

## Presentations :

how **not** to do them .

- Introduction
- Awful example
- Discussion

# Presentations :

how not to do them .

Why so negative ?

Kipling:

"There are nine and sixty ways  
of constructing tribal laws,  
And every single one of them is right!"

## Introduction:

Learn to do good presentations :

- (a) by presenting ;
- (b) by making mistakes ;
- (c) by enjoying good presentations ;
- (d) by enduring bad presentations .

## Useful reading

- D.E. Knuth, T. Larabee, P.M. Roberts  
*Mathematical Writing*, 1989, MAA  
"great prizes are ... interest  
and understanding ; all else  
is secondary."
- S. Senn, "On Thinking and Learning"  
*RSS News & Notes* 35.4 (2007) 1-3  
"Content is of primary importance"
- E. Tufte, "Powerpoint is Evil"  
[http://www.wired.com/wired/  
archive/11.09/ppt2.html](http://www.wired.com/wired/archive/11.09/ppt2.html)  
"Imagine a widely used and  
expensive prescription drug that  
promised to make us beautiful  
but didn't."

**Start of awful  
example**

A double integral

$$\eta = \sqrt{x^2 + y^2} + \sqrt{(n-x)^2 + y^2}$$

$$\alpha = \tan^{-1} \frac{y}{x} + \tan^{-1} \frac{y}{n-x}$$

$$\frac{1}{2} \iint (\alpha - \sin \alpha) \exp\left(-\frac{1}{2}(\eta - n)\right) dx dy$$

Obviously

we can ignore

$$\int_{\pi/2}^{\pi} \int_0^{\infty} (\alpha - s \cos \theta) \exp(-\frac{1}{2}(q-n)) r dr d\theta$$

We can deal with

$$\int_0^{\pi/2} \int_0^{\frac{n}{2} \sec \theta} (\alpha - smx) \exp(\eta - n) r dr d\theta$$

Theorem 7 (Asymptotic upper bound on mean perimeter length). The mean perimeter excess length  $J_m$  is subject to the following asymptotic upper bound:

$$J_m \leq O(\log m) \quad \text{as } m \rightarrow \infty. \quad (10)$$

*Proof.* Without loss of generality, place the points  $v_i$  and  $v_j$  at  $(-\frac{m}{2}, 0)$  and  $(\frac{m}{2}, 0)$ . The double integral in (4) possesses mirror symmetry in each of the two axes, so we can write

$$\begin{aligned} J_m &= 2 \iint_{[0, m]^2} (\phi - \sin \phi) \exp(-\frac{1}{2}(v - m)) \text{Leb}(dz) \\ &= 2 \int_0^{\pi/2} \int_0^{\frac{m}{2} \sec \theta} (\phi - \sin \phi) \exp(-\frac{1}{2}(v - m)) r dr d\theta + \\ &\quad + 2 \int_{\pi/2}^{\pi} \int_0^m (\phi - \sin \phi) \exp(-\frac{1}{2}(v - m)) r dr d\theta \quad (11) \end{aligned}$$

(using polar coordinates  $(r, \theta)$  about the second point  $v_j$  located at  $(\frac{m}{2}, 0)$ ). The integrand in the second summand is dominated by  $\pi \exp(-\frac{1}{2}) r$ , which is integrable over  $(r, \theta) \in [0, \infty) \times [\frac{\pi}{2}, \pi]$ . (In this region geometry shows that  $v - m > r(1 - \cos \theta) \geq r$ .) Thus we can apply Lebesgue's dominated convergence theorem to deduce that the second summand is  $O(1)$  as  $m \rightarrow \infty$ , hence may be neglected.

