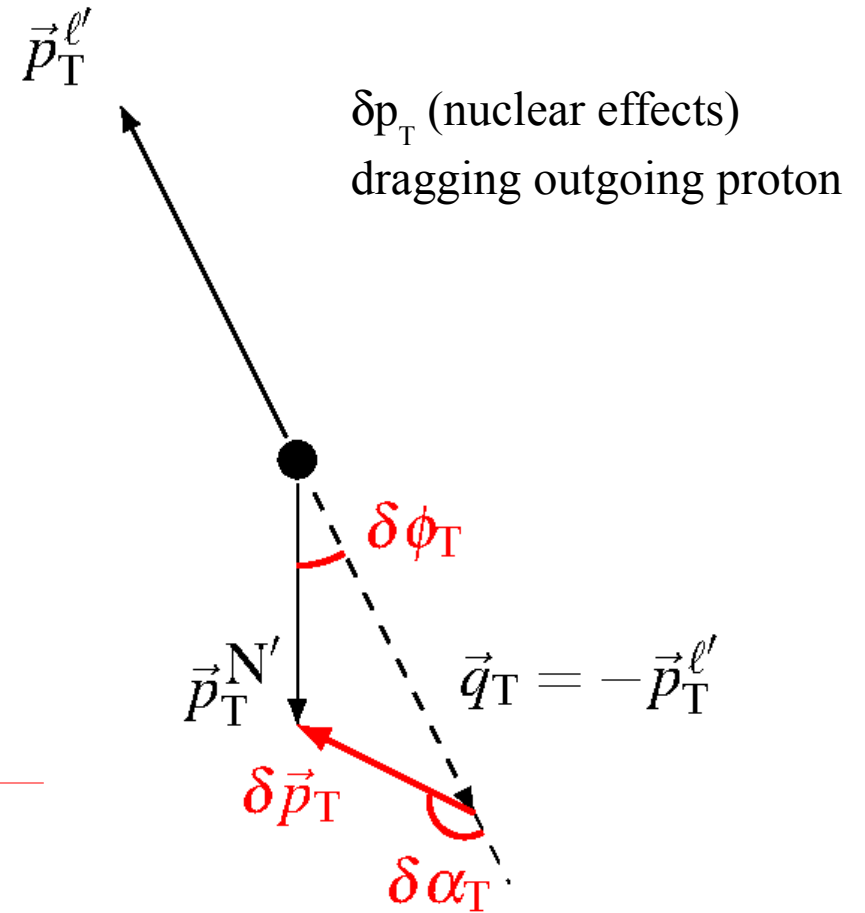
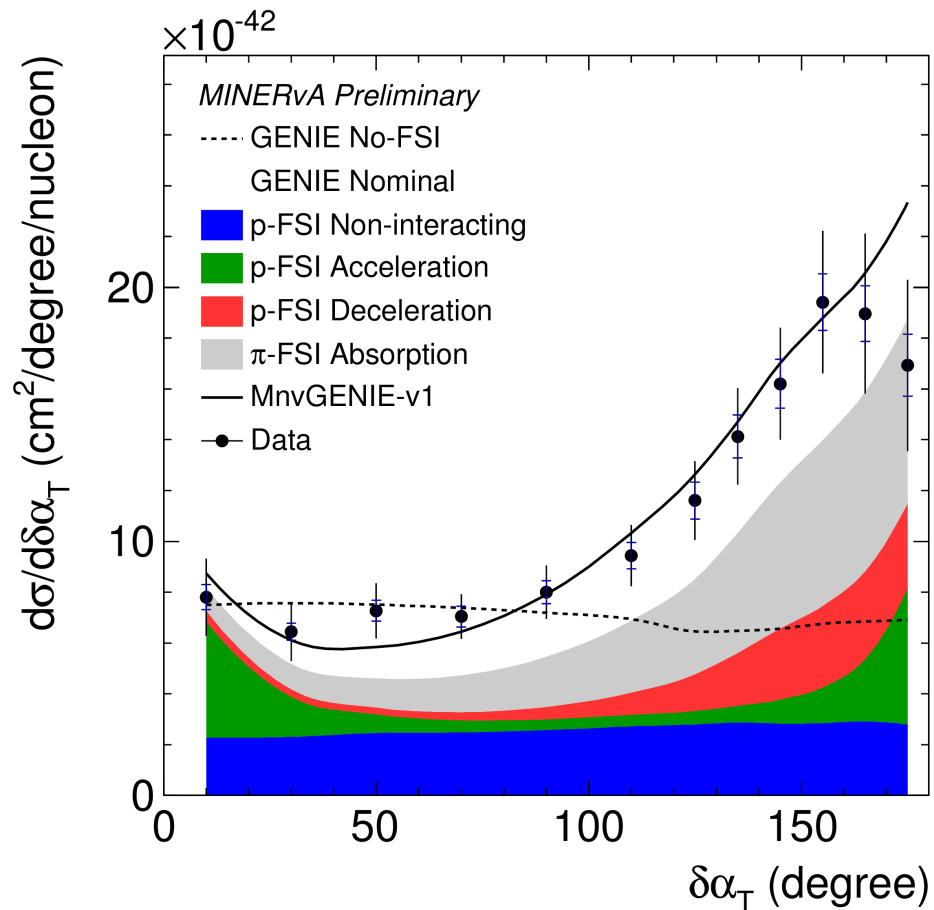
With full nuclear effects

$$\delta\vec{p}_T = \vec{p}_T^N - \Delta\vec{p}_T$$

Baseline for all non-Fermi motion effects
 Factor out Fermi motion uncertainty
 Complementary to p_n

Nuclear Effect Diagnostics

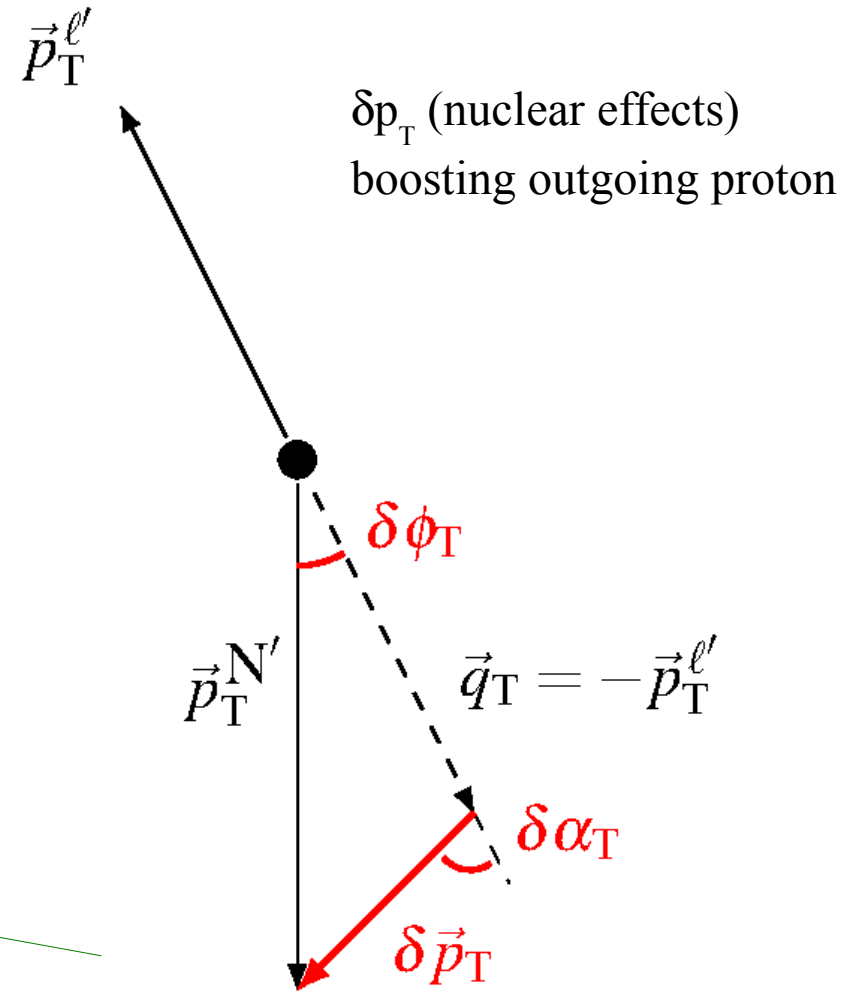
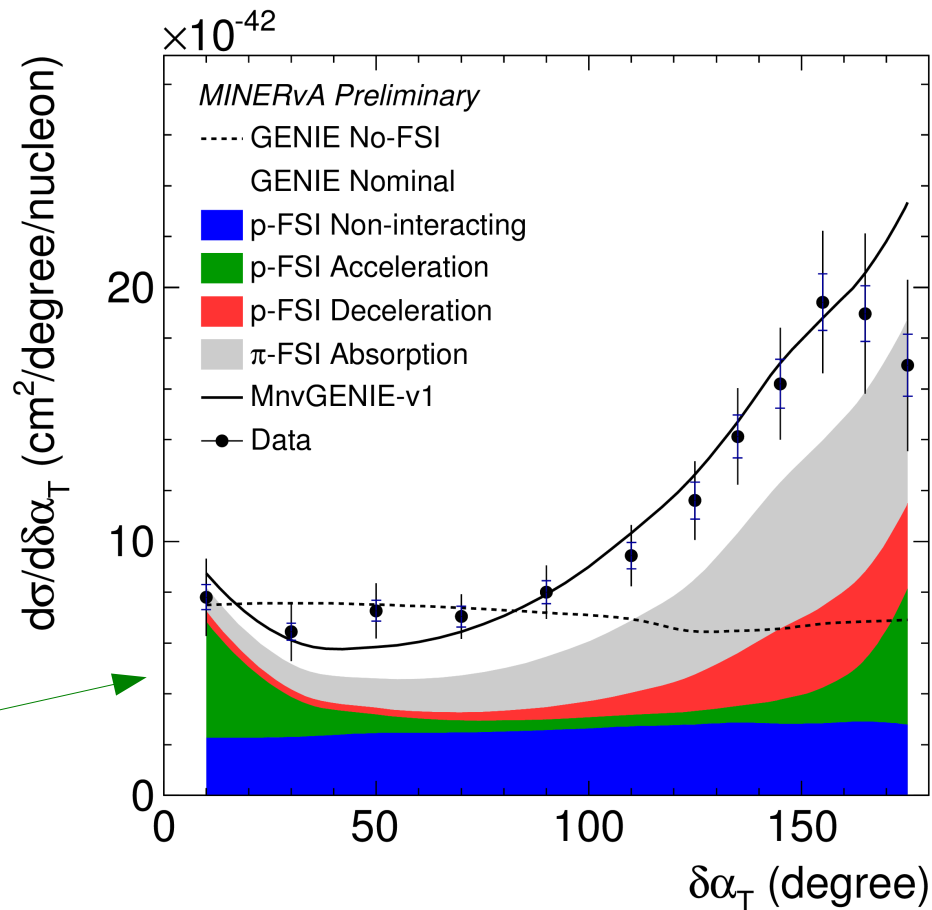
- CCQE with Fermi motion
- FSI deceleration vs. acceleration



Deceleration at large $\delta\alpha_T$

Nuclear Effect Diagnostics

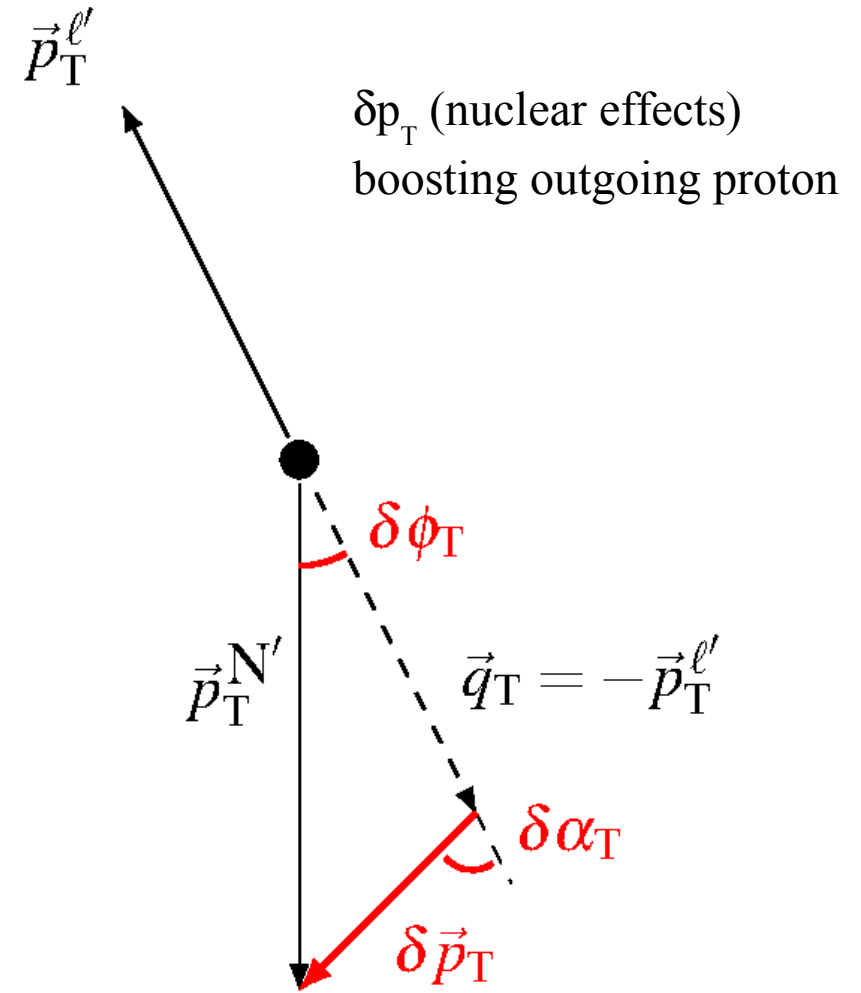
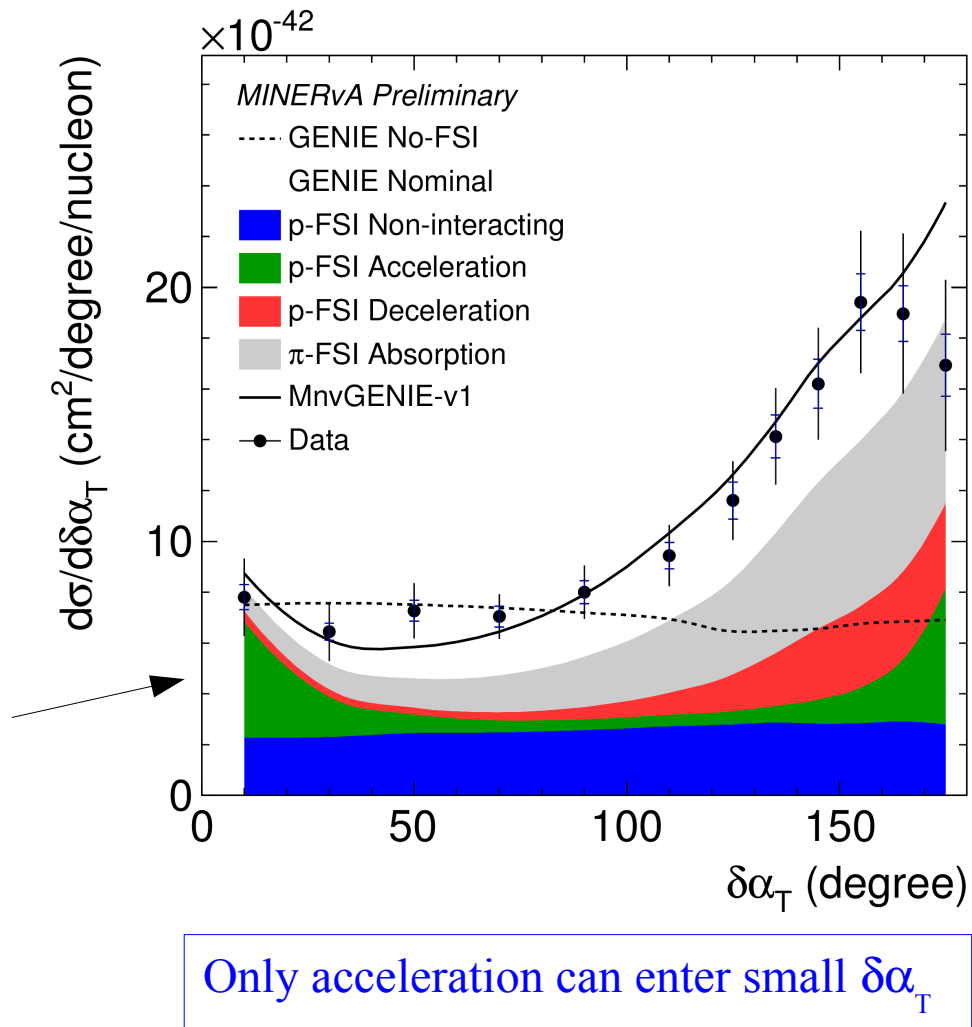
- CCQE with Fermi motion
- FSI deceleration vs. acceleration



