

Neutron transport from a probabilistic perspective

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University of Bath

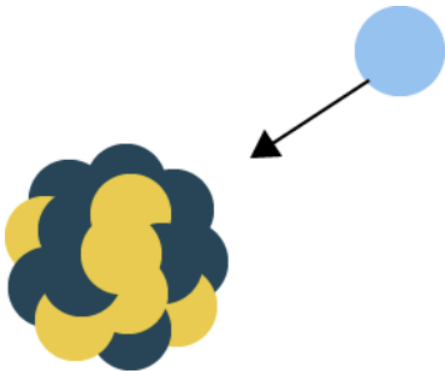
Contents

- 1 Neutron transport
- 2 Neutron transport equation
- 3 Neutron branching process
- 4 Monte Carlo

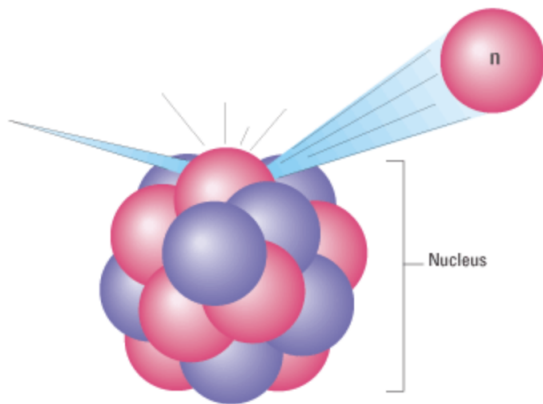
Neutron transport

- Neutron transport pertains to the study of the dynamics of neutrons and how they interact with surrounding material.
- Its origins date back to the 1800s when the Boltzmann equation was used to study the kinetic theory of gases.
- Large scale development of nuclear reactors during and after WWII caused an increase in focus on the mathematics of neutron transport.

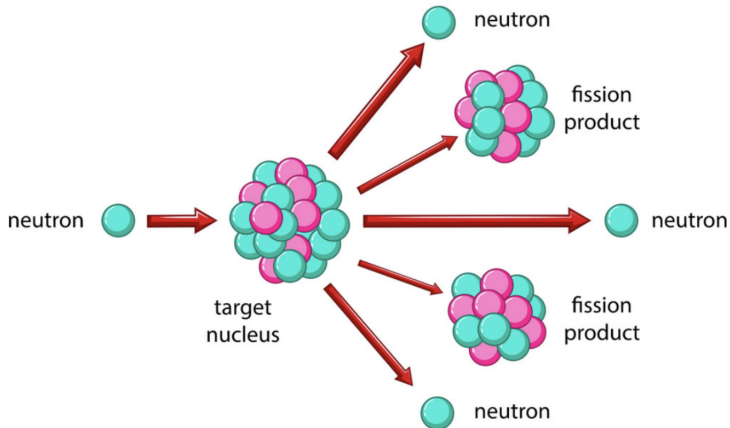
Neutron transport



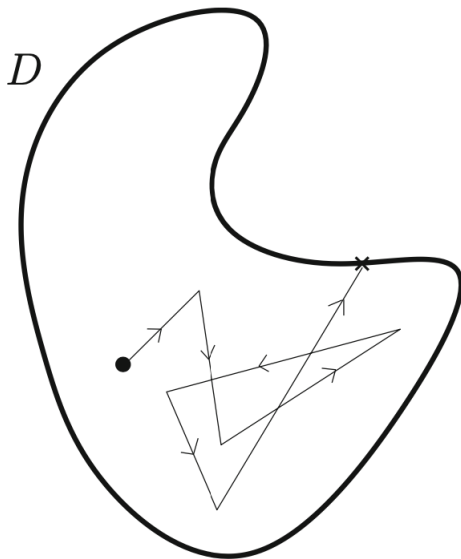
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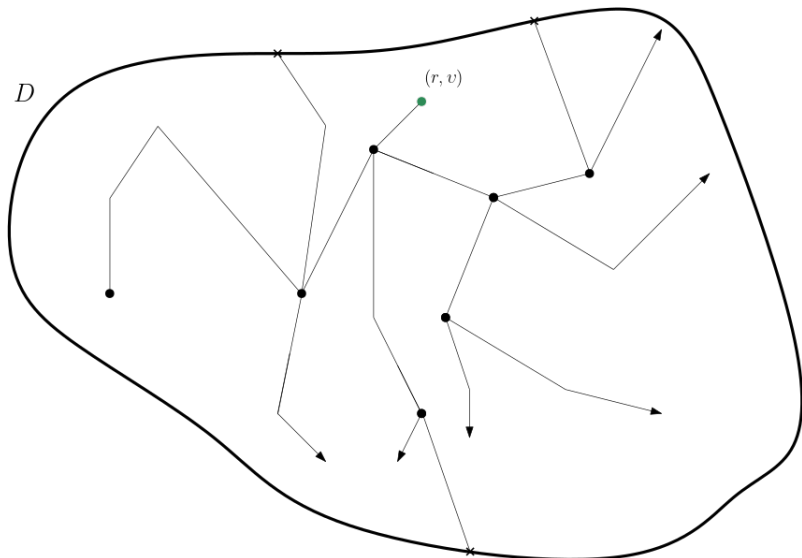
Neutron transport



