# Oscillatory Attraction and Repulsion from a Subset of the Unit Sphere or Hyperplane for Isotropic Stable Lévy Processes 

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#### Abstract

Suppose that $S$ is a closed set of the unit sphere $\mathbb{S}^{d-1}=\left\{x \in \mathbb{R}^{d}:|x|=\right.$ $1\}$ in dimension $d \geq 2$, which has positive surface measure. We construct the law of absorption of an isotropic stable Lévy process in dimension $d \geq 2$ conditioned to approach $S$ continuously, allowing for the interior and exterior of $\mathbb{S}^{d-1}$ to be visited infinitely often. Additionally, we show that this process is in duality with the unconditioned stable Lévy process. We can replicate the aforementioned results by similar ones in the setting that $S$ is replaced by $D$, a closed bounded subset of the hyperplane $\left\{x \in \mathbb{R}^{d}:(x, v)=0\right\}$ with positive surface measure, where $v$ is the unit orthogonal vector and where $(\cdot, \cdot)$ is the usual Euclidean inner product. Our results complement similar results of the authors [17] in which the stable process was further constrained to attract to and repel from $S$ from either the exterior or the interior of the unit sphere.


Keywords Stable process • Time reversal • Duality

[^0]
## 1 Introduction

Let $X=\left(X_{t}, t \geq 0\right)$ be a $d$-dimensional stable Lévy process $(d \geq 2)$ with probabilities $\left(\mathbb{P}_{x}, x \in \mathbb{R}^{d}\right)$. This means that $X$ has càdlàg paths with stationary and independent increments, and there exists an $\alpha>0$ such that, for $c>0$, and $x \in \mathbb{R}^{d}$,
under $\mathbb{P}_{x}$ the law of $\left(c X_{c^{-\alpha}}, t \geq 0\right)$ is equal to $\mathbb{P}_{c x}$.
The latter is the property of so-called self-similarity. It turns out that stable Lévy processes necessarily have $\alpha \in(0,2]$. The case $\alpha=2$ is that of standard $d$ dimensional Brownian motion, thus has a continuous path. All other $\alpha \in(0,2)$ have no Gaussian component and are pure jump processes. In this article we are specifically interested in phenomena that can only occur when jumps are present. We thus restrict ourselves henceforth to the setting $\alpha \in(0,2)$.

Although Brownian motion is isotropic, this need not be the case in the stable case when $\alpha \in(0,2)$. Nonetheless, we will restrict to the isotropic setting. To be more precise, this means, for all orthogonal transformations $U: \mathbb{R}^{d} \mapsto \mathbb{R}^{d}$ and $x \in \mathbb{R}^{d}$,
the law of $\left(U X_{t}, t \geq 0\right)$ under $\mathbb{P}_{x}$ is equal to $\left(X_{t}, t \geq 0\right)$ under $\mathbb{P}_{U x}$.
For convenience, we will henceforth refer to $X$ as a stable process.
As a Lévy process, our stable process of index $(0,2)$ has a characteristic triplet $(0,0, \Pi)$, where the jump measure $\Pi$ satisfies

$$
\begin{equation*}
\Pi(B)=\frac{2^{\alpha} \Gamma((d+\alpha) / 2)}{\pi^{d / 2}|\Gamma(-\alpha / 2)|} \int_{B} \frac{1}{|y|^{\alpha+d}} \ell_{d}(\mathrm{~d} y), \quad B \subseteq \mathcal{B}\left(\mathbb{R}^{d}\right) \tag{1}
\end{equation*}
$$

where $\ell_{d}$ is $d$-dimensional Lebesgue measure. ${ }^{1}$ This is equivalent to identifying its characteristic exponent as

$$
\Psi(\theta)=-\frac{1}{t} \log \mathbb{E}\left(\mathrm{e}^{\mathrm{i} \theta \cdot X_{t}}\right)=|\theta|^{\alpha}, \quad \theta \in \mathbb{R}^{d}
$$

where we write $\mathbb{P}$ in preference to $\mathbb{P}_{0}$.
In this article, we characterise the law of a stable process conditioned to continuously approach a closed subdomain of the surface of a unit sphere, say $S \subseteq \mathbb{S}^{d-1}=\left\{x \in \mathbb{R}^{d}:|x|=1\right\}$, which has non-zero surface measure. Moreover,