Deceleration at large $\delta\alpha_T$
 Acceleration at both small and (due to transverse projection) large $\delta\alpha_T$

Nuclear Effect Diagnostics

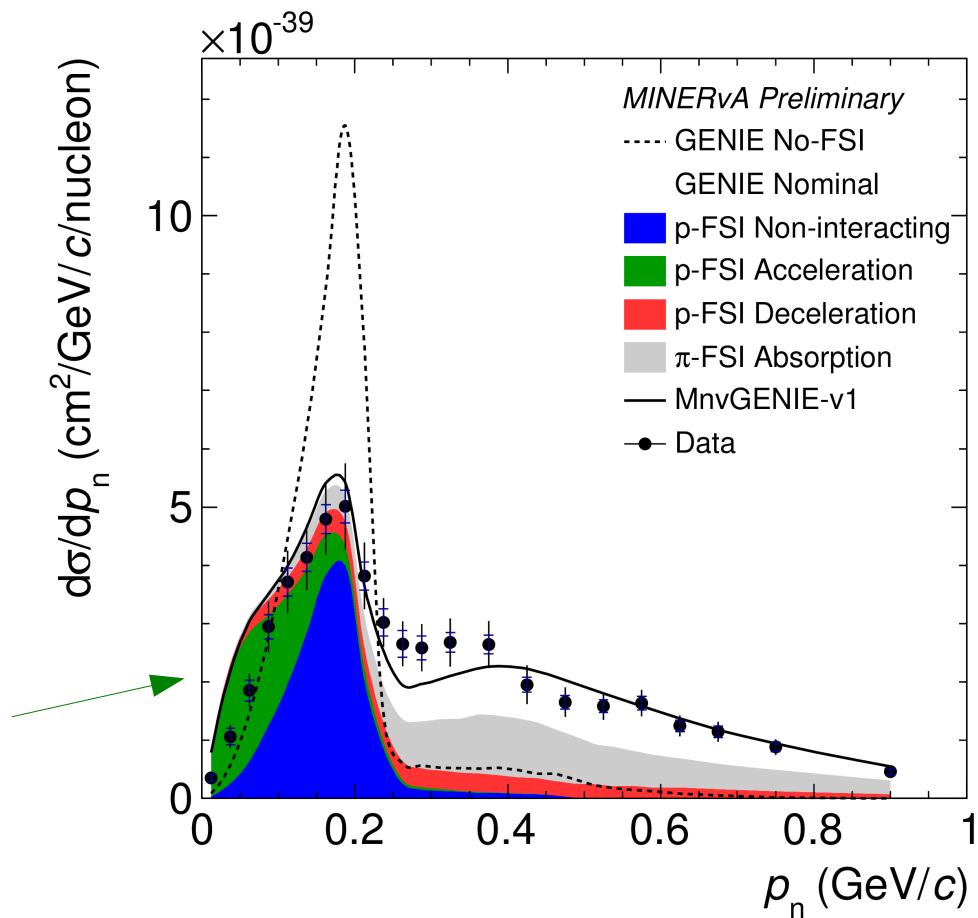
- CCQE with Fermi motion
- FSI deceleration vs. acceleration



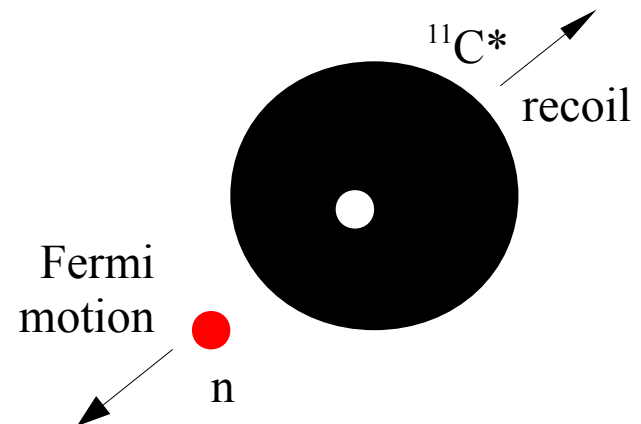
Deceleration at large $\delta\alpha_T$
Acceleration at both small and (due to transverse projection) large $\delta\alpha_T$

Nuclear Effect Diagnostics

- CCQE with Fermi motion
- FSI deceleration vs. acceleration



Acceleration to the left of QE peak
Strongly distort QE peak



A more general analysis of kinematic imbalance

Transverse: $0 = \vec{p}_T^{\ell'} + \vec{p}_T^{N'} - \delta\vec{p}_T$

Longitudinal: $E_\nu = p_L^{\ell'} + p_L^{N'} - \delta p_L$

New variable: $p_n \equiv \sqrt{\delta p_T^2 + \delta p_L^2}$

Neutrino energy is unknown (in the first place), equations are not closed.

For RES, DIS, 2p2h, no longer exclusive μ -p-A' final states
 p_n : smeared δp_T beyond QE peak

A. Furmanski, J. Sobczyk, Phys.Rev. C95 (2017) no.6, 065501

Xianguo Lu, Oxford

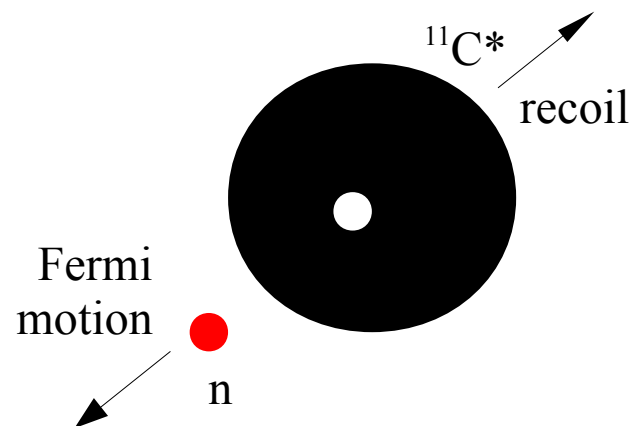
Assuming exclusive μ -p-A' final states

Use energy conservation to close the equations

$$E_\nu + m_A = E_{\ell'} + E_{N'} + E_{A'}$$

$$E_{A'} = \sqrt{m_{A'}^2 + p_n^2}$$

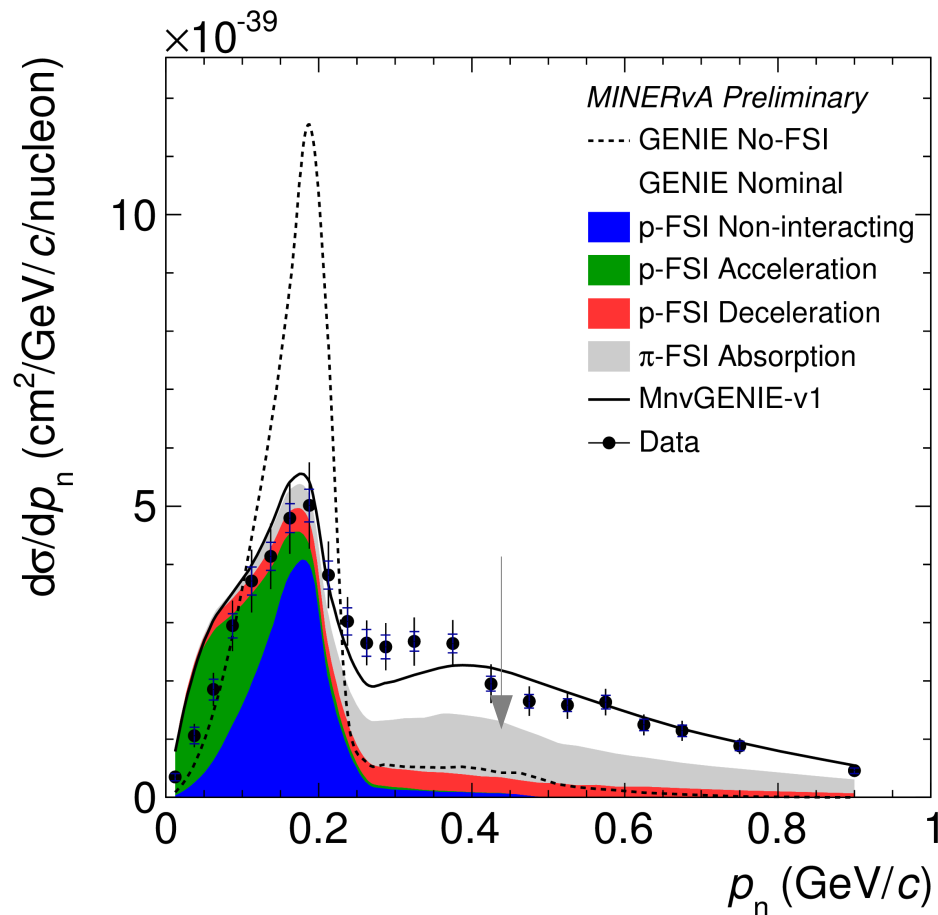
p_n : recoil momentum of the nuclear remnant



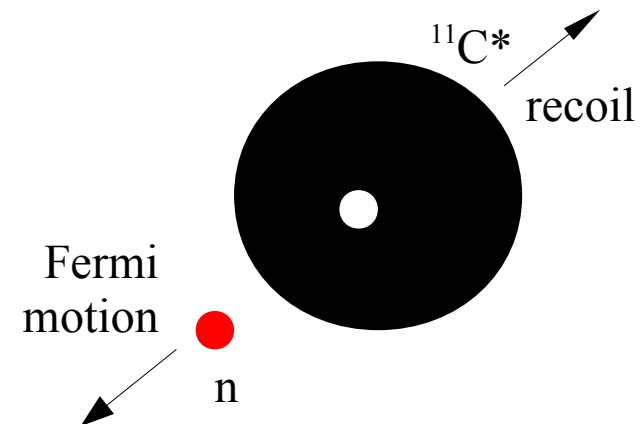
$$p_L = \frac{1}{2}(M(A) + k'_L + p'_L - E' - E_{p'}) - \frac{p_T^2 + M^*(A-1)^2}{2(M(A) + k'_L + p'_L - E' - E_{p'})}$$

Nuclear Effect Diagnostics

- CCQE with Fermi motion
- FSI deceleration vs. acceleration
- Pionless resonant production, pion absorption FSI, and 2p2h



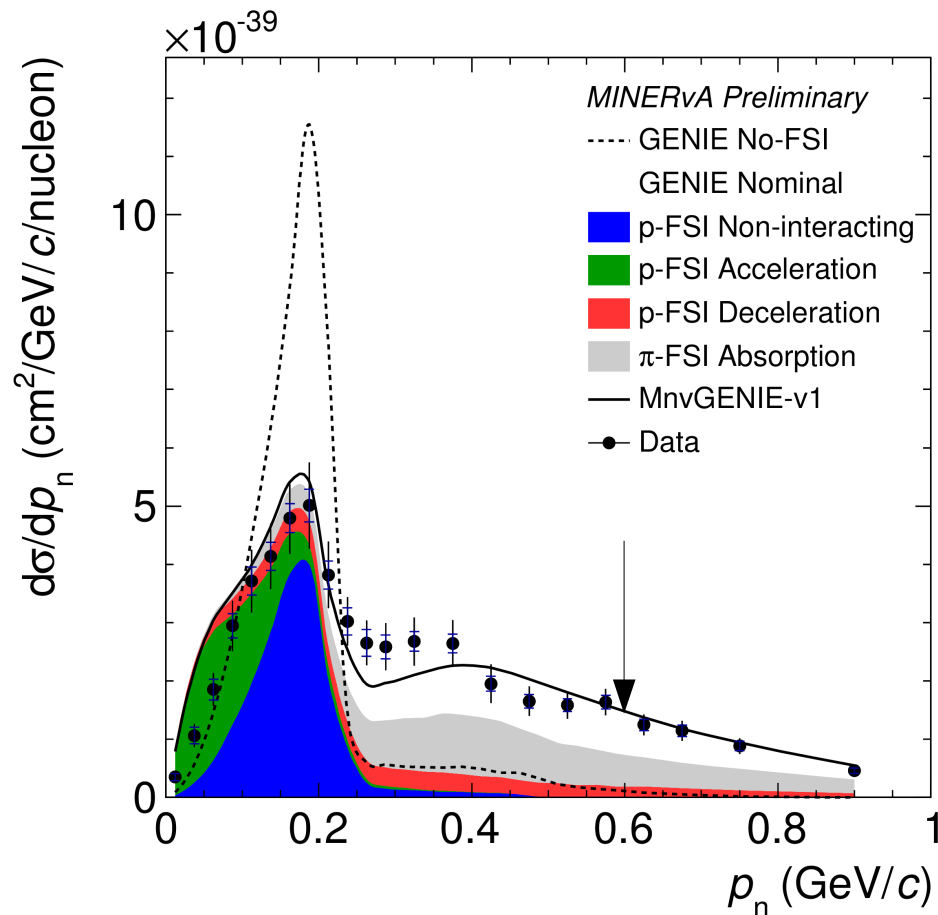
p_n : smeared δp_T beyond QE peak \rightarrow tail
 – π -FSI Absorption



Pion production and 2p2h process:
 strong intra-nuclear momentum transfer
 due to momentum sharing with proton

Nuclear Effect Diagnostics

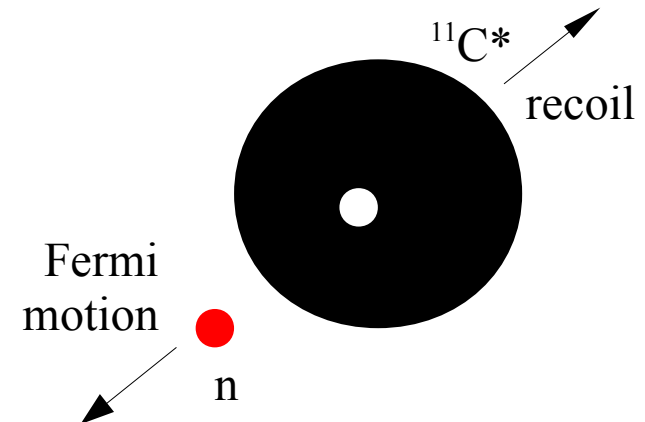
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p_n : smeared δp_T beyond QE peak \rightarrow tail