In fact we can also show that part of the first summand generates an  $O(1)$  term: the dominated convergence theorem can be applied for any  $\epsilon \in [0, \pi/2]$  to show that

$$2 \int_0^{\pi/2} \int_{\epsilon}^{\frac{m}{2} \sec \theta} (\phi - \sin \phi) \exp(-\frac{1}{2}(v - m)) r dr d\theta = O(1),$$

since the integrand is dominated by  $\pi \exp(-\frac{1}{2}(1 - \cos \theta)) r$  over the region  $(r, \theta) \in (0, \infty) \times (\epsilon, \frac{\pi}{2})$  (in this region geometry shows that  $v - m > r(1 - \cos \theta) > r(1 - \cos \epsilon)$ ). Thus for fixed  $\epsilon \in (0, \frac{\pi}{2})$  as  $m \rightarrow \infty$  we have the asymptotic expression

$$J_m = 2 \int_0^{\pi} \int_0^{\frac{m}{2} \sec \theta} (\phi - \sin \phi) \exp(-\frac{1}{2}(v - m)) r dr d\theta + O(1).$$

Now in the region  $(r, \theta) \in (0, \infty) \times [0, \epsilon]$  we know  $\phi < 2\theta < 2\epsilon$ , and moreover  $\phi - \sin \phi$  is an increasing function of  $\phi$  (so long as  $\epsilon < \frac{\pi}{2}$ ). Therefore there is a constant  $C_\epsilon$  such that

$$\phi - \sin \phi \leq 2\theta - \sin(2\theta) \leq \frac{C_\epsilon}{8} \frac{(2\theta)^2}{6} \leq C_\epsilon \frac{1 - \cos \theta}{3} \sin \theta.$$

Hence

$$\begin{aligned} &2 \int_0^{\pi} \int_0^{\frac{m}{2} \sec \theta} (\phi - \sin \phi) \exp(-\frac{1}{2}(v - m)) r dr d\theta \\ &\leq \frac{4C_\epsilon}{3} \int_0^{\pi} \int_0^{\frac{m}{2} \sec \theta} (1 - \cos \theta) \sin \theta \exp(-\frac{1}{2}(1 - \cos \theta)) r dr d\theta \\ &= \frac{4C_\epsilon}{3} \int_0^{\pi} \left( \int_0^{\frac{m}{2} \sec \theta - 1} e^{-v} dv \right) \frac{\sin \theta \, d\theta}{1 - \cos \theta} \quad (\text{using } v = \frac{m}{2}(m\cos \theta - 1)) \\ &\leq \frac{4C_\epsilon}{3} \int_0^{\frac{m}{2}(\sec \epsilon - 1)} \left( \int_0^v e^{-v} dv \right) \frac{1}{1 + 4\theta/m} \frac{dv}{\theta} \quad (\text{using } v = \frac{m}{2}(m\cos \theta - 1)) \\ &\leq \frac{4C_\epsilon}{3} \log(\frac{m}{2}(\sec \epsilon - 1)) + O(1). \end{aligned}$$

□

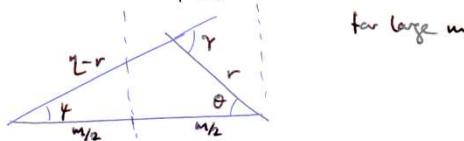
Actually

$$\frac{1}{2} \iint (x - smx) \exp\left(-\frac{1}{2}(y-n)\right) dx dy$$

$$= \frac{8}{3}(\log n + \text{constant})$$

and we can do  
higher-order asymptotics too.

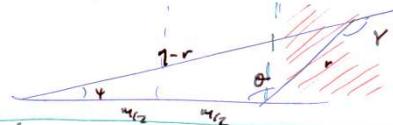
Asymptotic for  $J_m = 2 \int_{R_2}^{\infty} \int_0^{\pi} (\gamma - \sin \gamma) \exp(-\frac{1}{2}(\gamma - m)) r dr d\theta$



Using polar coordinates  $(r, \theta)$ , it is natural to split  $J_m$  into two double integrals:

$$J_m = 2 \int_0^{m/2} \int_0^{\pi/2} (\gamma - \sin \gamma) \exp(-\frac{1}{2}(\gamma - m)) r dr d\theta + 2 \int_{m/2}^{\pi} \int_0^{\infty} (\gamma - \sin \gamma) \exp(-\frac{1}{2}(\gamma - m)) r dr d\theta.$$