[^1]our conditioning will allow the stable process to approach $S$ by visiting the exterior and interior of $\mathbb{S}^{d-1}$ infinitely often. We note that when $\alpha \in(1,2)$, stable processes will hit the unit sphere with positive probability and otherwise, when $\alpha \in(0,1]$ it hits the unit sphere with probability zero; see e.g. [25] or [16]. The aforesaid conditioning is thus only of interest when $\alpha \in(0,1]$.

In addition to constructing the conditioned process, we develop an expression for the law of the limiting point of contact on S. Moreover, we show that, when time reversed from the strike point on $S$, the resulting process can be described as nothing more than the stable process itself.

It turns out that the methodology we use here is robust enough to cover a similar suite of results for the case of an isotropic stable process conditioned to a closed subdomain of an arbitrary hyperplane in $\mathbb{R}^{d}$ that is orthogonal to an arbitrary unitlength vector $v \in \mathbb{R}^{d}$.

Our results naturally complement those of the recent paper [17], which considers a similar type of conditioning, albeit requiring the stable process to additionally remain either inside or outside of the unit ball. Other related works include [9] and [14], who considered a real valued stable process conditioned to hit 0 continuously and a real valued stable process conditioned to continuously approach the boundary of the interval $[-1,1]$ from the outside, respectively. In order to make our results pertinent, we restrict ourselves to the case that $d \geq 2$.

## 2 Oscillatory Attraction Towards S

Let $\mathbb{D}\left(\mathbb{R}^{d}\right)$ denote the space of càdlàg paths $\omega:[0, \infty) \rightarrow \mathbb{R}^{d} \cup \partial$ with lifetime $\zeta(\omega)=\inf \{s>0: \omega(s)=\partial\}$, where $\partial$ is a cemetery point. The space $\mathbb{D}\left(\mathbb{R}^{d}\right)$ will be equipped with the Skorokhod topology, with its closed $\sigma$-algebra $\mathcal{F}$ and natural filtration $\left(\mathcal{F}_{t}, t \geq 0\right)$. The reader will note that we will also use a similar notion for $\mathbb{D}(E)$ later on in this text in the obvious way for an $E$-valued Markov process. We will always work with $X=\left(X_{t}, t \geq 0\right)$ to mean the coordinate process defined on the space $\mathbb{D}\left(\mathbb{R}^{d}\right)$. Hence, the notation of the introduction indicates that $\mathbb{P}=\left(\mathbb{P}_{x}, x \in \mathbb{R}^{d}\right)$ is such that $(X, \mathbb{P})$ is our stable process.

We want to construct the law of the stable process conditioned to continuously limit to $S \in \mathbb{S}^{d-1}$ whilst visiting both $\mathbb{B}_{d}:=\left\{x \in \mathbb{R}^{d}:|x|<1\right\}$ and $\overline{\mathbb{B}}_{d}^{c}:=\mathbb{R}^{d} \backslash \overline{\mathbb{B}}_{d}$ infinitely often at arbitrarily small times prior to striking $S$. We shall denote the associated probabilities by $\mathbb{P}^{S}=\left(\mathbb{P}_{x}^{S}, x \in \mathbb{R}^{d} \backslash S\right)$. For a more precise definition of what is meant by this form of conditioning, let us introduce the stopping times,

$$
\begin{equation*}
\tau_{\beta}=\inf \left\{t>0: \beta^{-1}<\left|X_{t}\right|<\beta\right\}, \quad \text { for } \beta>1 \tag{2}
\end{equation*}
$$