- π-FSI Absorption
- 2p2h

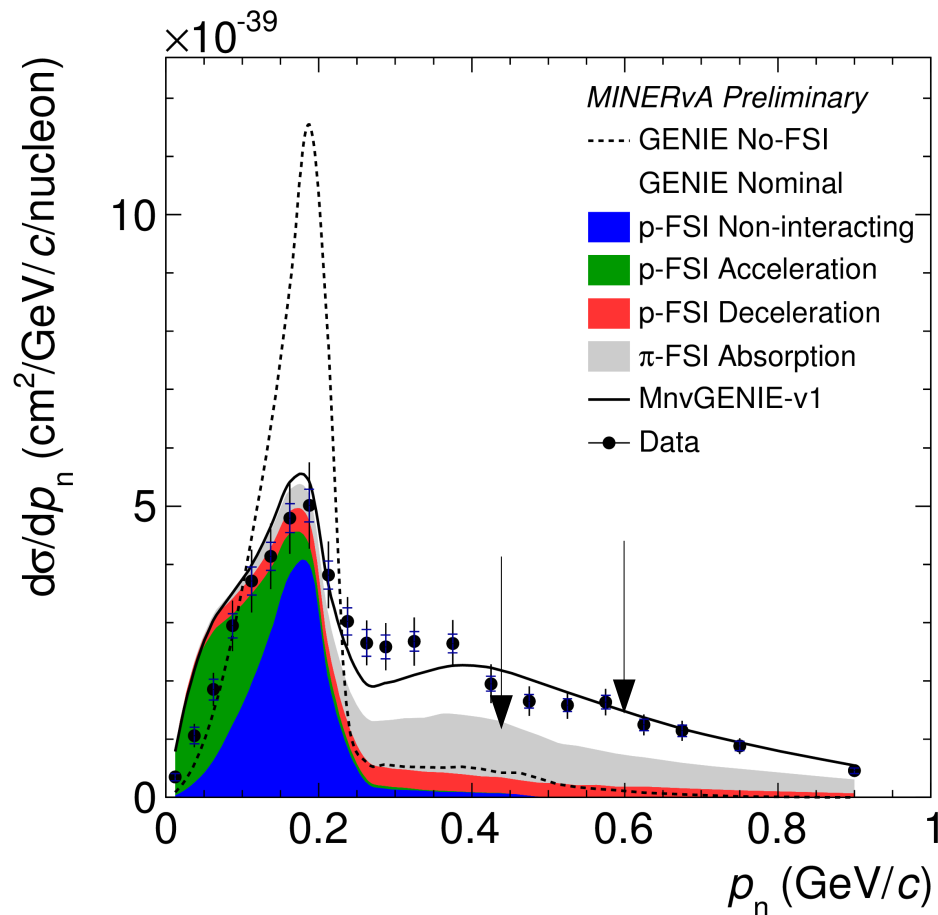
(= MnvGENIE-v1 – GENIE Nominal)



Pion production and 2p2h process:
strong intra-nuclear momentum transfer
due to momentum sharing with proton

Nuclear Effect Diagnostics

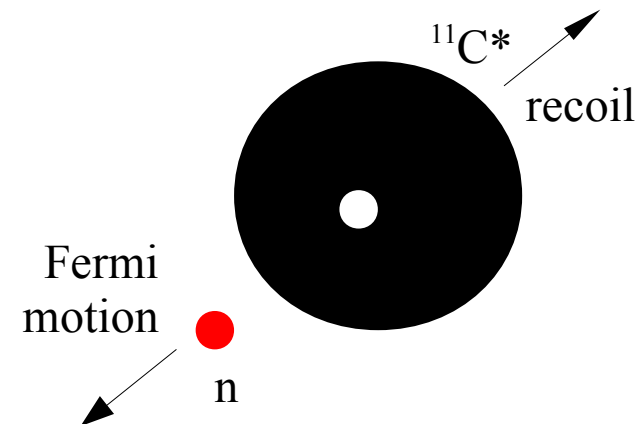
- CCQE with Fermi motion
- FSI deceleration vs. acceleration
- Pionless resonant production, pion absorption FSI, and 2p2h



GENIE describes the tail reasonably well due to large contribution from 2p2h tuned to MINERvA inclusive measurements

p_n : smeared δp_T beyond QE peak \rightarrow tail

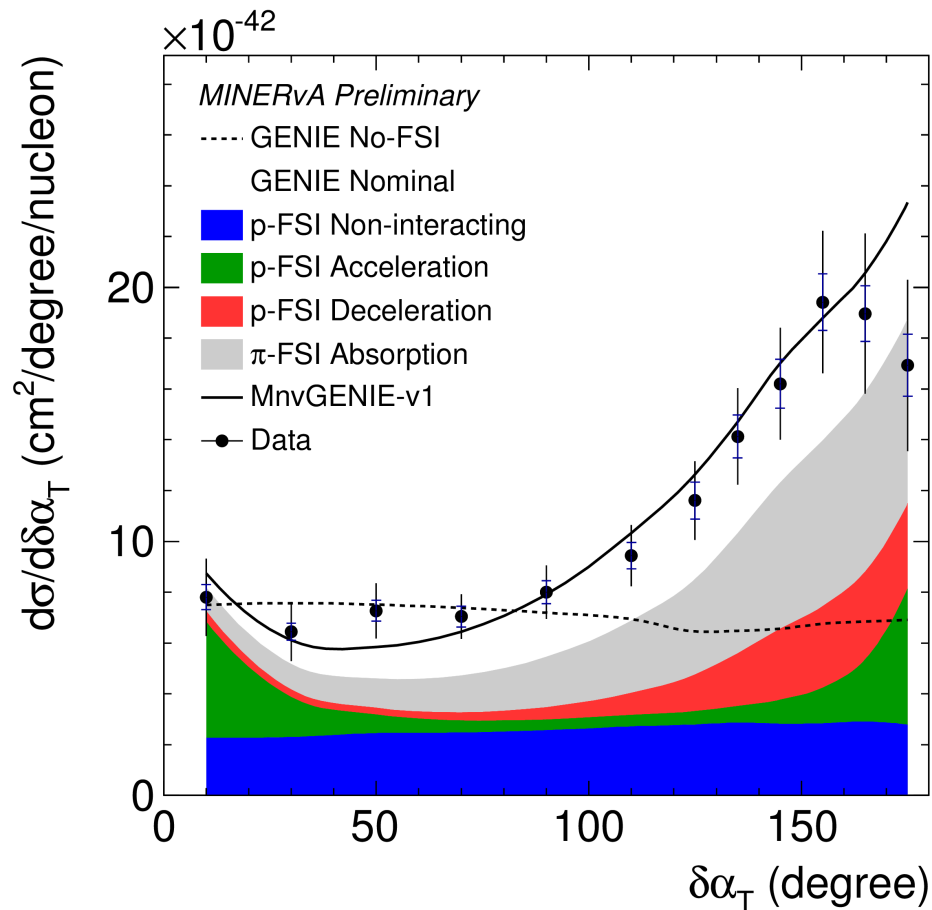
- π -FSI Absorption
- 2p2h
(= MnvGENIE-v1 – GENIE Nominal)



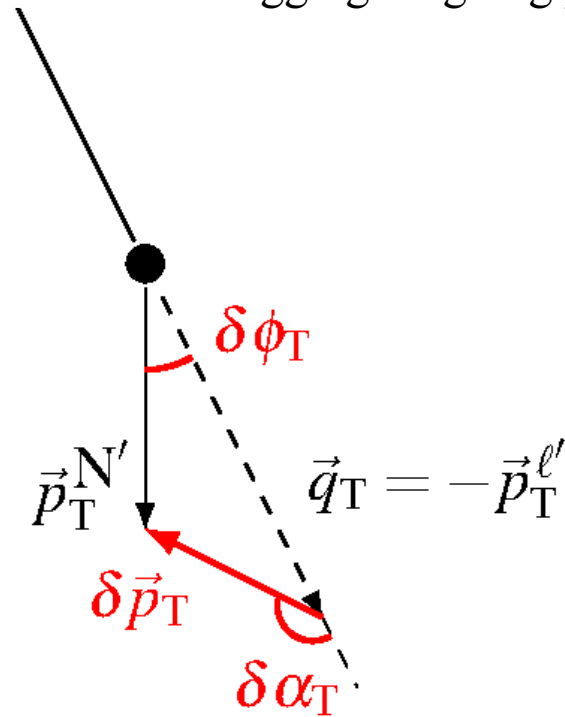
Pion production and 2p2h process:
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Nuclear Effect Diagnostics

- CCQE with Fermi motion
- FSI deceleration vs. acceleration
- Pionless resonant production, pion absorption FSI, and 2p2h



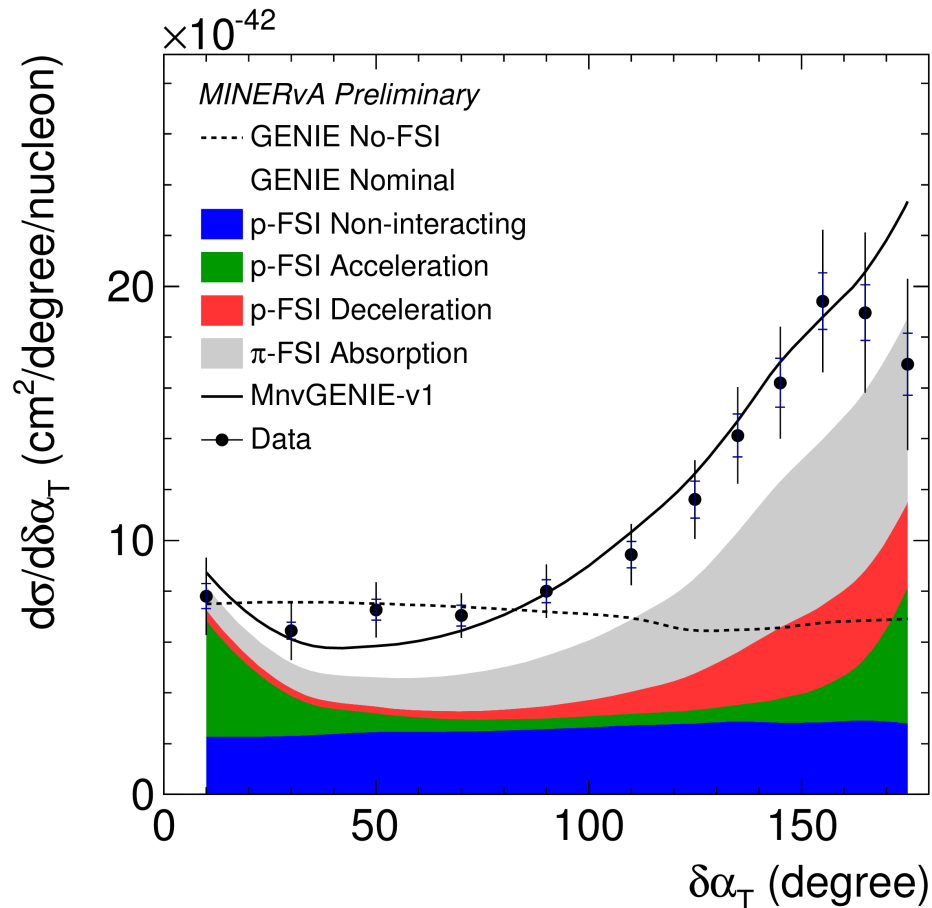
δp_T (nuclear effects)
dragging outgoing proton



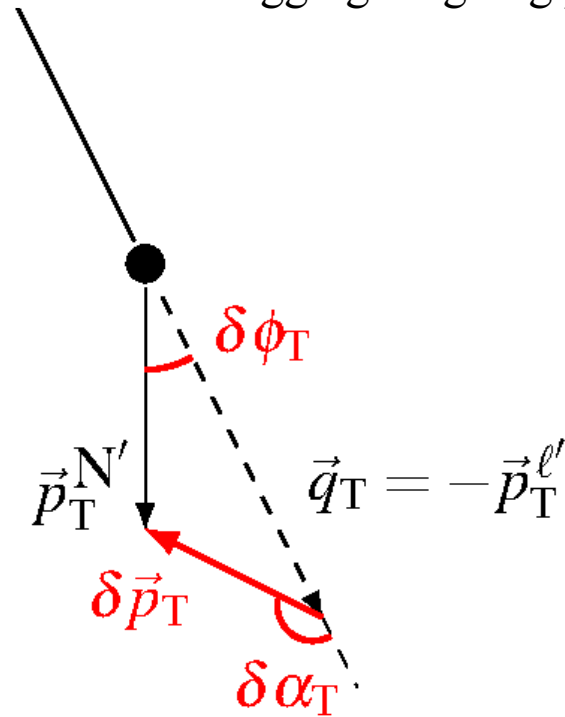
Proton momentum shared by others,
decelerated \rightarrow large $\delta \alpha_T$ region
– π -FSI Absorption

Nuclear Effect Diagnostics

- CCQE with Fermi motion
- FSI deceleration vs. acceleration
- Pionless resonant production, pion absorption FSI, and 2p2h



δp_T (nuclear effects)
dragging outgoing proton



Proton momentum shared by others,
decelerated \rightarrow large $\delta\alpha_T$ region

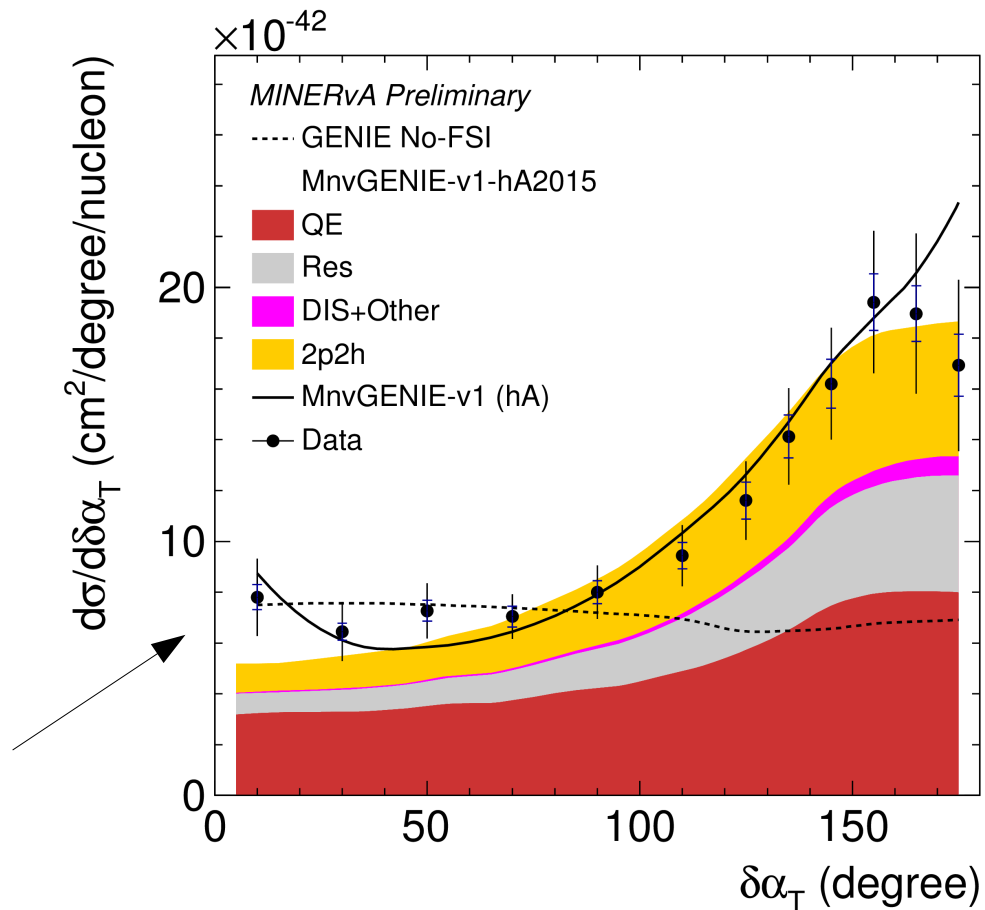
– π -FSI Absorption

– 2p2h (= MnvGENIE-v1 – GENIE Nominal)