Neutron transport



Neutron transport



- Reactor physics

- Understanding neutron interactions with the material.
- Mathematical modelling, simulations, thermal hydraulics, ...

- Shielding

- What proportion of neutrons will reach a certain area of the reactor?
- Rare event analysis, large deviations, sensitivity analysis, ...

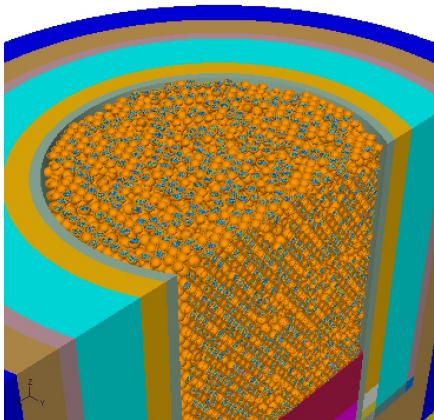
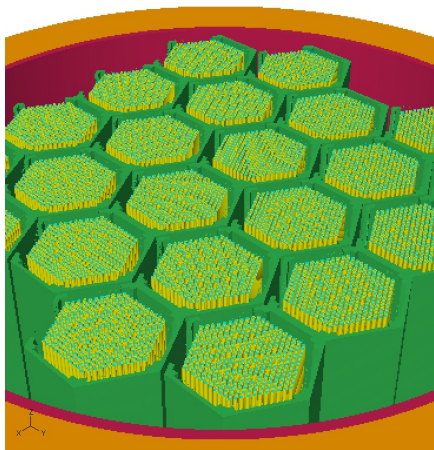
- Decommissioning

- Infer neutron sources from boundary readings.
- Inverse problems, constrained optimisation, non-linear programming, ...

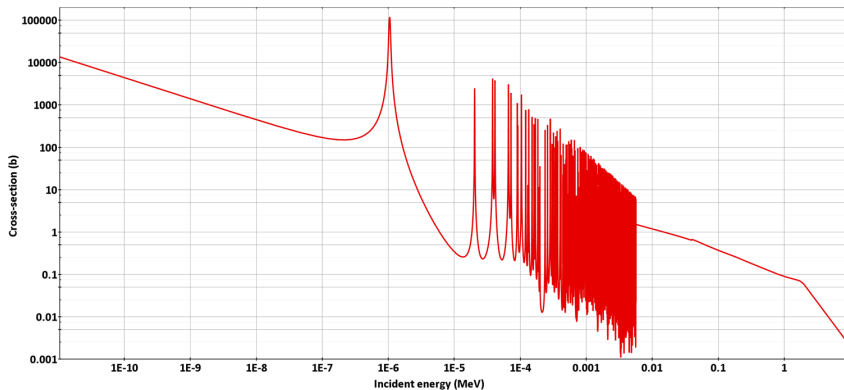
- Criticality

- Asymptotic growth rate of number of neutrons and their spatial distribution.
- PDEs, branching processes, stability analysis, ...

Neutron transport



Neutron transport



- We want to use mathematics to come up with clever ways of analysing and simulating these systems.
- For the rest of the talk, we will focus on criticality problems.
- Our aim will be to use tools from probability theory to understand the asymptotic behaviour of neutrons undergoing nuclear fission.

Contents

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- 2 Neutron transport equation
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- Let $D \subset \mathbb{R}^3$ be a open and bounded.
- Let $V := \{v \in \mathbb{R}^3 : v_{min} \leq |v| \leq v_{max}\}$, $0 < v_{min} \leq v_{max} < \infty$.
- If a particle has label $(r, v) \in D \times V$, we mean that the particle is at position r with velocity v .
- Let $\psi_t(r, v)$ denote the density of neutrons at position r with velocity v at time $t \geq 0$.