Whenever it is well defined, we will write, for $t \geq 0, \Lambda \in \mathcal{F}_{t}$ and $x \notin \mathrm{~S}$,

$$
\begin{equation*}
\mathbb{P}_{x}^{\mathrm{S}}(\Lambda, t<\zeta)=\lim _{\beta \rightarrow 1} \lim _{\varepsilon \rightarrow 0} \mathbb{P}_{x}\left(\Lambda, t<\tau_{\beta} \mid \tau_{\mathbf{S}_{\varepsilon}}<\infty\right) \tag{3}
\end{equation*}
$$

where

$$
\tau_{\mathrm{S}_{\varepsilon}}=\inf \left\{t>0: X_{t} \in \mathrm{~S}_{\varepsilon}\right\} \quad \text { and } \quad \mathrm{S}_{\varepsilon}:=\left\{x \in \mathbb{R}^{d}: 1-\varepsilon \leq|x| \leq 1+\varepsilon \text { and } \arg (x) \in \mathrm{S}\right\} .
$$

Our first main result clarifies that the process $\left(X, \mathbb{P}^{\mathrm{S}}\right)$ is well defined. In the theorem below, and thereafter, we will understand $\sigma_{1}$ to mean the Lebesgue surface measure on $\mathbb{S}^{d-1}$ normalised to have unit mass, i.e. $\sigma_{1}\left(\mathbb{S}^{d-1}\right)=1$.
Theorem 1 Suppose that $\alpha \in(0,1]$ and the closed set $S \subseteq \mathbb{S}^{d-1}$ is such that $\sigma_{1}(S)>0$. For $\alpha \in(0,1]$, the limit (3) makes sense. Therefore, the process $\left(X, \mathbb{P}^{S}\right)$ is well defined such that

$$
\begin{equation*}
\left.\frac{\mathrm{d} \mathbb{P}_{x}^{S}}{\mathrm{~d} \mathbb{P}_{x}}\right|_{\mathcal{F}_{t}}=\frac{H_{S}\left(X_{t}\right)}{H_{S}(x)}, \quad t \geq 0, x \notin S \tag{4}
\end{equation*}
$$

where

$$
H_{S}(x)=\int_{S}|x-\theta|^{\alpha-d} \sigma_{1}(\mathrm{~d} \theta), \quad x \notin S
$$

Although excluded from the conclusion of Theorem 1, it is worth dwelling for a moment on the extreme case $S=\{\theta\}$, for $\theta \in \mathbb{S}^{d-1}$. It has been shown in [20] that, when $\alpha \in(0,1)$, conditioning a stable process to continuously limit to a point (which, by stationary and independent increments, can always be arranged to be $\left.\theta \in \mathbb{S}^{d-1}\right)$ results in a family of probability measures $\left(\mathbb{P}_{x}^{\{\theta\}}, x \neq \theta\right)$ which can be identified via a Doob $h$-transform with $h_{\theta}(x)=|x-\theta|^{\alpha-d}$. Although the sense in which the conditioning is performed cannot be contextualised via (3), we see that the resulting $h$-transformation is consistent with the use of the harmonic function $H_{S}$.

The way in which we will prove Theorem 1 will be to prove the following subtle result which establishes the leading order behaviour of the probability of hitting the set $\mathrm{S}_{\varepsilon}$.
Theorem 2 Let $S \subseteq \mathbb{S}^{d-1}$ be a closed subset such that $\sigma_{1}(S)>0$.
(i) Suppose $\alpha \in(0,1)$. For $x \notin S$,

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0} \varepsilon^{\alpha-1} \mathbb{P}_{x}\left(\tau_{S_{\varepsilon}}<\infty\right)=2^{1-2 \alpha} \frac{\Gamma((d+\alpha-2) / 2)}{\pi^{d / 2} \Gamma(1-\alpha)} \frac{\Gamma((2-\alpha) / 2)}{\Gamma(2-\alpha)} H_{S}(x) \tag{5}
\end{equation*}
$$