ADVANCED TOPICS: GENIE FSI_s

Advanced Topics: GENIE FSI

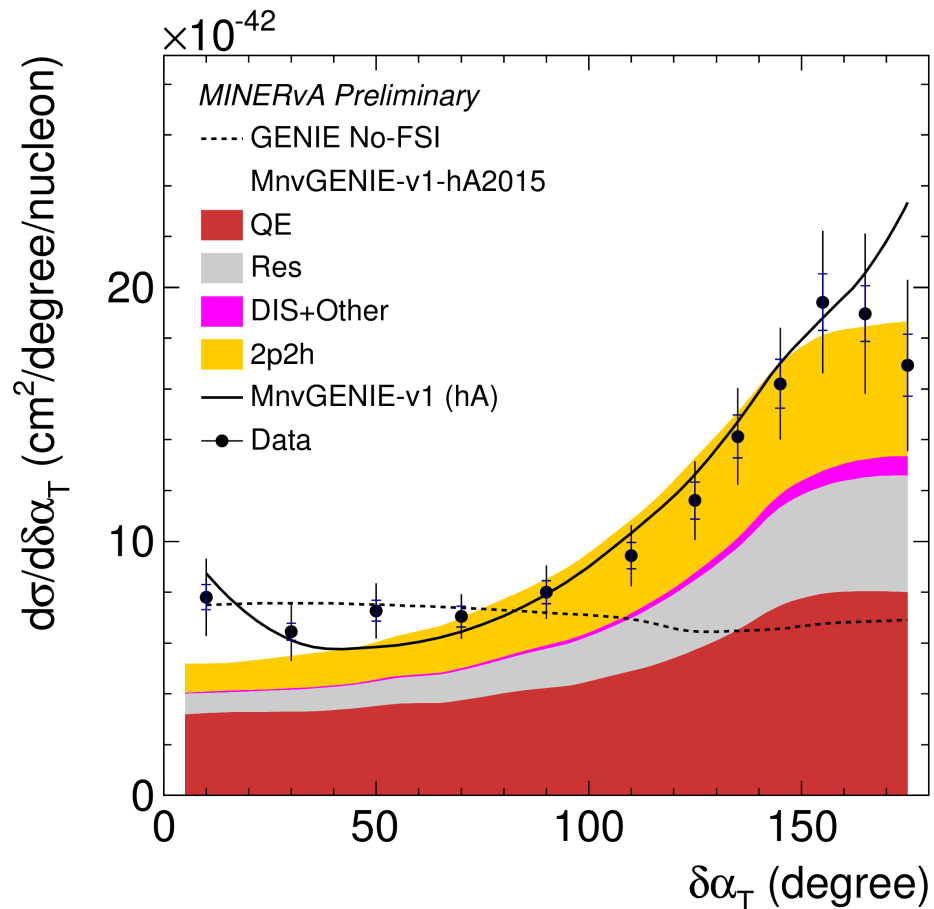
- (pre2015) hA: effective model, include “elastic component” in intranuclear scattering, used in GENIE MINERvA Tune (v1)
- hA2015: removed “elastic component”, replacing hA in MnvGENIE-v1-hA2015



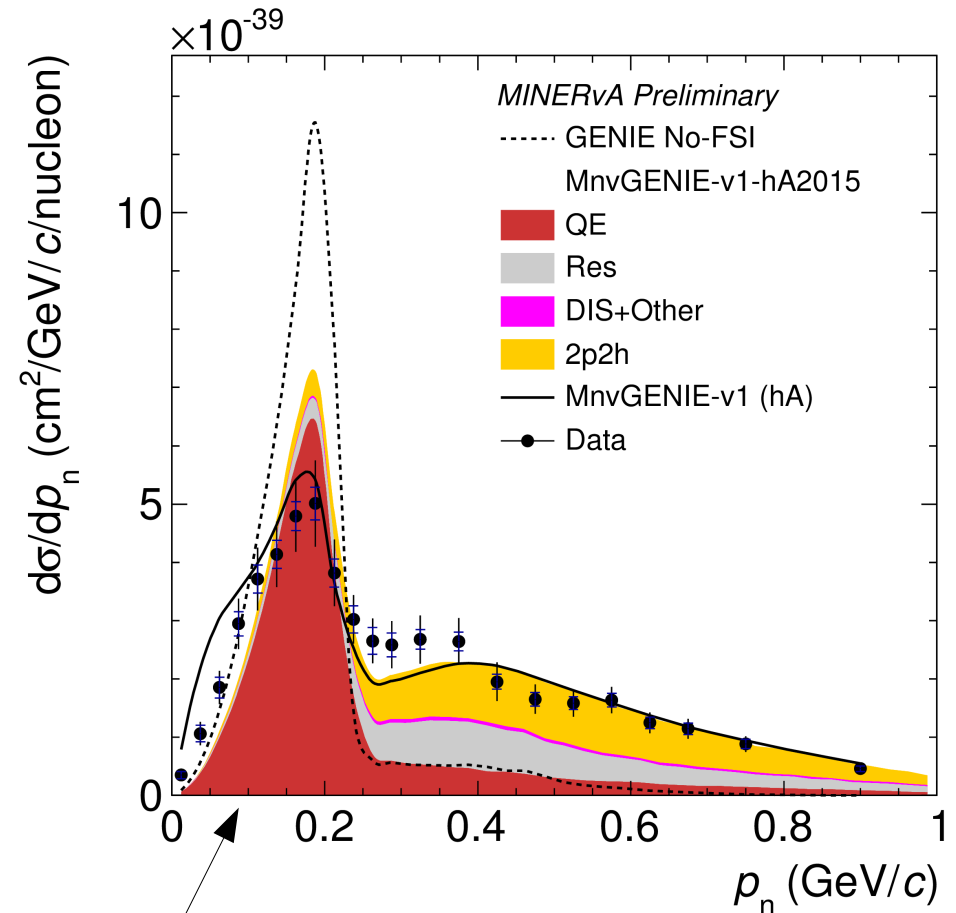
No p-FSI acceleration

Advanced Topics: GENIE FSI

- (pre2015) hA: effective model, include “elastic component” in intranuclear scattering, used in GENIE MINERvA Tune (v1)
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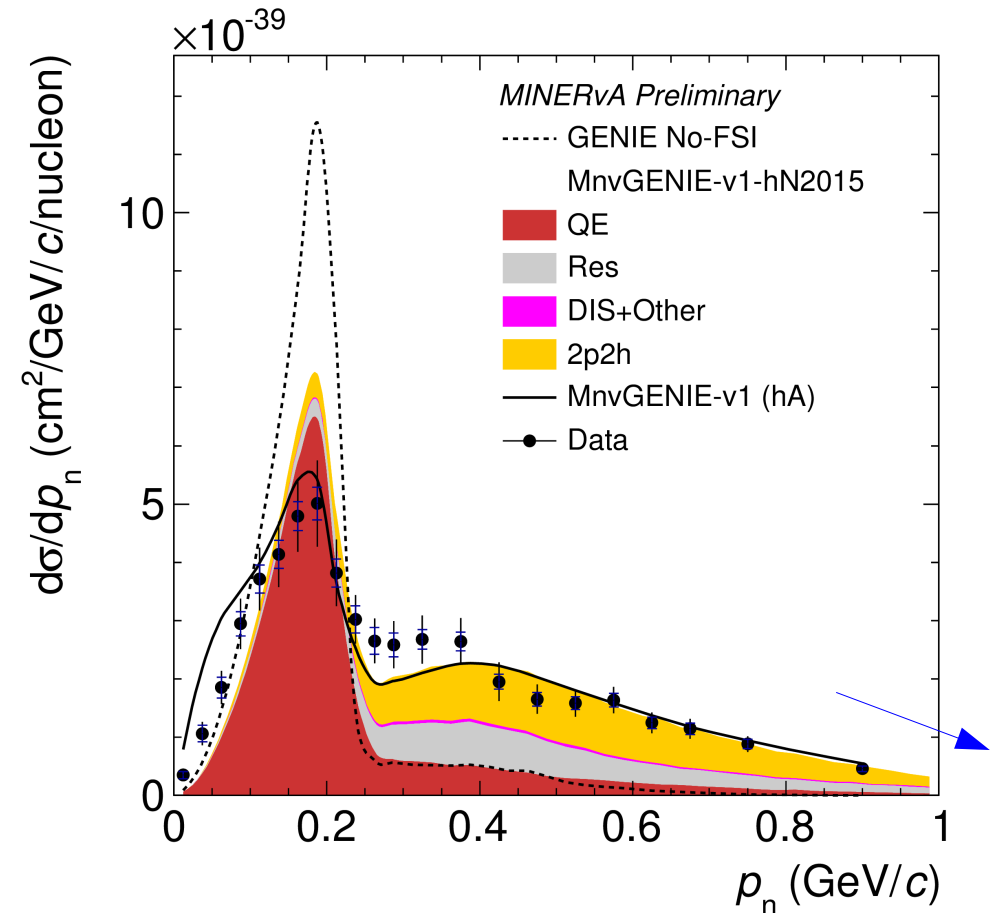
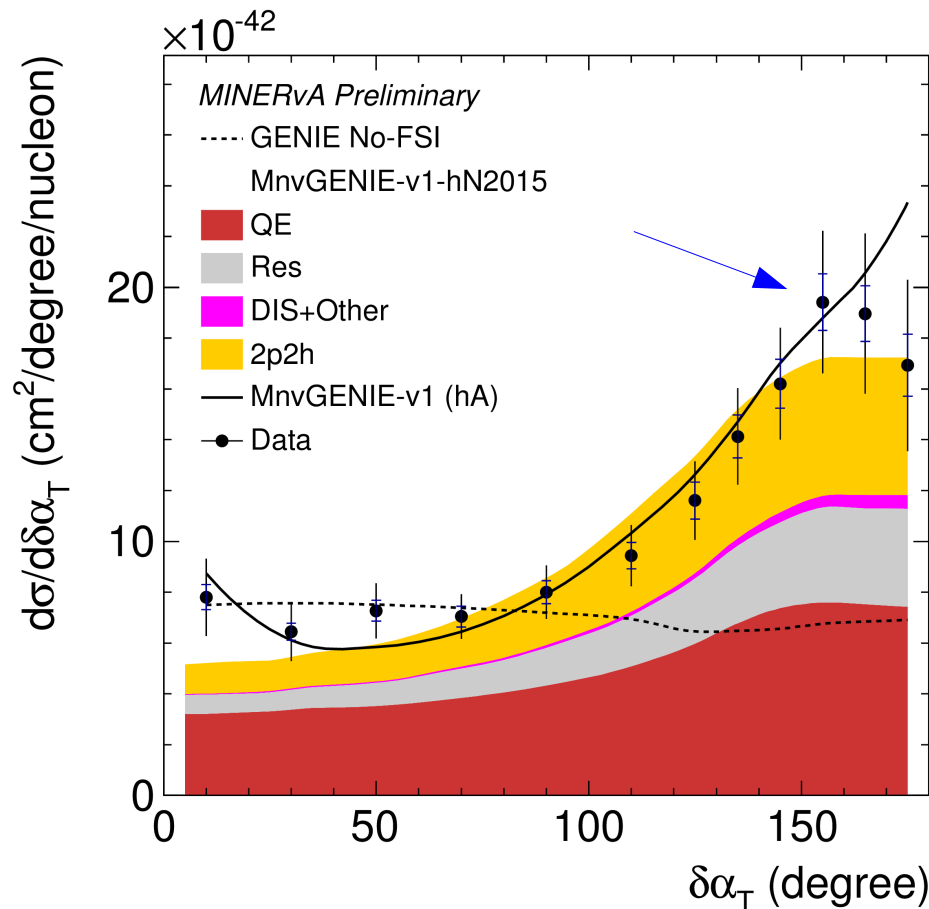
No p-FSI acceleration



QE peak not distorted, but much narrower

Advanced Topics: GENIE FSI

- (pre2015) hA: effective model, include “elastic component” in intranuclear scattering, used in GENIE MINERvA Tune (v1)
- hA2015: removed “elastic component”, replacing hA in MnvGENIE-v1-hA2015
- hN2015: full cascades + Oset, replacing hA in MnvGENIE-v1-hN2015



hA2015 and hN2015 difference in

- Large $\delta\alpha_T$ (concentrated)
- p_n tail (diluted)

ADVANCED TOPICS: NUWRO

Simulation: GENIE [Nucl.Instrum.Meth. A614 (2010) 87-104]

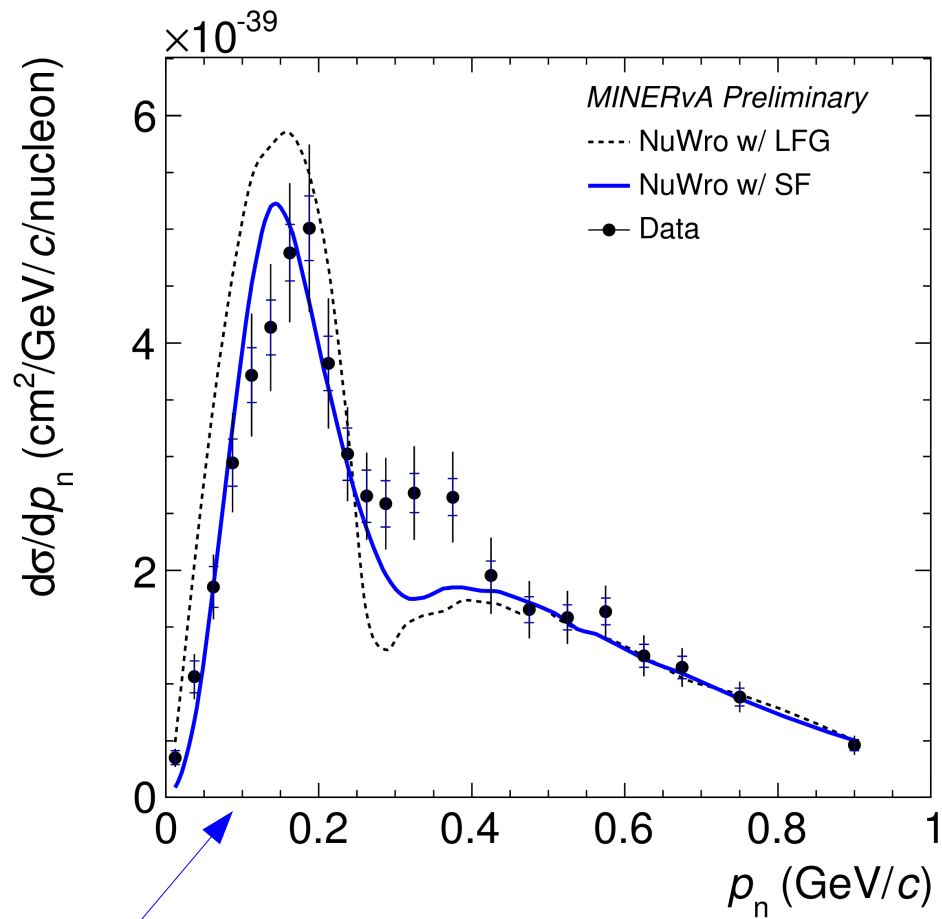
- **Nominal:** version 2.8.4
 - ✓ global Fermi Gas (RFG) model with Bodek-Ritchie (BR) tail [Phys. Rev. D 23, 1070 (1981)]
 - ✓ hA FSI [AIP Conf.Proc. 1405 (2011) 213-218]
- **No-FSI:** Nominal without FSI
- **MnvGENIE-v1: GENIE MINERvA Tune (v1) [only 2p2h relevant for this analysis]**
 - Added Random Phase Approximation (RPA) [Phys.Rev. C70 (2004) 055503]
 - Non-resonance pion production scaled down by 75% [Phys.Rev. D90 (2014) no.11, 112017]
 - Valencia 2p2h [Nieves *et al.*, Phys.Lett. B707 (2012) 72-75, Phys. Rev. C 86, 015504 (2012), Phys.Rev. D88 (2013) no.11, 113007, arXiv:1601.02038]
 - ✓ tuned to MINERvA inclusive data → significant enhancement in small 4-momentum transfer region [Phys.Rev.Lett. 116 (2016) 071802]