The second double integral is concerned with red-hatched region



In general we have  $0 < \gamma - \sin \gamma \geq \pi$  as  $\gamma \geq \pi$  (derivative is  $1 - \cos \gamma \geq 0$ ) and in the red-hatched region

$$\gamma - r > m$$

where for fixed  $r, \theta$  as  $m \rightarrow \infty$  so  $\gamma \rightarrow \theta$ ,  $\gamma - m \rightarrow r(1 - \cos \theta)$ . Hence the second double integral has integrand dominated by that in

$$2 \int_{m/2}^{\pi} \int_0^{\infty} \exp(-\frac{1}{2}) r dr d\theta < \infty.$$

Therefore we can apply Lebesgue's dominated convergence theorem:

$$\begin{aligned} & 2 \int_{m/2}^{\pi} \int_0^{\infty} (\gamma - \sin \gamma) \exp(-\frac{1}{2}(\gamma - m)) r dr d\theta \\ & \xrightarrow{m \rightarrow \infty} 2 \int_{m/2}^{\pi} \int_0^{\infty} (\theta - \sin \theta) \exp(-\frac{1}{2}(1 - \cos \theta)) r dr d\theta \\ & = \frac{8}{3} \int_{m/2}^{\pi} \frac{\theta - \sin \theta}{(1 - \cos \theta)^2} d\theta = \boxed{\frac{8}{3} (\pi + \log 2 - 1)} \end{aligned}$$

[Mathematica calculation]

In this range  $0 < \theta - \sin \theta \leq \gamma - \sin \gamma$  (using increasing relation of  $\theta - \sin \theta$ ) while  $0 < r(1 - \cos \theta) \leq \gamma - m$ .

So now analyse the error for this asymptotic: it is bounded by

$$\begin{aligned} & \left| 2 \int_{m/2}^{\pi} \int_0^{\infty} (\gamma - \sin \gamma) \exp(-\frac{1}{2}(\gamma - m)) r dr d\theta \right. \\ & \quad \left. - 2 \int_{m/2}^{\pi} \int_0^{\infty} (\theta - \sin \theta) \exp(-\frac{1}{2}(1 - \cos \theta)) r dr d\theta \right| \\ & \leq 2 \int_{m/2}^{\pi} \int_0^{\infty} |(\gamma - \sin \gamma) - (\theta - \sin \theta)| \exp(-\frac{1}{2}(1 - \cos \theta)) r dr d\theta \\ & \quad + 2 \int_{m/2}^{\pi} \int_0^{\infty} (\theta - \sin \theta) (\exp(-\frac{1}{2}(1 - \cos \theta)) - \exp(-\frac{1}{2}(\gamma - m))) r dr d\theta. \end{aligned}$$

Consider the first of these double integrals.

While  $m/2 \leq \gamma \leq \pi$  we have  $(\theta - \sin \theta)' = 1 - \cos \theta \in [1, 2]$  and  $\theta - \sin \theta \in [m/2 - 1, \pi]$ .

Moreover while  $\psi = \theta - \sin \theta \in [0, \pi/2]$  we have  $(\tan \psi - \psi)' = \sec^2 \psi - 1 \geq 0$  so  $\psi \leq \tan \psi \leq \frac{\tan \theta}{m - \cos \theta} \leq \frac{\sin \theta}{m}$  (when here we use the fact that  $\cos < 0$  in this range).