(ii) When $\alpha=1$, we have that, for $x \notin S$,

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0}|\log \varepsilon| \mathbb{P}_{x}\left(\tau_{S_{\varepsilon}}<\infty\right)=\frac{\Gamma((d-1) / 2)}{\pi^{(d-1) / 2}} H_{S}(x) \tag{6}
\end{equation*}
$$

Theorem 2 also gives us the opportunity to understand the strike position of the conditioned stable process. Indeed, let $\mathrm{S}^{\prime}$ be a closed subset of S . Define $\mathrm{S}_{\varepsilon}^{\prime}=\{x \in$ $\mathbb{R}^{d}: 1-\varepsilon \leq|x| \leq 1+\varepsilon$ and $\left.\arg (x) \in \mathrm{S}^{\prime}\right\}$ and $\tau_{\mathrm{S}_{\varepsilon}^{\prime}}:=\inf \left\{t>0: X_{t} \in \mathrm{~S}_{\varepsilon}^{\prime}\right\}$. Then, $\left\{\tau_{\mathrm{S}_{\varepsilon}^{\prime}}<\infty\right\} \subseteq\left\{\tau_{\mathbf{S}_{\varepsilon}}<\infty\right\}$ and thanks to Theorem 2, when $\alpha \in(0,1)$, we have

$$
\lim _{\varepsilon \rightarrow 0} \mathbb{P}_{x}\left(\tau_{\mathbf{S}_{\varepsilon}^{\prime}}<\infty \mid \tau_{\mathbf{S}_{\varepsilon}}<\infty\right)=\lim _{\varepsilon \rightarrow 0} \frac{\varepsilon^{\alpha-1} \mathbb{P}_{x}\left(\tau_{\mathbf{S}_{\varepsilon}^{\prime}}<\infty\right)}{\varepsilon^{\alpha-1} \mathbb{P}_{x}\left(\tau_{\mathbf{S}_{\varepsilon}}<\infty\right)}=\frac{H_{\mathbf{S}^{\prime}}(x)}{H_{\mathrm{S}}(x)}, \quad x \notin \mathrm{~S}
$$

A similar statement also holds when $\alpha=1$ by changing the scaling in $\varepsilon$ to $|\log \varepsilon|$. This gives us the following result.
Corollary 1 For a closed $S \subseteq \mathbb{S}^{d-1}$ such that $\sigma_{1}(S)>0$ and $\alpha \in(0,1]$, we have that for all closed $S^{\prime} \subseteq S$,

$$
\begin{equation*}
\mathbb{P}_{x}^{S}\left(X_{\zeta-} \in S^{\prime}\right)=\frac{H_{S^{\prime}}(x)}{H_{S}(x)}, \quad x \notin S \tag{7}
\end{equation*}
$$

In light of the above Corollary, it is worth remarking that we can also see the probabilities $\mathbb{P}^{S}$ as the result of first conditioning to continuously hit $\mathbb{S}^{d-1}$ and then conditioning the strike point to be in S . Indeed, we note that, for $A \in \mathcal{F}_{t}$ and $t \geq 0$,

$$
\begin{aligned}
\mathbb{P}_{x}^{\mathbb{S}^{d-1}}\left(A \mid X_{\zeta-} \in \mathrm{S}\right) & =\mathbb{E}_{x}^{\mathbb{S}^{d-1}}\left[\mathbf{1}_{A} \frac{\mathbb{P}_{X_{t}^{d-1}}^{\mathbb{P}^{d-1}}\left(X_{\zeta-} \in \mathrm{S}\right)}{\mathbb{P}_{x}^{\mathbb{S}^{d-1}}\left(X_{\zeta-} \in \mathrm{S}\right)}\right] \\
& =\mathbb{E}_{x}\left[\mathbf{1}_{A} \frac{H_{\mathbb{S}^{d-1}}\left(X_{t}\right)}{H_{\mathbb{S}^{d-1}}(x)} \frac{H_{\mathrm{S}}\left(X_{t}\right)}{H_{\mathbb{S}^{d-1}}\left(X_{t}\right)} \frac{H_{\mathbb{S}^{d-1}}(x)}{H_{\mathrm{S}}(x)}\right] \\
& =\mathbb{E}_{x}\left[\mathbf{1}_{A} \frac{H_{\mathrm{S}}\left(X_{t}\right)}{H_{\mathrm{S}}(x)}\right] \\
& =\mathbb{P}_{x}^{\mathrm{S}}(A) .
\end{aligned}
$$