Simulation: NuWro [Phys.Rev. C86 (2012) 015505]

- Version: 11q
 - ✓ Local Fermi Gas (LFG) or Spectral Function (SF) [Benhar *et al.*, Nucl.Phys. A579 (1994) 493-517]
 - ✓ FSI: intranuclear cascades of hadronic interactions + Oset model [Nucl.Phys. A484 (1988) 557-592]
- Valencia 2p2h [Nieves *et al.*, Phys.Lett. B707 (2012) 72-75, Phys. Rev. C 86, 015504 (2012)]

Advanced Topics: NuWro

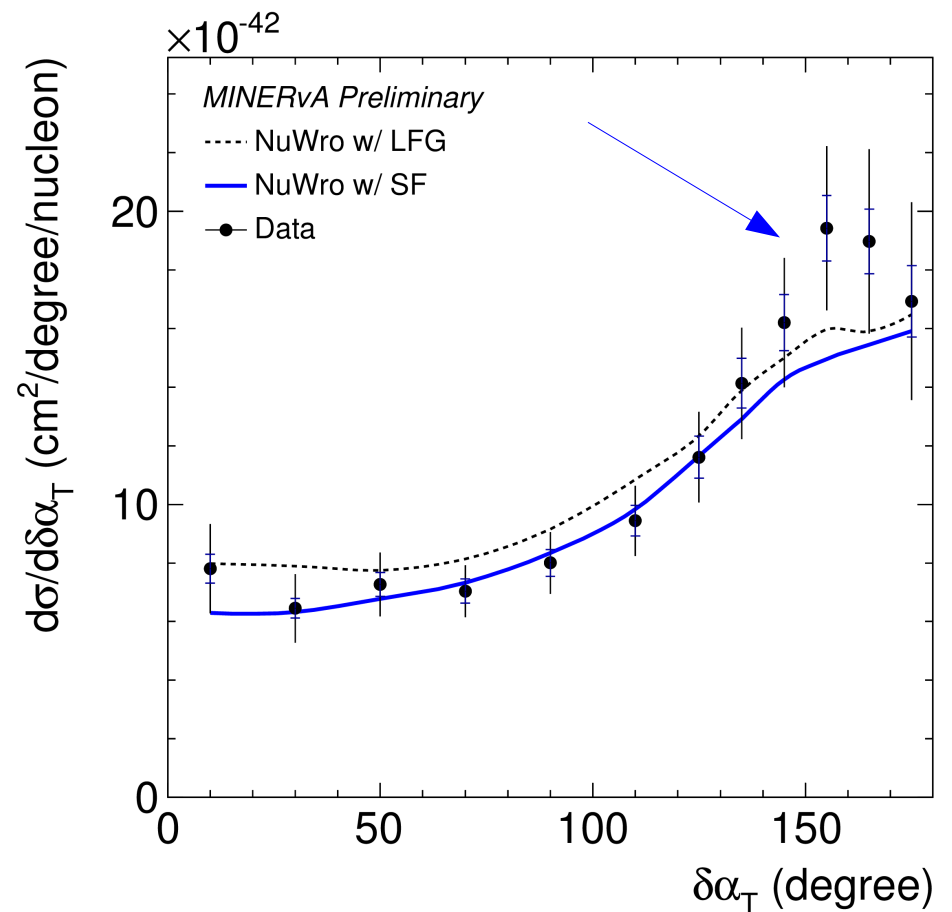
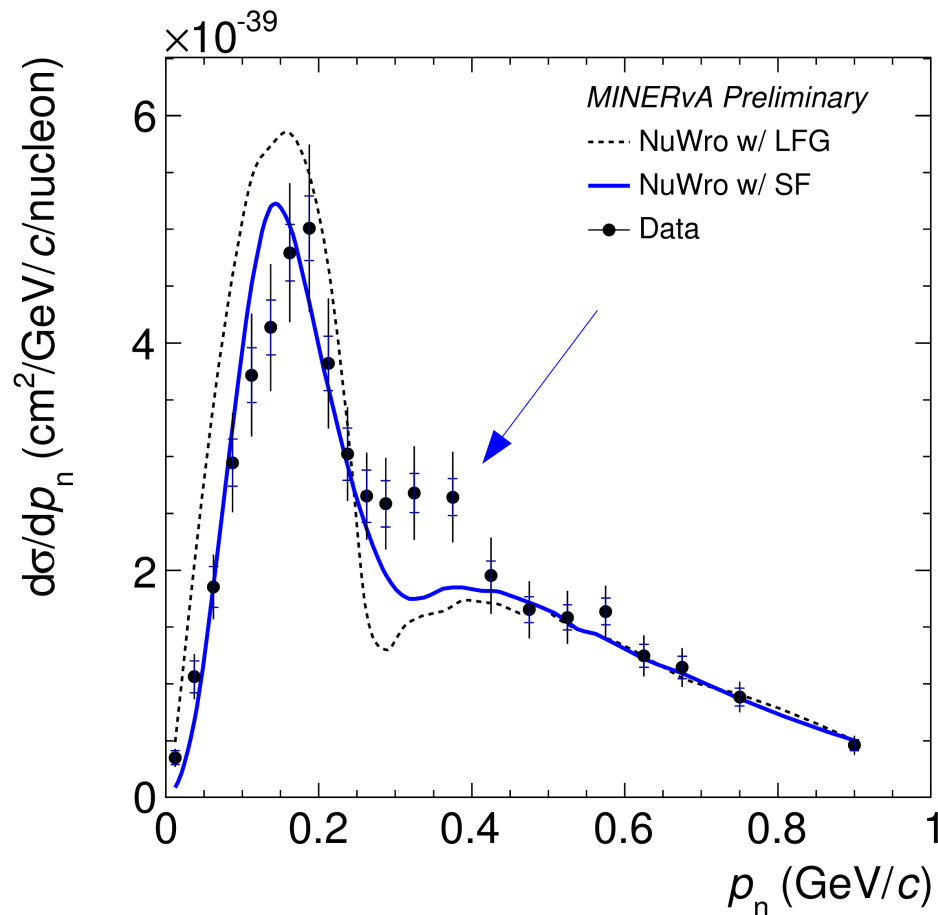
- Fermi motion



SF describes Fermi motion very well

Advanced Topics: NuWro

- Fermi motion
- Resonance / 2p2h strength



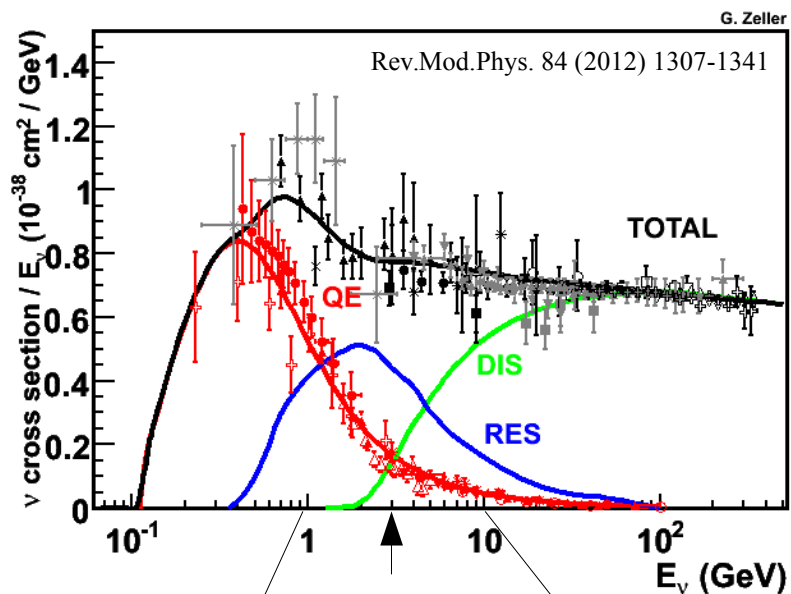
SF describes Fermi motion very well
Resonance / 2p2h lacks of strength in small regions

ADVANCED TOPICS: COMPARISON TO T2K

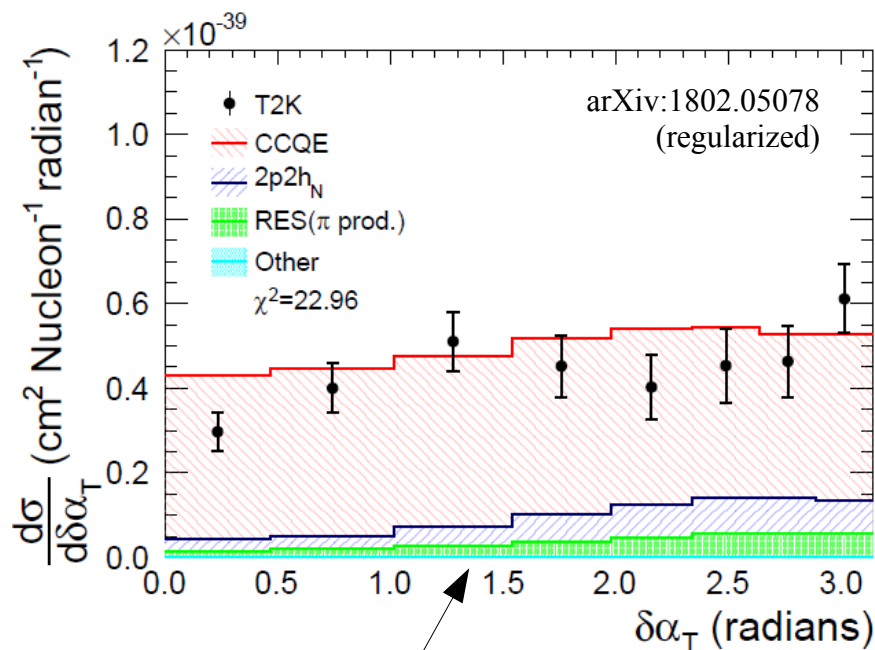
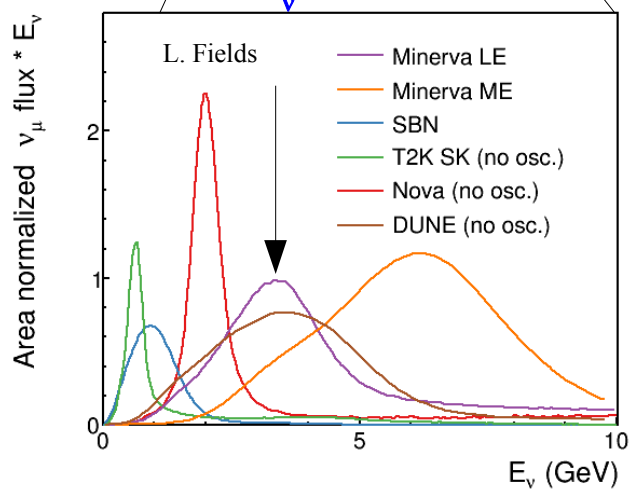
Advanced Topics: Comparison to T2K

[arXiv:1802.05078] *same target, slight difference in signal phase space definition

- $\delta\alpha_T$



NuMI LE
 $\langle E_\nu \rangle \sim 3$ GeV

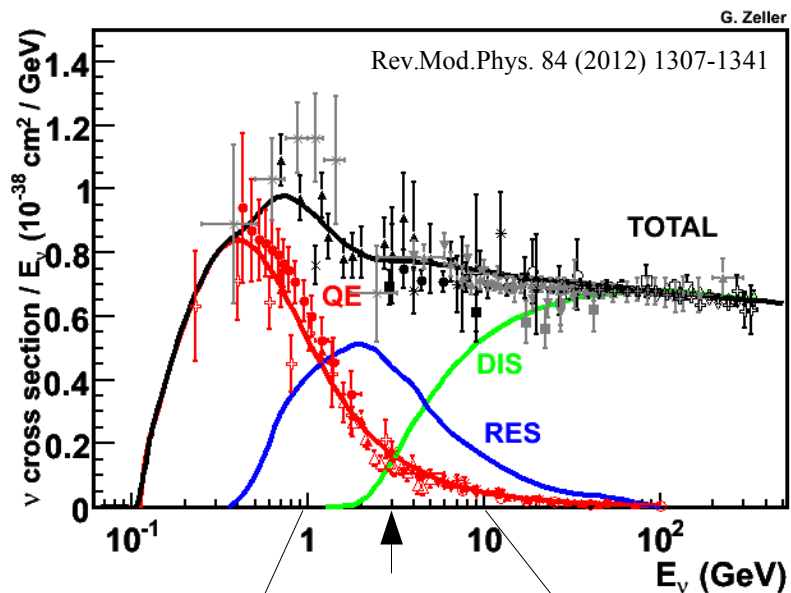


MINERvA-T2K difference mainly due to RES:
Very small resonance contribution at T2K