So in this range  $(\gamma - \sin \gamma) - (\theta - \sin \theta) \leq \max_{m/2 \leq \gamma \leq \pi} (\theta - \sin \theta)' \times \leq 2 \tan \psi \leq 2 \frac{\sin \theta}{m}$

So the first integral is bounded above by

$$\begin{aligned} & \frac{4}{m} \int_{m/2}^{\pi} \int_0^{\infty} \sin \theta \exp(-\frac{1}{2}(1 - \cos \theta)) r^2 dr d\theta \\ & \leq \frac{4}{m} \int_{m/2}^{\pi} \int_0^{\infty} \sin \theta e^{-r/2} r^2 dr d\theta \\ & = \frac{4}{m} \left( \sqrt{\pi/2} \sin \theta \right) \left( \int_0^{\infty} e^{-r/2} r^2 dr \right) \\ & = \frac{4}{m} \times 1 \times \frac{\pi}{8} = \boxed{\frac{8\pi}{64}} \end{aligned}$$

(In fact one could be precise here: the integral equals

$$\frac{64\pi}{m} \int_{m/2}^{\pi} \frac{\sin \theta}{(1 - \cos \theta)^2} d\theta = \frac{16\pi}{m}$$

written  
6.7.07  
included  
with Mma

Use the fact,  $\ln a \leq \ln a^2/4$  so  $\sqrt{\ln a} \leq \ln a^{1/2}$  if  $a \geq 1$ .  
Also note  $(Y-u) = r(1-\cos\theta)$   
 $= \sqrt{(m-r\cos\theta)^2 + r^2 \sin^2\theta} - (m-r\cos\theta)$   
 $\leq (m-r\cos\theta) \cdot \frac{1}{2} \frac{r^2 \sin^2\theta}{(m-r\cos\theta)^2} = \frac{r^2 (1-\cos^2\theta)}{2(m-r\cos\theta)}$ .

The second integral is bounded above by

$$\begin{aligned} & 2 \int_{\pi/2}^{\pi} \int_0^{\infty} (R - r\cos\theta) \exp(-\frac{s}{2}(1-\cos\theta)) (1 - \exp(-\frac{r^2(1-\cos^2\theta)}{4(m-r\cos\theta)})) r dr d\theta \\ & \leq 2\pi \int_{\pi/2}^{\pi} \int_0^{\infty} \exp(-\frac{s}{2}(1-\cos\theta)) (1 - \exp(-\frac{r^2 \sin^2\theta}{4m})) r dr d\theta \quad (\text{since } \theta < 0 \text{ without loss!}) \\ & \leq 2\pi \int_{\pi/2}^{\pi} \int_0^{\infty} \exp(-\frac{s}{2}(1-\cos\theta)) \frac{\sin^2\theta}{4m} r^3 dr d\theta \quad \theta - \pi/2 \leq 0 \\ & \quad + 2\pi \int_{\pi/2}^{\pi} \int_{r_m(\theta)}^{\infty} \exp(-\frac{s}{2}(1-\cos\theta)) r dr d\theta \quad (1-e^{-u} \leq u) \\ & \quad \text{factor } \frac{\sin^2\theta}{4m} r_m(\theta)^2 \leq 1 \text{ so } r_m(\theta) = \frac{2\sqrt{m}}{\sin\theta} \\ & = \frac{\pi}{2m} \int_{\pi/2}^{\pi} \int_0^{\infty} \exp(-\frac{s}{2}(1-\cos\theta)) r^3 dr \sin^2\theta d\theta \\ & \quad + 2\pi \int_{\pi/2}^{\pi} \int_{r_m(\theta)}^{\infty} \exp(-\frac{s}{2}(1-\cos\theta)) r dr d\theta \\ & = \frac{\pi}{2m} \int_{\pi/2}^{\pi} \frac{16}{(1-\cos\theta)^4} \sin^2\theta d\theta \int_0^{\infty} e^{-u} u^3 du \\ & = \frac{64\pi}{m} \int_{\pi/2}^{\pi} \frac{\sin^2\theta d\theta}{(1-\cos\theta)^4} \quad \therefore \boxed{\frac{64\pi}{5} \frac{1}{m}} \quad [\text{numerically calculated}] \end{aligned}$$