Moreover, by shrinking $\mathrm{S}^{\prime} \subseteq \mathrm{S} \subseteq \mathbb{S}^{d-1}$ to a singleton $\theta \in \mathbb{S}^{d-1}$, one can similarly show that

$$
\mathbb{P}_{x}^{S}\left(A \mid X_{\zeta-}=\theta\right)=\mathbb{P}_{x}^{\{\theta\}}(A)
$$

This has the flavour of a Williams' type decomposition that was shown for general Lévy processes conditioned to stay positive and subordinators conditioned to remain in an interval; see e.g [11] and [19].

## 3 Oscillatory Repulsion from S and Duality

Roughly speaking, we want to describe what we see when we time reverse the process $\left(X, \mathbb{P}^{\mathrm{S}}\right)$ from its strike point on S , i.e. its so-called dual process. Such a process will necessarily avoid visiting S. Recalling that, for $\alpha \in(0,1]$, the stable process hits spherical surfaces with probability zero (cf. [16, 25]), a heuristic guess for the aforesaid dual process is the stable process itself (see Fig. 1). This turns out to be precisely the case. In order to make this rigorous, we will use the language of Hunt-Nagasawa duality for Markov processes.

Suppose that $Y=\left(Y_{t}, t \leq \zeta\right)$ with probabilities $\mathrm{P}_{x}, x \in E$, is a regular Markov process on an open domain $E \subseteq \mathbb{R}^{d}$ (or more generally, a locally compact Hausdorff space with countable base), with cemetery state $\Delta$ and killing time $\zeta=\inf \{t>0$ : $\left.Y_{t}=\Delta\right\}$. Let us additionally write $\mathrm{P}_{\nu}=\int_{E} \nu(\mathrm{~d} a) \mathrm{P}_{a}$, for any probability measure $v$ on the state space of $Y$.

Suppose that $\mathcal{G}$ is the $\sigma$-algebra generated by $Y$ and write $\mathcal{G}\left(\mathrm{P}_{\nu}\right)$ for its completion by the null sets of $\mathrm{P}_{\nu}$. Moreover, write $\overline{\mathcal{G}}=\bigcap_{\nu} \mathcal{G}\left(\mathrm{P}_{\nu}\right)$, where the intersection is taken over all probability measures on the state space of $Y$, excluding the cemetery state. A finite random time k is called an $L$-time (generalized last exit time) if, given a coordinate process $\omega=\left(\omega_{t}, t \geq 0\right)$ on $\mathbb{D}(E)$,
(i) k is measurable in $\overline{\mathcal{G}}$, and $\mathrm{k} \leq \zeta$ almost surely with respect to $\mathrm{P}_{\nu}$, for all $\nu$,
(ii) $\{s<\mathrm{k}(\omega)-t\}=\left\{s<\theta_{t} \circ \mathrm{k}\right\}$ for all $t, s \geq 0$,


Fig. 1 The process $\left(X, \mathbb{P}^{S}\right)$ when time reversed is stochastically equal in law to $(X, \mathbb{P})$
where $\theta_{t}$ is the Markov shift of $\omega$ to time $t$. The most important examples of $L$-times are killing times and last exit times from closed sets.
Theorem 3 Suppose that $\alpha \in(0,1]$. For a given closed set $S \subset \mathbb{S}^{d-1}$ with $\sigma_{1}(S)>$ 0 , write

$$
\begin{equation*}
\nu(\mathrm{d} a):=\frac{\sigma_{1}(\mathrm{~d} a)}{\sigma_{1}(S)}, \quad a \in S . \tag{8}
\end{equation*}
$$