- Let $\sigma_s(r, v)$ denote the rate at which a scattering event occurs for a particle with configuration (r, v) . More precisely,

$$P_{(r_0, v_0)}(T_s > t) = e^{-\int_0^t \sigma_s(r_0 + v_0 u, v_0) du}.$$

- Given a scattering event occurs at r with incoming velocity v , $\pi_s(r, v, v')$ is the probability that the outgoing velocity is v' .
- $\sigma_f(r, v)$ denotes the rate at which a fission event occurs.
- Given a fission event occurs at r with incoming velocity v , $\pi_f(r, v, v')$ is the average number of neutrons produced in a fission event with outgoing new velocity v' .

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Neutron transport equation

$$\begin{aligned}\frac{\partial \psi_t}{\partial t}(r, v) &= v \cdot \nabla \psi_t(r, v) \\ &+ \sigma_s(r, v) \left(\int_V \psi_t(r, v') \pi_s(r, v, v') dv' - \psi_t(r, v) \right) \\ &+ \sigma_f(r, v) \left(\int_V \psi_t(r, v') \pi_f(r, v, v') dv' - \psi_t(r, v) \right), \\ &= (T + S + F)[\psi_t](r, v)\end{aligned}$$

Neutron transport equation

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We also impose the following boundary and initial conditions:

- $\psi_t(r, v) = 0, \quad r \in \partial D, \mathbf{n}_r \cdot v > 0,$
- $\psi_0(r, v) = g(r, v)$

Eigenvalue problem

Want to find $\lambda \in \mathbb{R}$ and a function $\varphi : D \times V \rightarrow [0, \infty)$ such that

$$(T + S + F)\varphi = \lambda\varphi.$$

Eigenvalue problem

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- λ tells us about the asymptotic growth rate (per unit time) of the number of neutrons.
- φ tells us which regions are favourable.

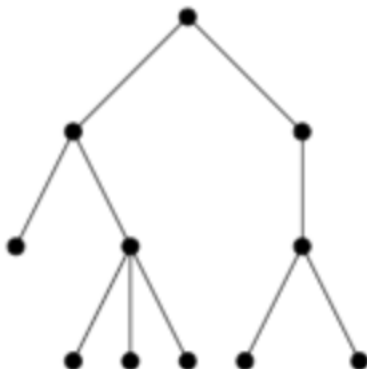
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- The simplest example of a branching process is the Bienaymé-Galton-Watson (BGW) process.
- The BGW process is a stochastic process $(Z_n)_{n \geq 0}$ where Z_n represents the number of individuals in the n -th generation of a randomly evolving population that undergoes asexual reproduction.
- The process evolves according to the recurrence formula: $Z_0 = 1$ and for $n \geq 0$,

$$Z_{n+1} = \sum_{i=1}^{Z_n} \xi_i^{(n)}, \quad \xi_i^{(n)} \overset{i.i.d.}{\sim} \xi.$$

BGW process



$$Z_0 = 1$$

$$Z_1 = 2$$

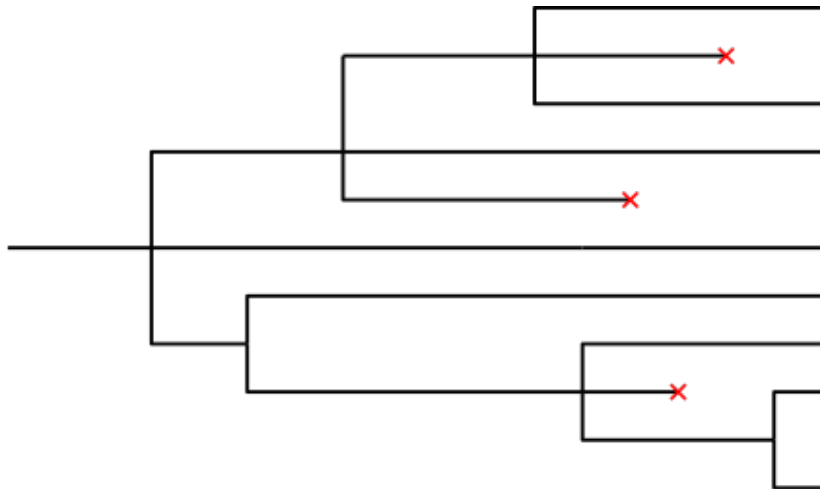
$$Z_2 = 3$$

$$Z_3 = 5$$

Continuous time BGW process

- The continuous time BGW process is a stochastic process $(Z_t)_{t \geq 0}$, where Z_t represents the number of individuals alive at time t in a randomly evolving population that undergoes asexual reproduction.
- With each individual, we associate an independent $Exp(1)$ random variable (clock). When an individual's clock rings, they give birth to a random number of individuals $\overset{i.i.d.}{\sim} \xi$.