Then we have established

$$J_m = 2 \int_{\pi/2}^{\pi} \int_0^{\infty} (Y - u) \exp(-\frac{s}{2}(Y-u)) r dr d\theta \\ + \frac{8}{3} (\pi + \log 2 - 1) + O\left(\frac{1}{m}\right)$$

where in fact the  $O\left(\frac{1}{m}\right)$  term is bounded in absolute value by

$$\frac{2\pi}{m} + \frac{64\pi}{5} \frac{1}{m} = \boxed{\left(3 + \frac{8}{5}\pi\right) \frac{2\pi}{m}}.$$

numerically  
6.7.07

So now we must control the second double integral

$$\int_0^{\pi/2} \int_0^{\infty} f_m(\theta, r) r dr d\theta \quad \text{when } f_m(\theta, r) = 2(Y - u) \exp(-\frac{s}{2}(Y-u)).$$

First of all, consider the approximation derived by replacing  $f_m(\theta, r)$  with  $g(\theta, r) = \frac{2}{3}(1-\cos\theta) \sin\theta \exp(-\frac{s}{2}(1-\cos\theta))$ .

(Motivation: agrees in first factor to  $O(\theta^3)$ , in exponent and as limit, when  $\theta$  is small and  $m$  is large!)

$$\begin{aligned} & \frac{2}{3} \int_0^{\pi/2} \int_0^{\infty} (1-\cos\theta) \sin\theta \exp(-\frac{s}{2}(1-\cos\theta)) r dr d\theta \\ & = \frac{8}{3} \int_0^{\pi/2} \int_0^{\infty} e^{-s(\cos\theta-1)} e^{-s} s ds \frac{\sin\theta}{1-\cos\theta} d\theta \quad \text{when } s = \frac{2}{3}(1-\cos\theta) \\ & \quad \text{so } ds = \frac{2}{3}(1-\cos\theta) d\theta \\ & = \frac{8}{3} \int_0^{\infty} \int_0^v e^{-s} s ds \frac{1+4v/m}{4v/m} \frac{4}{m} \frac{dv}{(1+4v/m)^2} \quad \text{when } v = \frac{m}{4}(\cos\theta-1) \\ & = \frac{8}{3} \int_0^{\infty} \int_0^v e^{-s} s ds \left( \frac{1}{4v/m} - \frac{1}{1+4v/m} \right) \frac{4}{m} dv \quad (partial fraction) \\ & \quad \text{(numerically part)} \\ & = \frac{8}{3} \int_0^v e^{-s} s ds \left( \log \frac{4v}{m} - \log \left(1 + \frac{4v}{m}\right) \right) \Big|_0^\infty \quad (integrated} \rightarrow 0 \text{ as } v \rightarrow \infty \\ & \quad - \frac{8}{3} \int_0^\infty e^{-v} v \left( \log \frac{4v}{m} - \log \left(1 + \frac{4v}{m}\right) \right) dv \\ & = \frac{8}{3} \int_0^\infty e^{-s} s ds \log \frac{4v}{m} \\ & \quad - \frac{8}{3} \int_0^\infty e^{-v} v \log \frac{4v}{m} dv \\ & \quad + \frac{8}{3} \int_0^\infty e^{-v} v \log \left(1 + \frac{4v}{m}\right) dv \\ & = \frac{8}{3} \left( \log \frac{4v}{m} - 1 + \gamma \right) \\ & \quad - \frac{8}{3} \int_0^\infty (1+v) e^{-v} \log \left(1 + \frac{4v}{m}\right) dv + \frac{8}{3} \int_0^\infty e^{-v} \frac{4v}{1+4v/m} \frac{4}{m} dv \quad \text{integrated by parts again} \\ & = \frac{8}{3} \left( \log \frac{4v}{m} - 1 + \gamma \right) + \frac{32}{3m} \int_0^\infty e^{-v} \frac{4v}{1+4v/m} dv. \quad \text{use } \int_0^\infty e^{-v} \log v dv = 1 - \gamma \text{ for } \gamma \text{ the Euler-Mascheroni constant} \end{aligned}$$

and consequently

$$\begin{aligned}
 & \frac{8}{3} \int_0^\infty \int_0^v (e^{-s} - e^{-\frac{1}{2}(y-u)}) s ds \frac{1}{1+4uv} \frac{du}{v} \\
 &= \frac{8}{3} \int_0^\infty \int_0^v (1 - \exp(-\frac{m}{2}(1-\frac{2}{v})) (\sqrt{1 + \frac{(s/2v)^2}{(1-s/2v)^2}} \frac{4v}{m} (2+\frac{4v}{m}) - 1)) \\
 &\quad \leftarrow \boxed{\text{see Ch 10 ex 1}}
 \end{aligned}$$