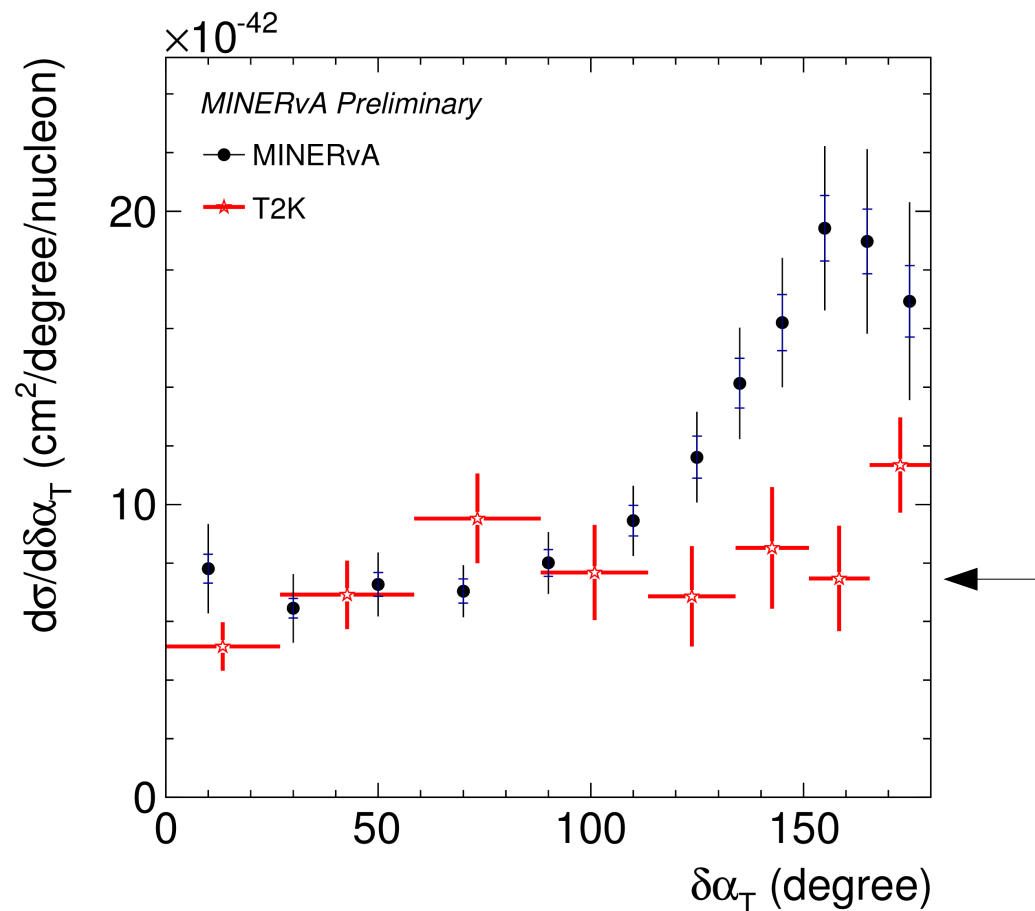
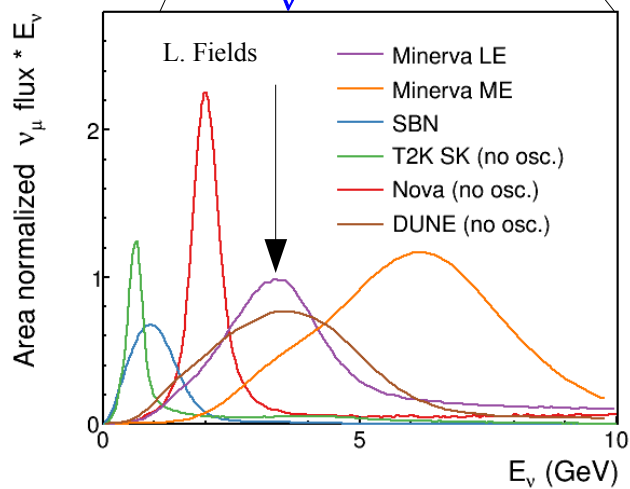
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NuMI LE
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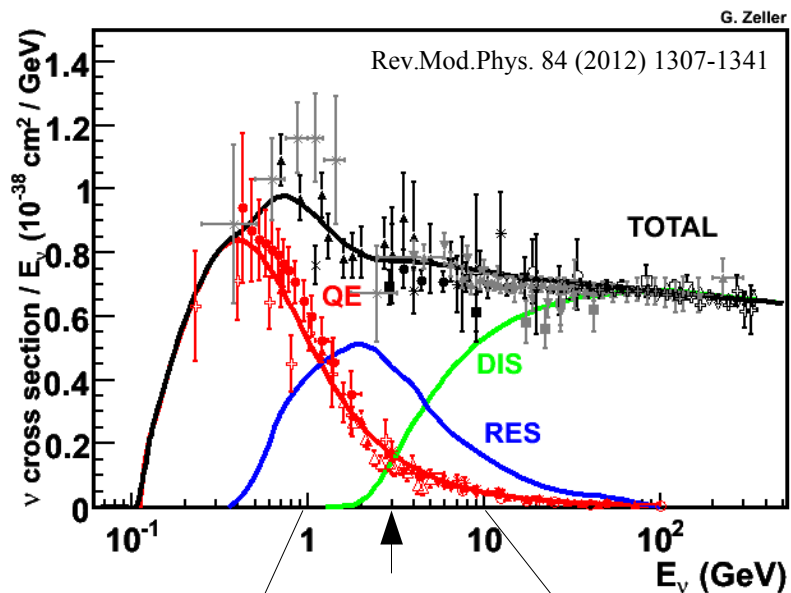


MINERvA-T2K difference mainly due to RES
Fermi motion (isotropic) baseline consistent

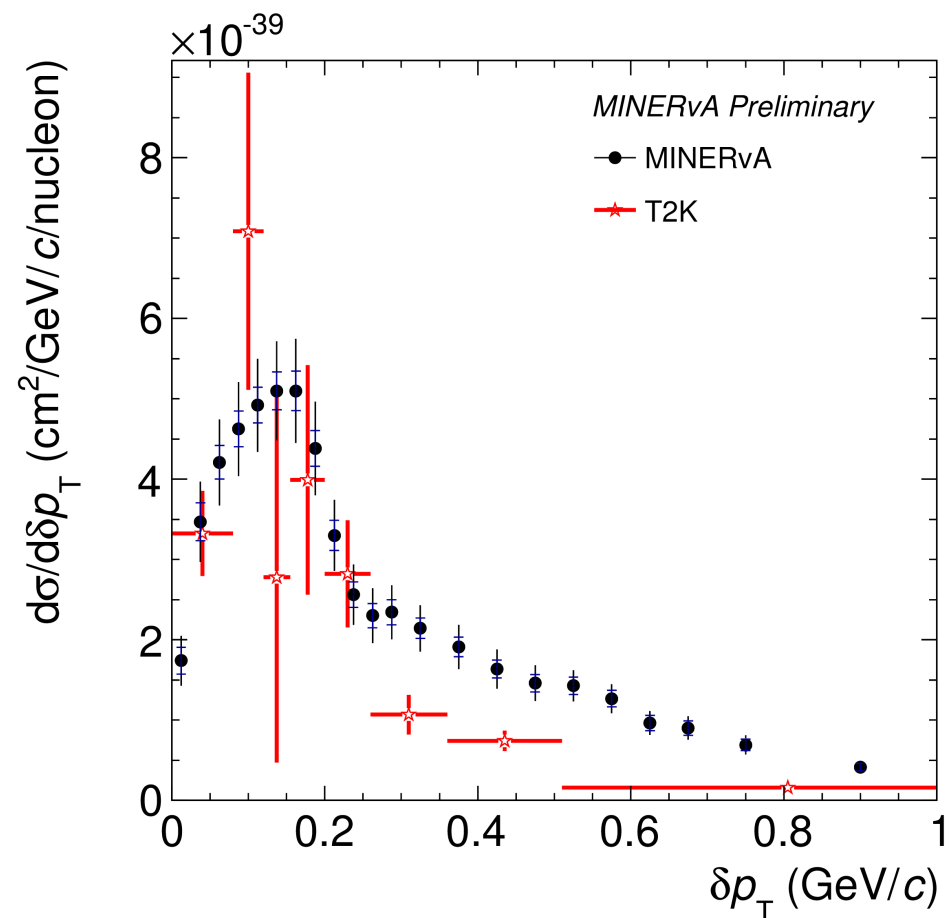
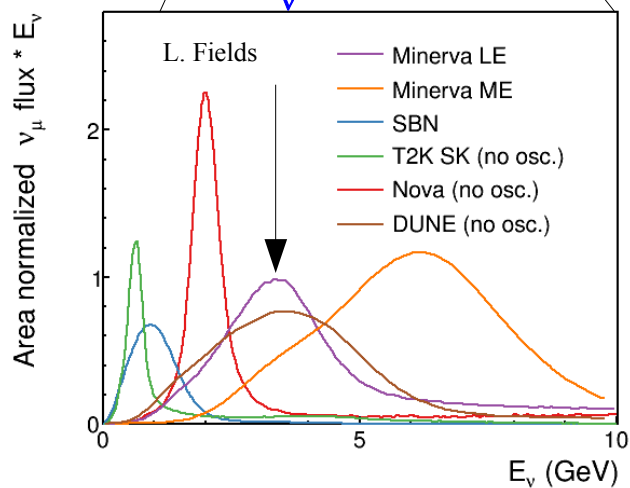
Advanced Topics: Comparison to T2K

[arXiv:1802.05078] *same target, slight difference in signal phase space definition

- $\delta\alpha_T$
- δp_T



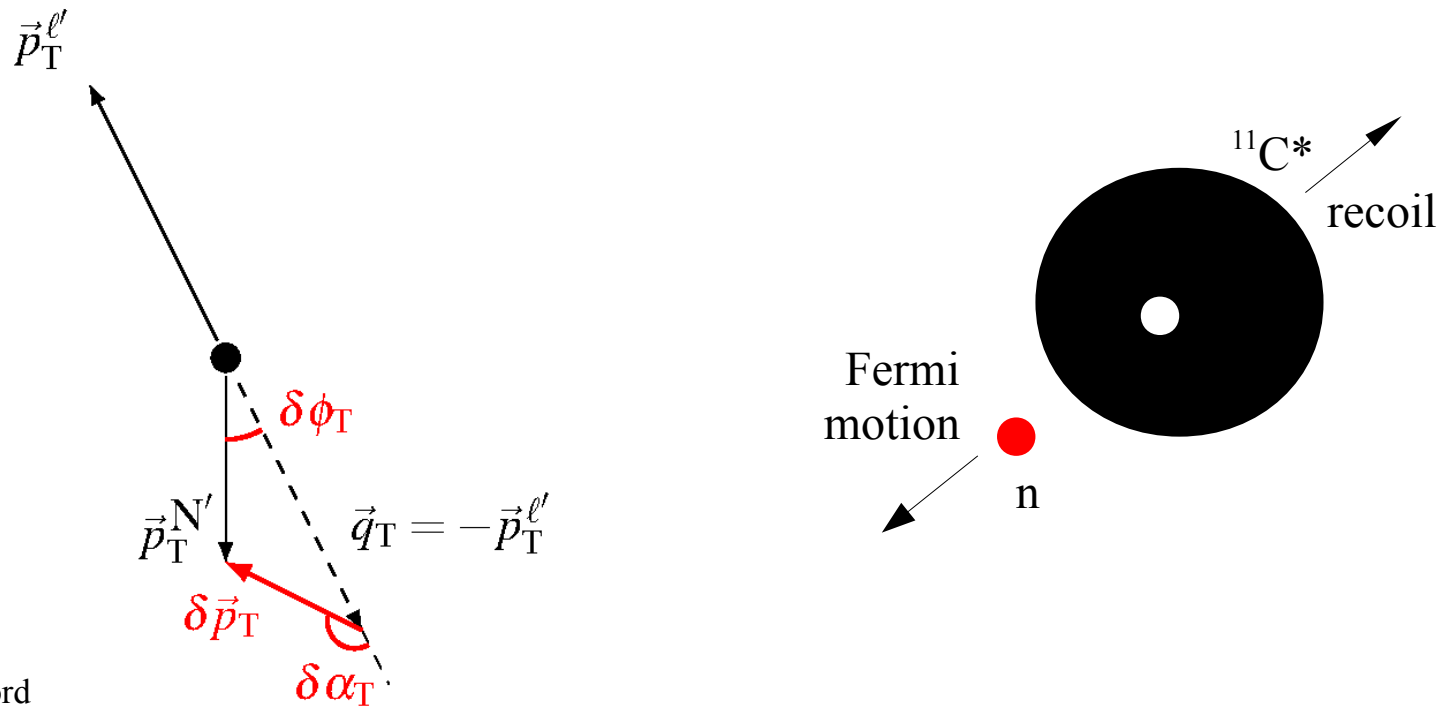
NuMI LE
 $\langle E_\nu \rangle \sim 3$ GeV



MINERvA-T2K difference mainly due to RES
 The QE peaks are consistent

Summary and Outlook

- Muon-proton mesonless production at MINERvA
 - 2014: LE neutrino beam, CH target
 - 2016: LE neutrino beam, CH + nuclear targets
 - **This analysis**: LE neutrino beam, CH (3.28×10^{20} POT)
 - **Future: medium energy neutrino beam CH + nuclear targets** ($E_\nu \sim 6$ GeV, 12×10^{20} POT)
- In this analysis, we have shown
 - ✓ Single-particle kinematics (muon and proton momentum and angle)
 - ✓ Transverse kinematic imbalances ($\delta\alpha_T$, δp_T)
 - ✓ Initial neutron momentum (p_n)

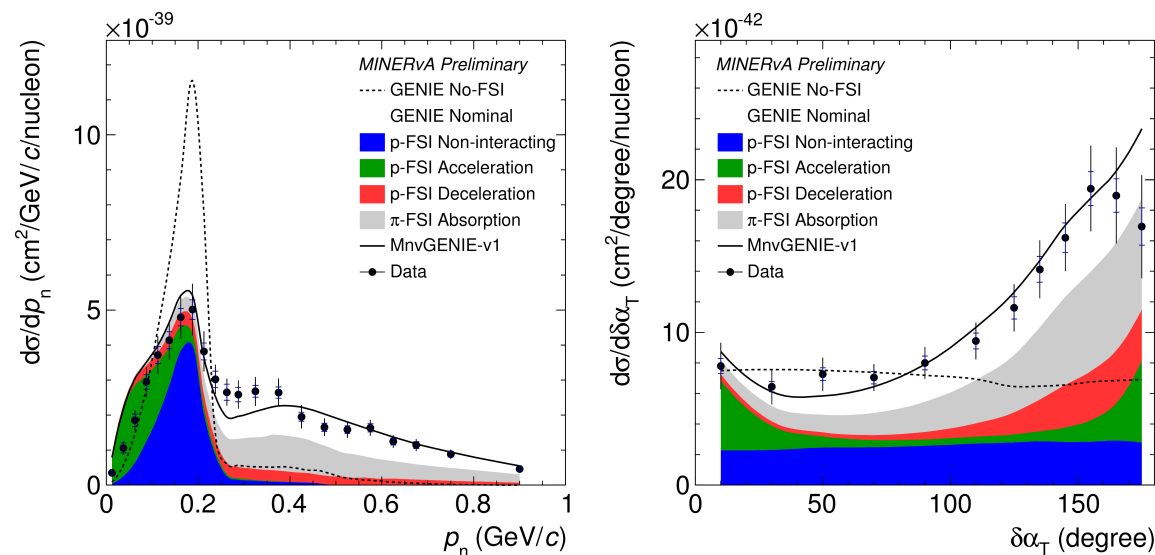


Summary and Outlook

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 - **Future: medium energy neutrino beam CH + nuclear targets** ($E_\nu \sim 6$ GeV, 12×10^{20} POT)
- In this analysis, we have shown
 - ✓ Single-particle kinematics (muon and proton momentum and angle)
 - ✓ Transverse kinematic imbalances ($\delta\alpha_T$, δp_T)
 - ✓ Initial neutron momentum (p_n)

By rearranging final-stat kinematics, nuclear effects can be diagnosed:

- ✓ p_n strong constraint to Fermi motion
- ✓ $\delta\alpha_T$ factors out Fermi motion uncertainty and have direct sensitivity to FSI



Summary and Outlook

- Muon-proton mesonless production at MINERvA
 - 2014: LE neutrino beam, CH target
 - 2016: LE neutrino beam, CH + nuclear targets
 - **This analysis**: LE neutrino beam, CH (3.28×10^{20} POT)
 - **Future: medium energy neutrino beam CH + nuclear targets** ($E_\nu \sim 6$ GeV, 12×10^{20} POT)
- In this analysis, we have shown
 - ✓ Single-particle kinematics (muon and proton momentum and angle)
 - ✓ Transverse kinematic imbalances ($\delta\alpha_T$, δp_T)
 - ✓ Initial neutron momentum (p_n)

By rearranging final-stat kinematics, nuclear effects can be diagnosed:

- ✓ p_n strong constraint to Fermi motion
- ✓ $\delta\alpha_T$ factors out Fermi motion uncertainty and have direct sensitivity to FSI

Interesting observation:

- GENIE MINERvA Tune (v1)
 - ✓ Describes data well to first order
 - ✓ Critical component is Valencia 2p2h tuned to MINERvA inclusive data
- NuWro
 - ✓ SF provides very good description of data

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 - ✓ Single-particle kinematics (muon and proton momentum and angle)
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- **New developments**
 - Transverse kinematic imbalances
 - ✓ New system to solve the nuclear effect problem in neutrino interaction most relevant for oscillation measurements
 - ✓ Radical approach → double transverse kinematic imbalance [Phys. Rev. D 92, 051302(R)]
 - First measurement of Furmanski-Sobczyk initial neutron momentum
 - ✓ diagnostic power
 - ✓ Practically efficient way to select pure CCQE events (beyond the scope of this talk)
 - ESC (Elastically Scattered Contained) proton selection
 - ✓ Powerful to enhance the proton reconstruction quality
 - ✓ Application in other homogeneous non-magnetized detectors, e.g. LAr