BGW process



Neutron branching process

- Let $\{(r_i(t), v_i(t)) : i = 1, \dots, N_t\}$ denote the collection of particles alive at time $t \geq 0$ in $D \times V$.

- Define

$$\phi_t[g](r, v) = \mathbb{E}_{\delta_{(r,v)}} \left[\sum_{i=1}^{N_t} g(r_i(t), v_i(t)) \right].$$

- E.g., taking $g = \mathbf{1}_{D \times V}$, $\phi_t[g](r, v) = \mathbb{E}_{\delta_{(r,v)}}[N_t]$.
Similarly taking $g = \mathbf{1}_{D_\epsilon \times V}$, $\phi_t[g](r, v) = \mathbb{E}_{\delta_{(r,v)}}[\# \text{ particles in } D_\epsilon]$.

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Theorem

$\phi_t[g](r, v)$ “solves” the neutron transport equation with initial condition $\psi_0 = g$,
i.e.

$$\frac{\partial \phi_t}{\partial t}(r, v) = (T + S + F)\phi_t[g](r, v).$$

Neutron branching process

The proof uses generalisations of the following facts:

- If $T_i \sim \text{Exp}(\lambda_i)$, $i = 1, 2$ are independent, then $\min\{T_1, T_2\} \sim \text{Exp}(\lambda_1 + \lambda_2)$.
- $P(\min\{T_1, T_2\} = T_i) = \lambda_i/(\lambda_1 + \lambda_2)$.
- $\mathbb{E}[X] = \sum_i \mathbb{E}[X|A_i]\mathbb{P}(A_i)$.

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In our case,

- $T_1 = \text{first scatter time}$, $T_2 = \text{first fission time}$.
- X is replaced by $\sum_{i=1}^{N_t} f(r_i(t), v_i(t))$.
- $A_1 = \{\min\{T_1, T_2\} > t\}$.
- $A_2 = \{\min\{T_1, T_2\} \leq t, \min\{T_1, T_2\} = T_1\}$.
- $A_3 = \{\min\{T_1, T_2\} \leq t, \min\{T_1, T_2\} = T_2\}$.

Neutron branching process

- By showing ϕ_t “solves” the NTE, we’ve justified working with the branching process to inform us about solutions to the NTE.
- Why is this useful?
- Returning to the BGW process, if $m = E[\xi]$, we can see that $E[Z_n] = m^n$.
- From this, we have three cases: $m > 1$, $m = 1$ and $m < 1$.
- The continuous time version is $E[Z_t] = e^{(m-1)t}$ and we have the same three cases.

Theorem

There exists $\lambda \in \mathbb{R}$ and a function $\varphi : D \times V \rightarrow [0, \infty)$ such that

$$\phi_t[\varphi](r, v) = e^{\lambda t} \varphi(r, v)$$

and, for any bounded function g , there exists a constant $C_g > 0$ such that

$$\phi_t[g](r, v) \sim C_g e^{\lambda t} \varphi(r, v), \quad t \rightarrow \infty.$$

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- Monte Carlo was first suggested by Stanisław Ulam who was attempting to calculate the probability of winning a game of solitaire.
- The method was developed during World War II as part of the Manhattan project to simulate neutron diffusion paths for the hydrogen bomb.

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Monte Carlo for neutron transport

- Recall the asymptotic

$$\phi_t[g](r, v) \sim C_g e^{\lambda t} \varphi(r, v).$$

- Setting $g = \mathbf{1}_{D \times V}$, we have $C_g = 1$ and hence

$$\mathbb{E}_{\delta_{(r,v)}}[N_t] \sim e^{\lambda t} \varphi(r, v).$$

- Manipulating this gives

$$\lambda = \lim_{t \rightarrow \infty} \frac{1}{t} \log \mathbb{E}_{\delta_{(r,v)}}[N_t] \approx \frac{1}{T} \log \frac{1}{N} \sum_{i=1}^N N_T^{(i)}.$$

and

$$\varphi(r, v) = \lim_{t \rightarrow \infty} e^{-\lambda t} \mathbb{E}_{\delta_{(r,v)}}[N_t] \approx e^{-\hat{\lambda} t} \frac{1}{N} \sum_{i=1}^N N_T^{(i)}.$$