For every L-time k of $(X, \mathbb{P})$, the process $\left(X_{\left.(\mathrm{k}-t)_{-}, t<k\right)}\right.$ under $\mathbb{P}_{v}$ is a timehomogeneous Markov process whose transition probabilities agree with those of $\left(X, \mathbb{P}^{S}\right)$.

## 4 The Setting of a Subset in an $\mathbb{R}^{\boldsymbol{d - 1}}$ Hyperplane

As alluded to in the introduction, the methods used in Sects. 2 and 3 are robust enough to deal with the setting of an arbitrary $(d-1)$-dimensional hyperplane in $\mathbb{R}^{d}$. Without loss of generality, we can describe such a hyperplane with unit orthogonal vector $v \in \mathbb{S}^{d-1}$ via

$$
\mathbb{H}^{d-1}=\left\{x \in \mathbb{R}^{d}:(x, v)=0\right\},
$$

where $(\cdot, \cdot)$ is the usual Euclidean inner product. Henceforth, we will assume that $v \in \mathbb{S}^{d-1}$ is given, as it otherwise plays no role in the forthcoming. We are interested in defining the law of the stable process conditioned to hit $\mathrm{D} \subseteq \mathbb{H}^{d-1}$ in a similar spirit to the discussion in Sect. 2.

To this end, let us define

$$
\kappa_{\beta}=\inf \left\{t>0:-\beta<\left(v, X_{t}\right)<\beta\right\}, \quad \text { for } \beta>0
$$

Whenever it is well defined, we will write, for $t \geq 0, \Lambda \in \mathcal{F}_{t}$ and $x \notin \mathrm{D}$,

$$
\begin{equation*}
\mathbb{P}_{x}^{\mathrm{D}}(\Lambda, t<\zeta)=\lim _{\beta \rightarrow 0} \lim _{\varepsilon \rightarrow 0} \mathbb{P}_{x}\left(\Lambda, t<\kappa_{\beta} \mid \tau_{\mathrm{D}_{\varepsilon}}<\infty\right) \tag{9}
\end{equation*}
$$

where

$$
\tau_{\mathrm{D}_{\varepsilon}}=\inf \left\{t>0: X_{t} \in \mathrm{D}_{\varepsilon}\right\} \quad \text { and } \quad \mathrm{D}_{\varepsilon}:=\left\{x \in \mathbb{R}^{d}:-\varepsilon \leq(v, x) \leq \varepsilon \text { and } \hat{x} \in \mathrm{D}\right\} .
$$

Here $\hat{x}$ denotes the orthogonal projection of $x$ onto $\mathbb{H}^{d-1}$; in other words. $\hat{x}=$ $x-v(v, x)$. We can gather the analogous conclusions of Theorems 1, 2, 3 and Corollary 1 into one theorem.

Theorem 4 Suppose that $\alpha \in(0,1]$ and the closed and bounded set $D \subseteq \mathbb{H}^{d-1}$ is such that $0<\ell_{d-1}(D)<\infty$, where we recall that $\ell_{d-1}$ is $(d-1)$-dimensional Lebesgue measure.
(i) Suppose $\alpha \in(0,1)$. For $x \notin D$,

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0} \varepsilon^{\alpha-1} \mathbb{P}_{x}\left(\tau_{D_{\varepsilon}}<\infty\right)=2^{1-\alpha} \pi^{-(d-2) / 2} \frac{\Gamma\left(\frac{d-2}{2}\right) \Gamma\left(\frac{d-\alpha}{2}\right) \Gamma\left(\frac{2-\alpha}{2}\right)^{2}}{\Gamma\left(\frac{1-\alpha}{2}\right) \Gamma\left(\frac{d-1}{2}\right) \Gamma(2-\alpha)} M_{D}(x) \tag{10}
\end{equation*}
$$

where

$$
M_{D}(x)=\int_{D}|x-y|^{\alpha-d} \ell_{d-1}(\mathrm{~d} y), \quad x \notin D
$$