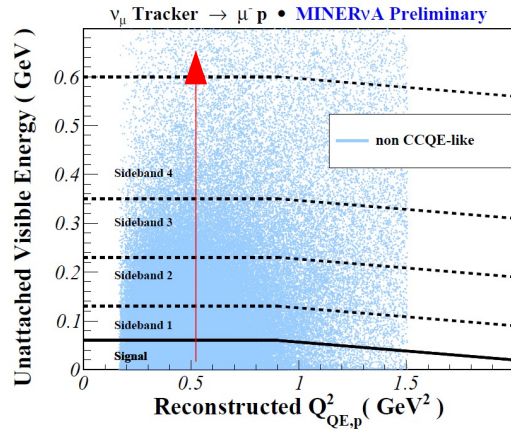


Source: <http://www.cnhubei.com/ztmjys-pyts>

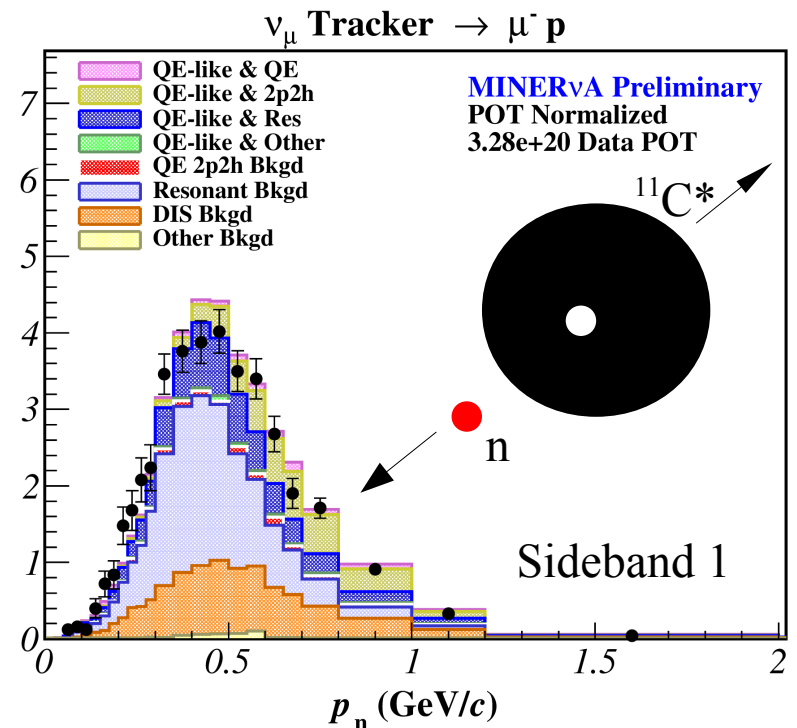
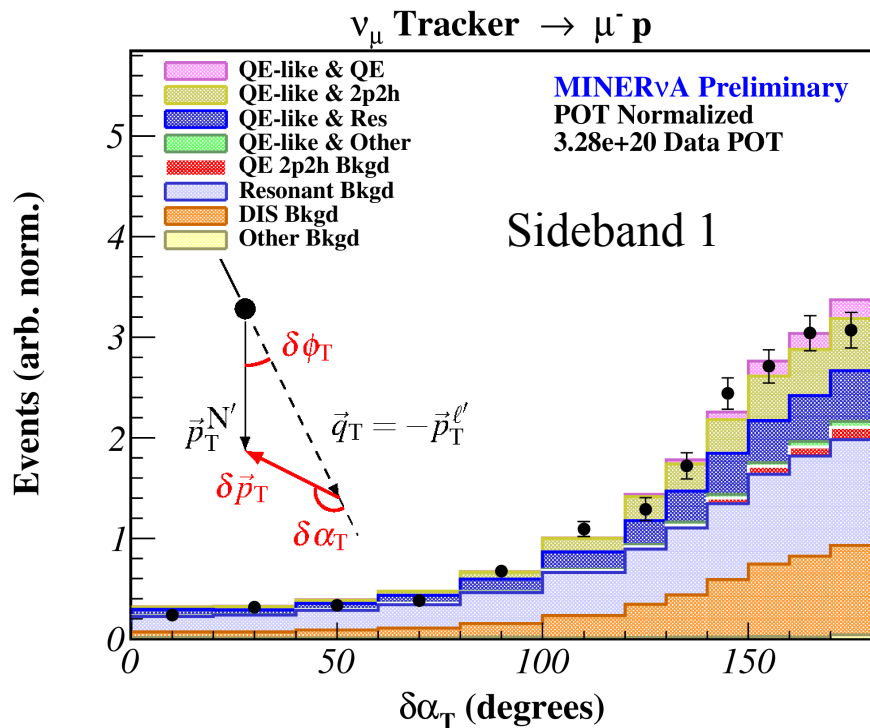
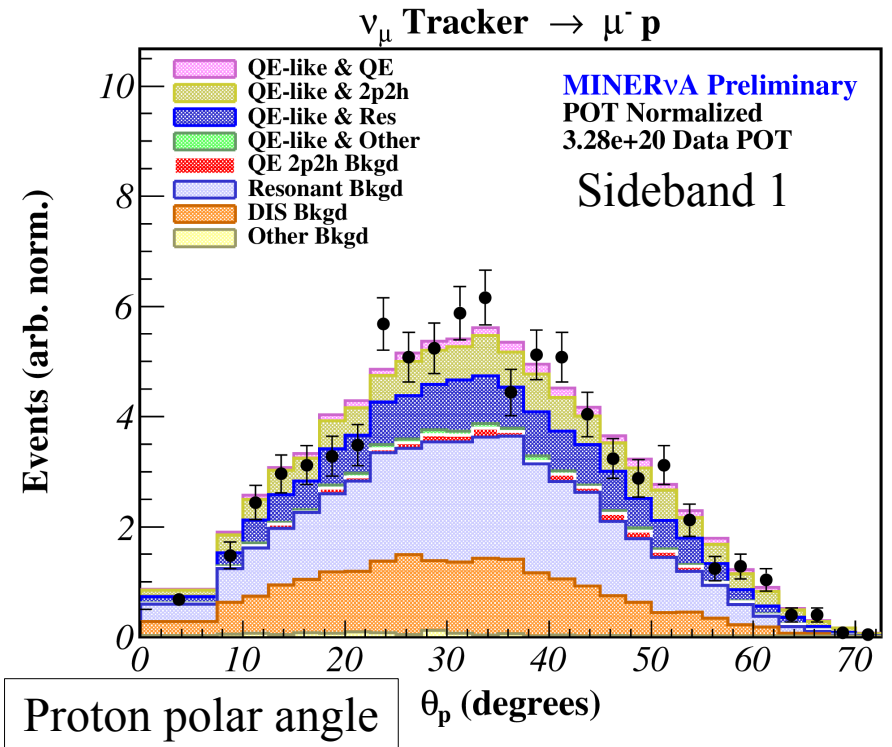
BACKUP

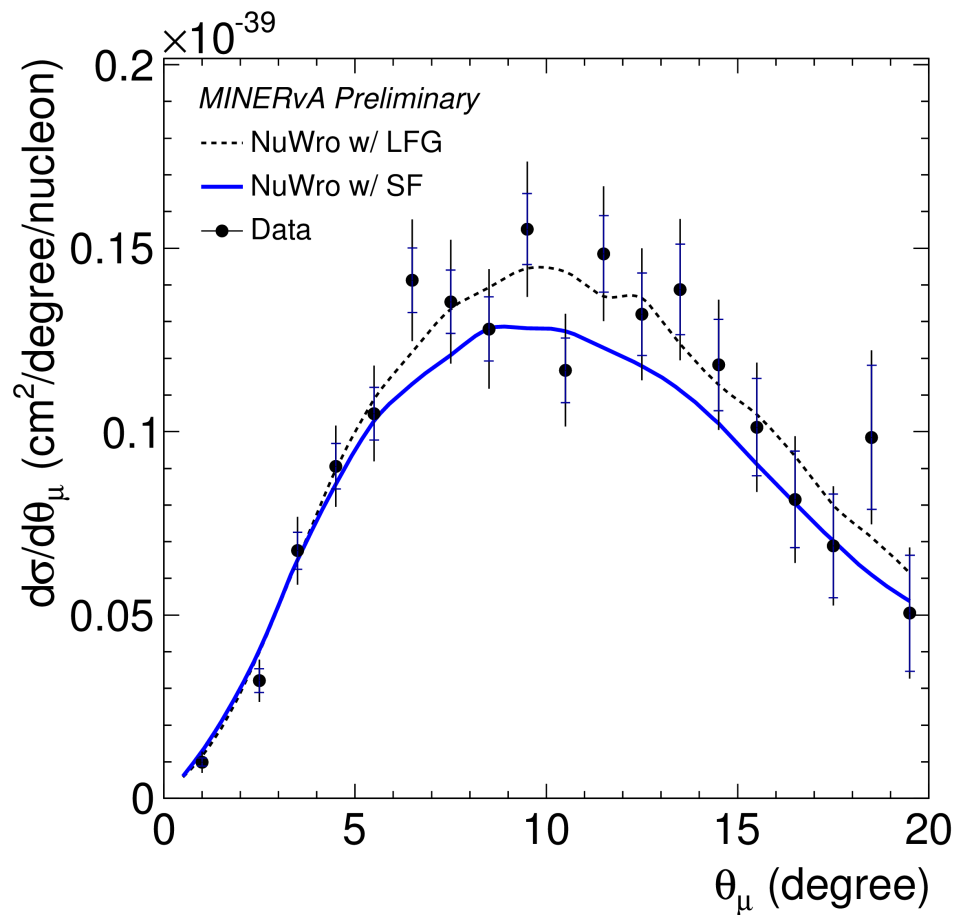
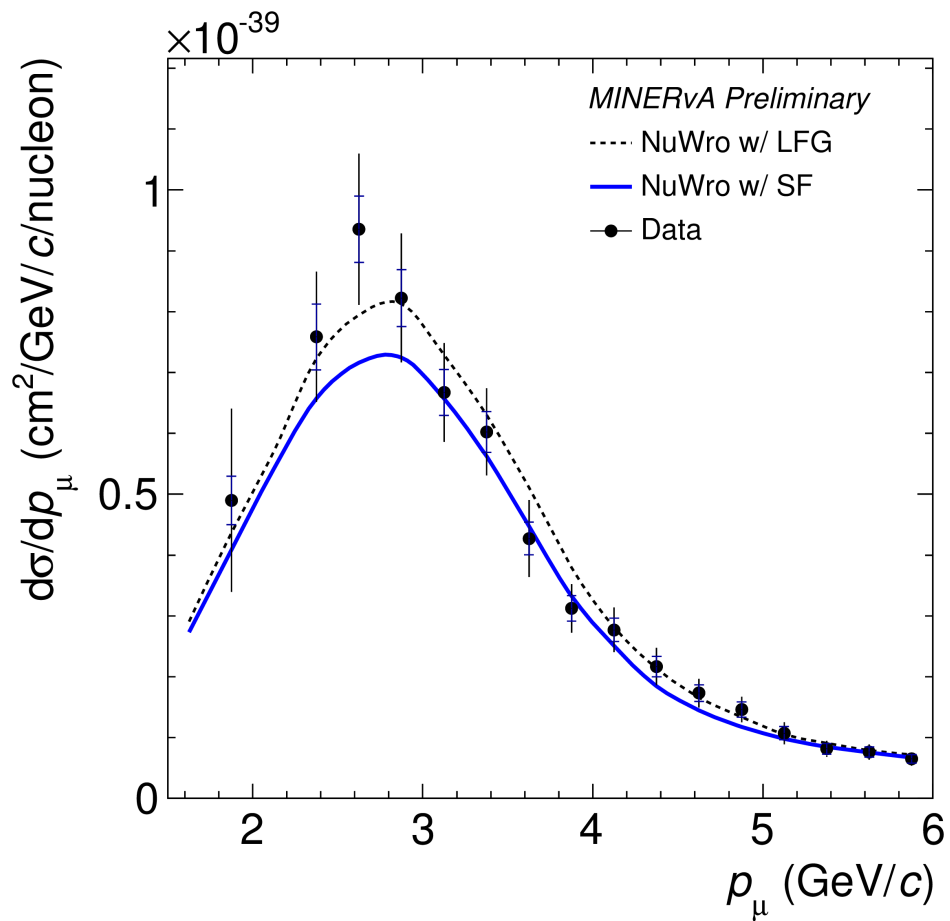
Background Estimation

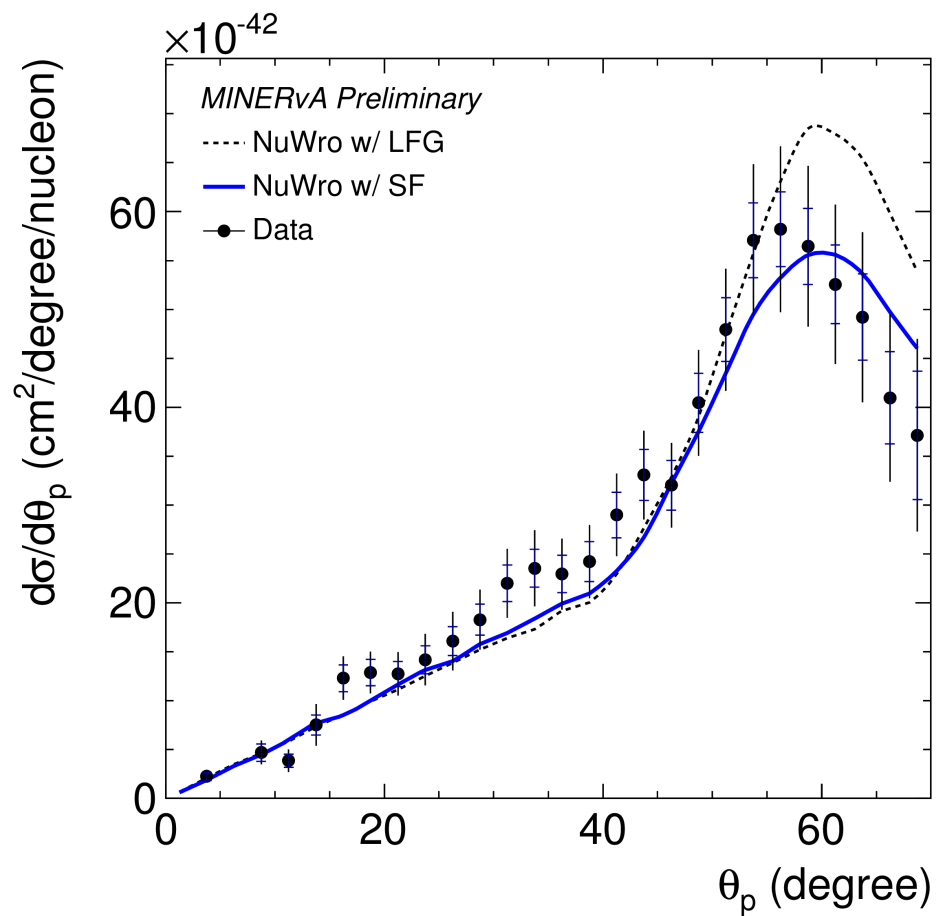
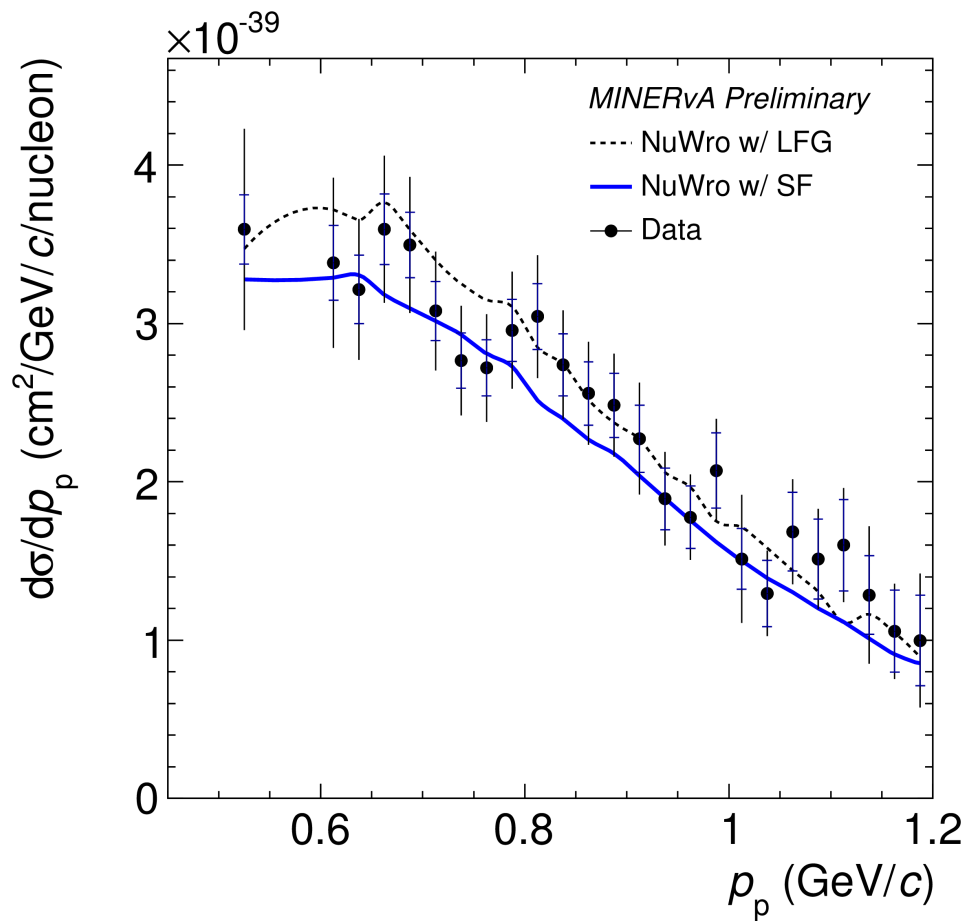
- Data-MC comparison at reconstructed level after sideband fit