Now the integrand is dominated by the integrand on the LHS.

$$\begin{aligned}
 & \frac{8}{3} \int_0^\infty \int_0^v (1 - \exp(-\frac{1}{4} \frac{(s/2v)^2}{1-s/2v} \cdot 4v (2+\frac{4v}{m}))) e^{-s} s ds \frac{du}{v} \\
 &\rightarrow \frac{8}{3} \int_0^\infty \int_0^v (1 - \exp(-\frac{s^2}{2v-s})) e^{-s} s ds \frac{du}{v} \\
 &\leq \frac{8}{3} \int_0^\infty \int_0^1 (1 - \exp(-\frac{t^2}{2-t} v)) e^{-vt} v dt ds \quad \boxed{(s=vt)} \\
 &= \frac{8}{3} \int_0^1 \int_0^\infty (1 - \exp(-\frac{t^2}{2-t} v)) e^{-tv} t^2 v dv \frac{dt}{t} \\
 &= \frac{8}{3} \int_0^1 \int_0^\infty (e^{-tv} - \exp(-\frac{2}{2-t} tv)) t^2 v dv \frac{dt}{t} \\
 &= \frac{8}{3} \int_0^1 (1 - (\frac{2-t}{2})^2) \frac{dt}{t} = \frac{8}{3} \int_0^1 (\frac{4t-4}{4}) \frac{dt}{t} \\
 &= \frac{8}{3} \int_0^1 (1 - t \cdot \frac{1}{4}) dt = \frac{8}{3} (1 - \frac{1}{8}) = \frac{7}{3} < \infty.
 \end{aligned}$$

However the dominated integrand is also the limit as  $m \rightarrow \infty$  with  $s, v$  (or  $t, v$ ) held fixed! We deduce

$$\begin{aligned}
 & \frac{8}{3} \int_0^{\pi/2} \int_0^{m \sec \theta} (1 - \cos \theta) \sin \theta (\exp(-\frac{r}{2}(\cos \theta)) - \exp(-\frac{r}{2}(y-u))) r dr d\theta \\
 &\rightarrow \boxed{\frac{7}{3}} \text{ as } m \rightarrow \infty.
 \end{aligned}$$

There are perfect for extracting error bounds, using  $\sqrt{1+a} \geq 1 + \frac{a}{2} - \frac{a^2}{8}$ ; however for now we turn to the other double integral!

Consider first  $2 \int_0^{\pi/2} \int_0^{m \sec \theta} ((Y - \sin Y) - \frac{(1 - \cos \theta) \sin \theta}{3}) \exp(-\frac{r}{2}(y-u)) r dr d\theta$  for fixed  $z > 0$ .

The integrand here is dominated by the integrand in

$$\begin{aligned}
 & 2 \left( \frac{\pi}{2} - 1 + \frac{1}{3} \right) \int_0^{\pi/2} \int_0^\infty \exp(-\frac{r}{2}(1 - \cos \theta)) r dr d\theta \\
 &= \left( \pi - \frac{4}{3} \right) \cdot 4 \int_0^{\pi/2} \frac{d\theta}{(1 - \cos \theta)^2} \\
 &= 4 \left( \pi - \frac{4}{3} \right) \left( \frac{(2 - \cos(\theta)) \sin(\theta)}{3(1 - \cos(\theta))^2} - \frac{2}{3} \right) < \infty.
 \end{aligned}$$

Hence we may apply the dominated convergence theorem:

The double integral in question converges to a constant which is approximated for small  $\varepsilon > 0$  by

$$\begin{aligned}
 & 2 \int_0^{\pi/2} \int_0^\infty ((\theta - \sin \theta) - \frac{(1 - \cos \theta) \sin \theta}{3}) \exp(-\frac{r}{2}(1 - \cos \theta)) r dr d\theta \\
 &= \frac{8}{3} \int_0^{\pi/2} \frac{(3(\theta - \sin \theta) - (1 - \cos \theta) \sin \theta)}{(1 - \cos \theta)^2} d\theta \quad \boxed{1.7707} \\
 &= \boxed{\frac{8}{3} \left( \frac{11}{3} - \pi \right)} \quad \boxed{\text{numerical calculation}}
 \end{aligned}$$

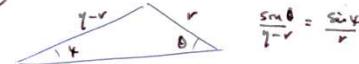
So the final task is to establish control of

$$2 \int_0^{\pi/2} \int_0^{m \sec \theta} ((Y - \sin Y) - \frac{(1 - \cos \theta) \sin \theta}{3}) \exp(-\frac{r}{2}(y-u)) r dr d\theta.$$

Here is a key argument:

$$\begin{aligned}
 Y - \sin Y &= \frac{Y^3}{6} + O(Y^5) = \frac{Y^3}{6} + O(\theta^5) \quad \text{as } 2\theta \geq Y \\
 &= \frac{\theta^3}{6} (1 + \frac{Y}{\theta})^3 + O(\theta^5) \quad \text{as } Y = \theta + r \\
 &= \frac{\theta^3}{6} (1 + \frac{\sin \theta + r}{\theta + r})^3 + O(\theta^5) \quad \text{as } Y < \theta, \sin \theta = \theta + O(\theta^3), \\
 &\quad \sin \theta = \theta + O(\theta^3) \\
 &= \frac{\theta^3}{6} (1 + \frac{r}{\theta + r})^3 + O(\theta^5) \\
 &= \frac{\theta^3}{6} \left( \frac{q}{q+r} \right)^3 + O(\theta^5) \\
 &= \frac{(1 - \cos \theta) \sin \theta}{3} \left( \frac{q}{q+r} \right)^3 + O(\theta^5) \quad \text{using } \theta < \frac{q}{q+r} \leq 1 + \frac{r}{q+r} \\
 &\quad \text{no uniformly!} \quad \leq 1 + \frac{m \sec \theta}{2(q+r)} \leq m \sec \theta
 \end{aligned}$$

overleaf we control result of  $O(\theta^5)$



$$\begin{aligned} \text{so } \frac{r}{1-r} &= \frac{1+4\sqrt{m}}{2\sqrt{v}} s / \sqrt{\left((1-\frac{s}{2\sqrt{v}})^2 + \left(\frac{s}{2\sqrt{v}}\right)^2 \frac{4v}{m} (2+\frac{4v}{m})\right)} \\ &\rightarrow \frac{\frac{s}{2\sqrt{v}}}{\sqrt{\left(1-\frac{s}{2\sqrt{v}}\right)^2}} = \frac{s}{2\sqrt{v-s}} \quad \checkmark \\ \frac{1}{1-r} &\rightarrow \frac{2\sqrt{v}}{2\sqrt{v-s}} \quad \checkmark \\ \frac{m}{1-r} &= \frac{1}{\sqrt{\left((1-\frac{s}{2\sqrt{v}})^2 + \left(\frac{s}{2\sqrt{v}}\right)^2 \frac{4v}{m} (2+\frac{4v}{m})\right)}} \rightarrow \frac{2\sqrt{v}}{2\sqrt{v-s}} \quad \checkmark \end{aligned}$$