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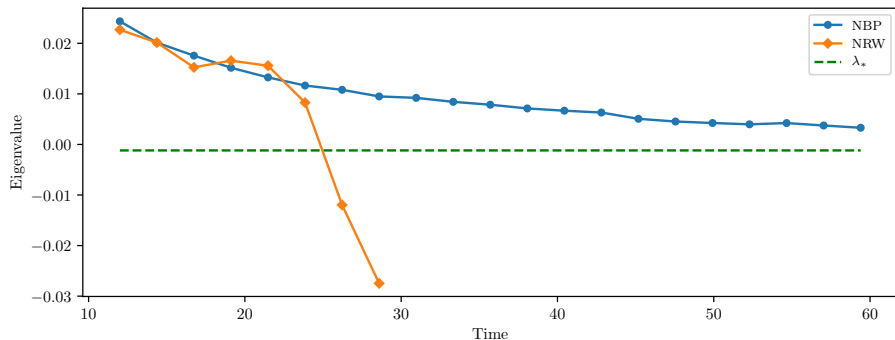
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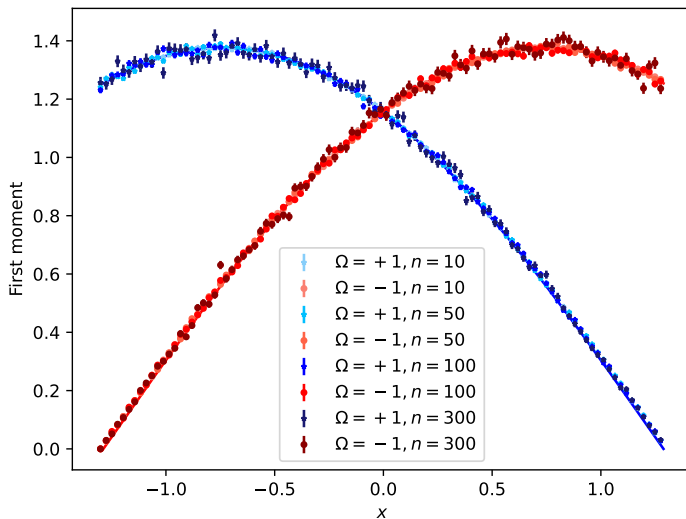
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- $D = (-L, L)$, $V = \{-1, +1\}$.
- Fission occurs at a constant rate:
 - 0 offspring w.p. p_0 ,
 - 2 offspring w.p. p_2 , isotropic velocities.
- Scattering occurs at a constant rate and is isotropic.

Monte Carlo for neutron transport



Simulations



Other considerations

What about a left eigenfunction?

Other considerations

What about a left eigenfunction?

Theorem

There exists $\lambda \in \mathbb{R}$ and functions $\varphi, \tilde{\varphi} : D \times V \rightarrow [0, \infty)$ such that

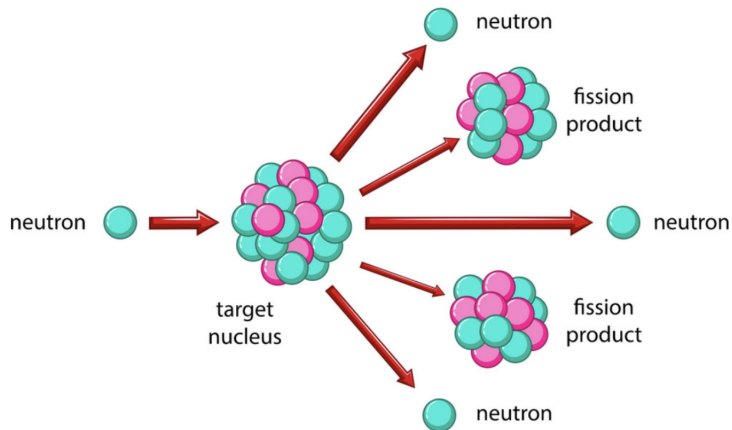
$$\phi_t[\varphi](r, v) = e^{\lambda t} \varphi(r, v), \quad \langle \tilde{\varphi}, \phi_t[g] \rangle = e^{\lambda t} \langle \tilde{\varphi}, g \rangle,$$

and, for any bounded function g ,

$$\phi_t[g](r, v) \sim e^{\lambda t} \varphi(r, v) \langle g, \tilde{\varphi} \rangle, \quad t \rightarrow \infty.$$

Here, $\langle f, g \rangle = \int_{D \times V} f(r, v) g(r, v) dr dv$.

Other considerations



Other considerations

- Suppose I choose a random variable but I only tell you its mean. How many random variables can you come up with that have the same mean?
- How does that apply in this setting? Why might that be useful?

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Other (related) application areas

- Radiotherapy
- Cosmic rays
- Fusion
- ...



Thank you!