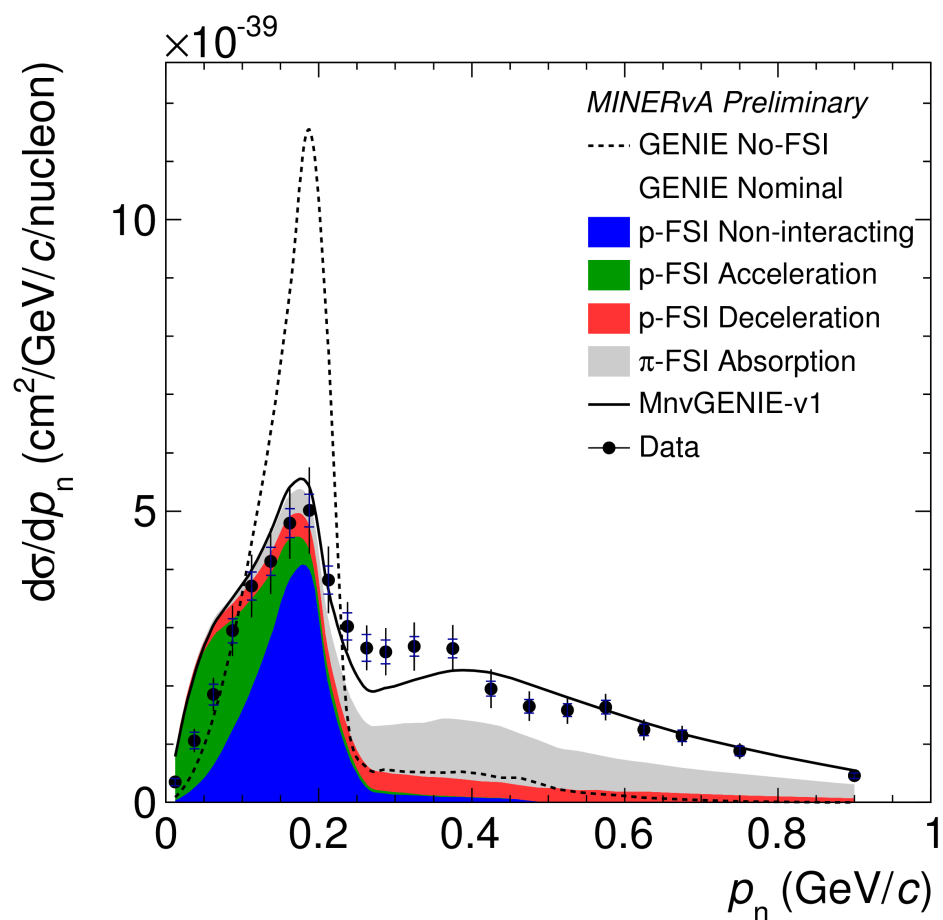
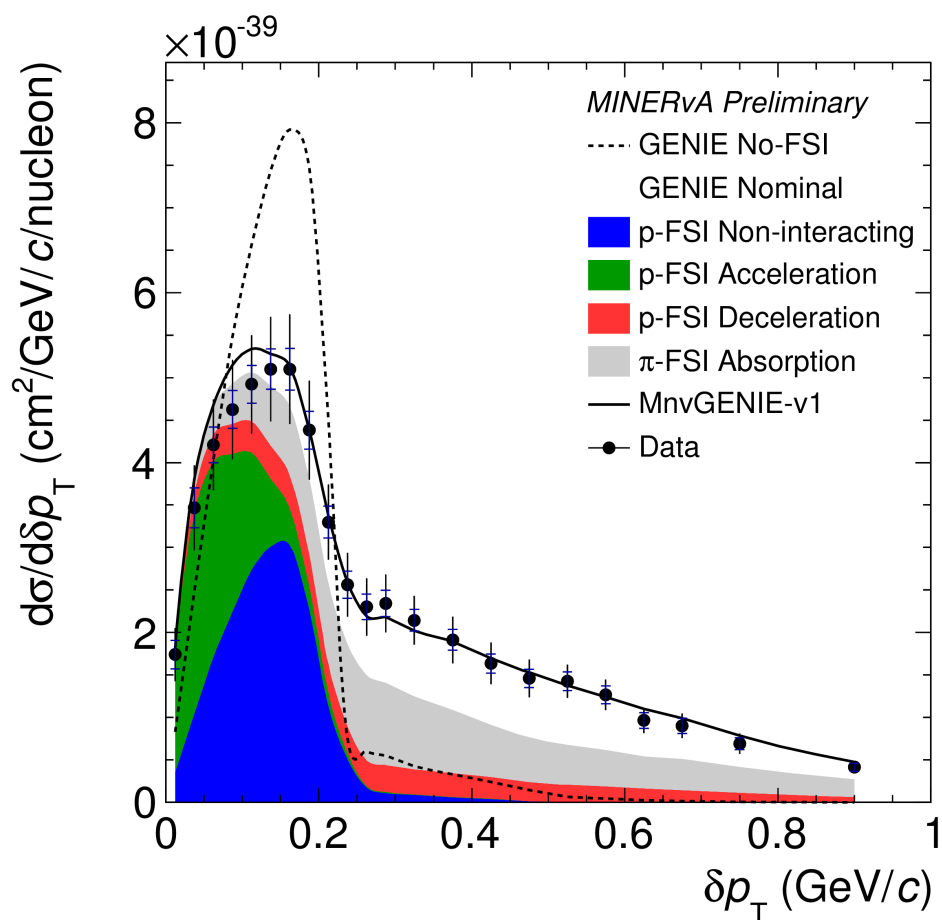
Sideband 1 \rightarrow 4 high background direction







Interpretation of δp_T



$$\delta \vec{p}_T = \vec{p}_T^N - \Delta \vec{p}_T$$

$$p_n \equiv \sqrt{\delta p_T^2 + \delta p_L^2}$$

Only differ by longitudinal momentum imbalance
 p_n has better resolution

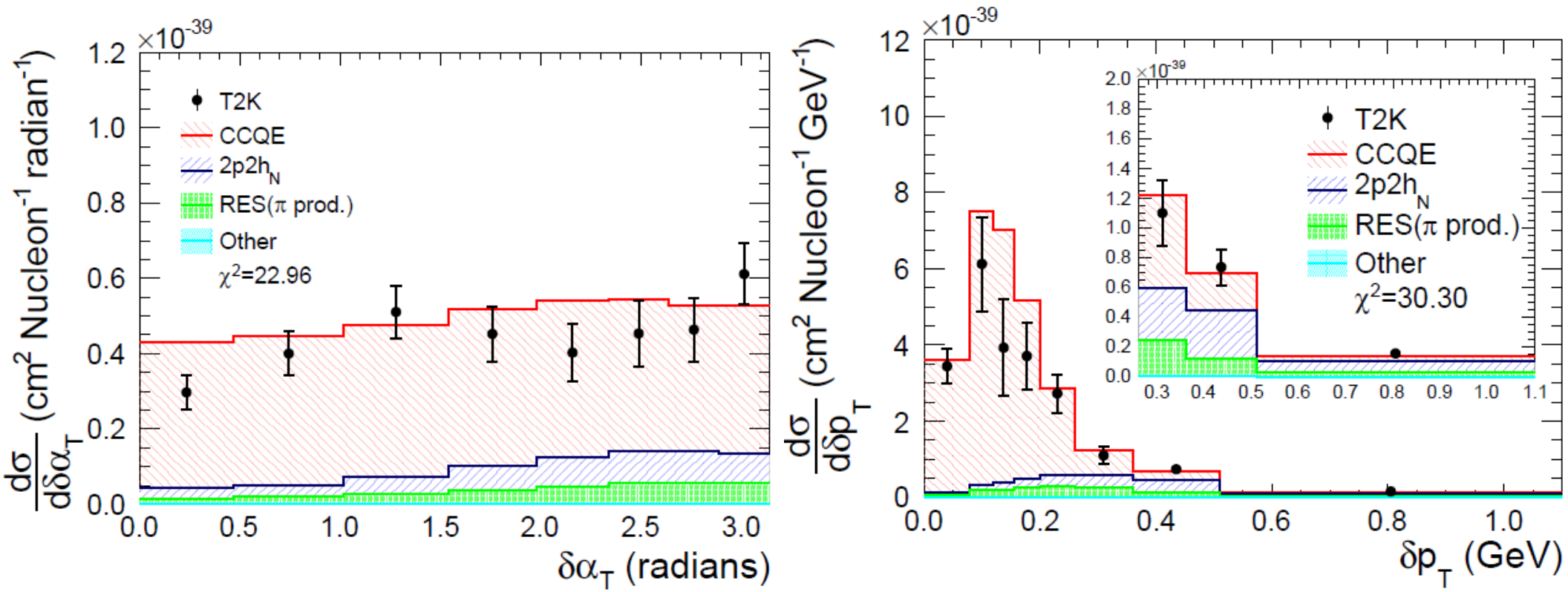
66.4 deg



Analysis	p_p	$\cos\theta_p$	p_μ	$\cos\theta_\mu$
Multi-dimensional	> 500 MeV	-	-	-
STV	450-1000 MeV	> 0.4	> 250 MeV	> -0.6
Inferred kinematics	> 450 MeV	> 0.4	-	-

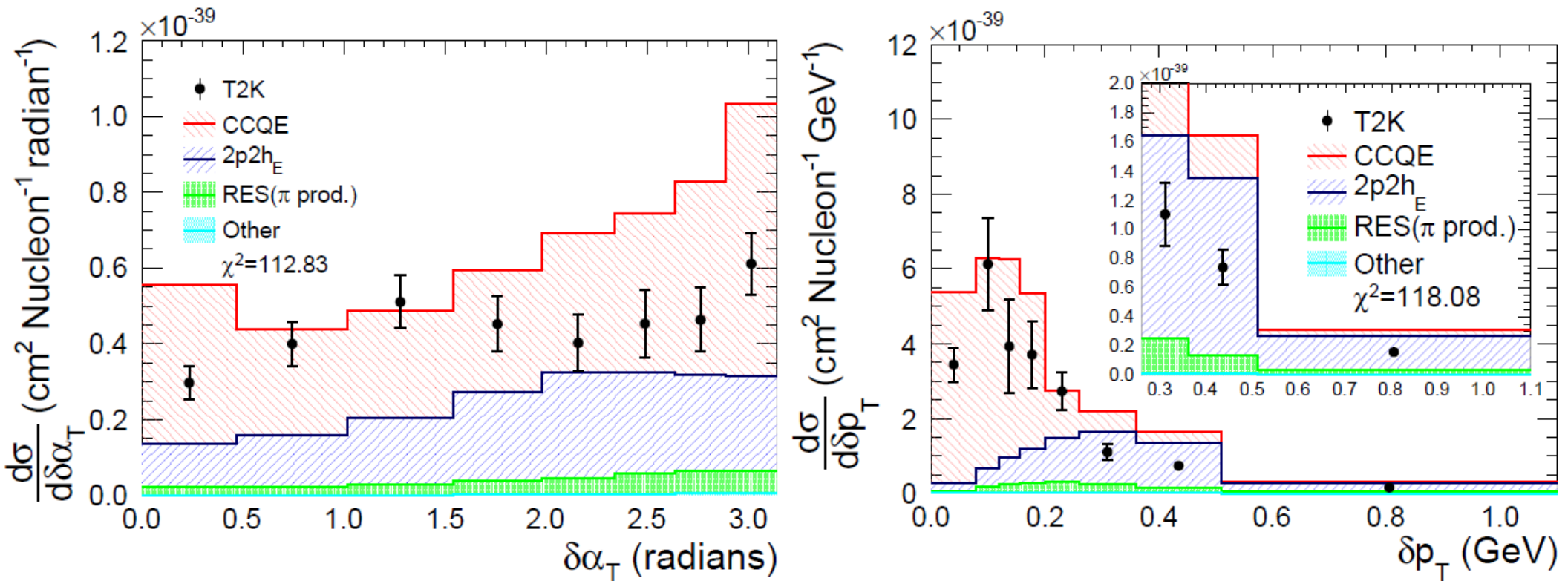
TABLE I. Signal phase space restrictions for the three analyses.

[arXiv:1802.05078](https://arxiv.org/abs/1802.05078)



NEUT for T2K has very small RES

FIG. 16. The extracted differential cross section as a function of the single transverse variables compared to: the NEUT 5.3.2.2 simulation with the SF initial state model and an ad hoc 2p2h model (**left**); the same NEUT simulation with various scalings of the mean free path of nucleons undergoing FSI processes to simulate different FSI strengths (**right**). 2p2h_N indicates the Nieves et. al. model of Ref. [76] implemented in NEUT. A comparison of the NEUT prediction without a 2p2h contribution is also shown. More details of these models can be found in Sec. IV A. The ‘N’ subscript after LFG indicates that the model is using both a 1p1h and 2p2h prediction from the aforementioned model of Nieves et. al. The inlays on the left plots show a close-up of the tail regions of δp_T and $\delta \alpha_T$ whilst those on the right show the same comparisons on a logarithmic scale.



GENIE for T2K has very small RES

FIG. 17. The extracted differential cross section as a function of the single transverse variables compared to: the GENIE 2.12.4 simulation (**left**) and the GiBUU 2016 simulation (**right**). GENIE uses the Bodek and Richie RFG initial state model and this prediction also includes GENIE’s empirical 2p2h prediction ($2p2h_E$). This GENIE prediction is similar that used as a starting point for the NO ν A experiment’s oscillation analyses. More details of these models can be found in Sec. IV A. The inlays on the plots show a close-up of the tail regions of δp_T and $\delta\alpha_T$.

Both GENIE taken from left column

END