Hence

$$\begin{aligned} &2 \int_0^{\infty} \int_0^{\frac{\pi}{2}} \sin^2 \theta \left( (\gamma - \sin \theta) - \left(1 - \cos \theta\right) \sin^2 \theta \right) \exp\left(-\frac{1}{2}(q-u)\right) r dr d\theta \\ &\rightarrow \frac{4}{3} \int_0^{\infty} \int_0^v \frac{2\sqrt{v}}{2\sqrt{v-s}} \left(1 + \frac{2\sqrt{v}}{2\sqrt{v-s}} + \left(\frac{2\sqrt{v}}{2\sqrt{v-s}}\right)^2\right) e^{-s} \\ &\quad \exp\left(-\frac{s^2}{2\sqrt{v-s}}\right) s^2 ds \frac{dv}{\sqrt{v}} \quad \checkmark \\ &= \frac{4}{3} \int_0^{\infty} \int_0^v \frac{2\sqrt{v}}{2\sqrt{v-s}} \left(1 + \frac{2\sqrt{v}}{2\sqrt{v-s}} + \left(\frac{2\sqrt{v}}{2\sqrt{v-s}}\right)^2\right) \exp\left(-\frac{2vs}{2\sqrt{v-s}}\right) s^2 ds \frac{dv}{\sqrt{v}} \\ &\quad \text{use } s=vt \text{ as above} \\ &= \frac{4}{3} \int_0^{\infty} \int_0^1 \frac{2}{2-t} \left(1 + \frac{2}{2-t} + \left(\frac{2}{2-t}\right)^2\right) \exp\left(-\frac{2t}{2-t} v\right) v t^2 dt dv \quad \checkmark \\ &= \frac{4}{3} \int_0^1 \int_0^{\infty} \exp\left(-\frac{2t}{2-t} v\right) v dv \frac{2t^2}{2-t} \left(1 + \frac{2}{2-t} + \left(\frac{2}{2-t}\right)^2\right) dt \\ &= \frac{2}{3} \int_0^1 (2-t) \left[1 + \frac{2}{2-t} + \left(\frac{2}{2-t}\right)^2\right] dt \\ &= \frac{2}{3} \left( \int_0^1 (2-t) dt + 2 \int_0^1 dt + 4 \int_0^1 \frac{dt}{2-t} \right) \\ &= \frac{2}{3} \left( 2 - \frac{1}{2} + 2 + 4 \left[ \log \frac{1}{2-t} \right]_0^1 \right) \\ &= \frac{2}{3} \left( 2 - \frac{1}{2} + 2 + 4 \log 2 \right) \sim \boxed{\frac{2}{3} \left( \frac{7}{2} + 4 \log 2 \right)}. \end{aligned}$$

Putting this all together we find

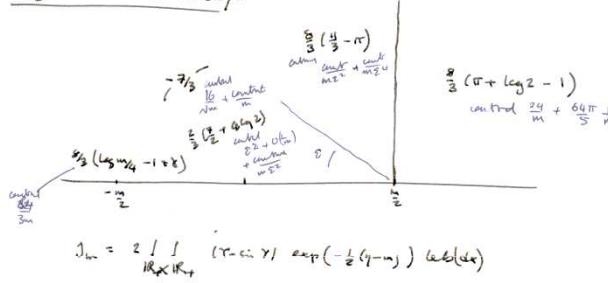
$$\begin{aligned} J_m &\sim \frac{8}{3} (\pi + \log 2 - 1) + \frac{2}{3} (\log \frac{m}{4} - 1 + 8) - \frac{2}{3} + \frac{8}{3} \left( \frac{11}{2} - \pi \right) \\ &\quad + \frac{2}{3} \left( \frac{7}{2} + 4 \log 2 \right) \end{aligned}$$

in the slightly extended sense that we get within  $\varepsilon$  of this as an upper bound!

Hence we can approximate  $J_m$  by

$$\boxed{\frac{8}{3} (\log m + \gamma + \frac{5}{3})} \quad \text{for large } m$$

Diagrammatic summary:



$$J_m = 2 \int_{R \times R^+} (\gamma - \sin y) \exp\left(-\frac{1}{2}(q-u)\right) u du dy$$

**End of awful  
example**

What  
went  
wrong?

How  
to do  
better?