(ii) Suppose $\alpha=1$. For $x \notin D$,

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0}|\log \varepsilon| \mathbb{P}_{x}\left(\tau_{D_{\varepsilon}}<\infty\right)=\frac{\Gamma\left(\frac{d-2}{2}\right)}{\pi^{(d-2) / 2}} M_{D}(x) \tag{11}
\end{equation*}
$$

(iii) The limit (9) makes sense, therefore the process $\left(X, \mathbb{P}^{D}\right)$ is well defined and

$$
\begin{equation*}
\left.\frac{\mathrm{d} \mathbb{P}_{x}^{D}}{\mathrm{~d} \mathbb{P}_{x}}\right|_{\mathcal{F}_{t}}=\frac{M_{D}\left(X_{t}\right)}{M_{D}(x)}, \quad t \geq 0, x \notin D \tag{12}
\end{equation*}
$$

(iv) We have for all closed $D^{\prime} \subseteq D$,

$$
\begin{equation*}
\mathbb{P}_{x}^{D}\left(X_{\zeta-} \in D^{\prime}\right)=\frac{M_{D^{\prime}}(x)}{M_{D}(x)}, \quad x \notin D \tag{13}
\end{equation*}
$$

(v) Write $\nu(\mathrm{d} a):=\ell_{d-1}(\mathrm{~d} a) / \ell_{d-1}(D), a \in D$. For every L-time k of $(X, \mathbb{P})$, the process $\left(X_{(\mathrm{k}-t)-}, t<\mathrm{k}\right)$ under $\mathbb{P}_{v}$ is a time-homogeneous Markov process whose transition probabilities agree with those of $\left(X, \mathbb{P}^{D}\right)$.

Roughly speaking, Theorem 4 are to be expected as, following the ideas of [22] one may map $\mathbb{S}^{d-1}$ onto $\mathbb{H}^{d-1}$ via a standard sphere inversion transformation, which, thanks to the Riesz-Bogdan-Żak transform, also transforms the paths of the stable processes into that of a $h$-transformed stable processes; see [8]. The proofs we have given below, however, are direct nonetheless, following similar steps to those of Theorems 1, 2 and 3, as well as Corollary 1.

## 5 Heuristic for the Proof of Theorem 2

Let us begin with a sketch of the proof of Theorem 2. We start by recalling an identity that is known in quite a general setting from the potential analysis literature; see for example Section 13.11 of [13] and Section VI. 2 of [7]. Suppose that $A$ is a bounded closed set and let $\tau_{A}=\inf \left\{t>0: X_{t} \in A\right\}$. Let $\mu_{A}$ be a finite measure supported on $A$, which is absolutely continuous with respect to Lebesgue measure and define its potential by

$$
U \mu_{A}(x):=\int_{A}|x-y|^{\alpha-d} \mu_{A}(\mathrm{~d} y), \quad x \in \mathbb{R}^{d}
$$

On account of the fact that $\mu_{A}$ is absolutely continuous, recalling that $|x|^{\alpha-d}$ is the potential of the stable process issued from the origin, stationary and independent increments allows us to identify

$$
U \mu_{A}(x)=\int_{A}|x-y|^{\alpha-d} m_{A}(y) \ell_{d}(\mathrm{~d} y)=\mathbb{E}_{x}\left[\int_{0}^{\infty} m_{A}\left(X_{t}\right) \mathrm{d} t\right], \quad x \notin A
$$

where $m_{A}$ is the density of $\mu_{A}$ with respect to Lebesgue measure, $\ell_{d}$. As the support of $\mu_{A}$ is precisely $A$, we must have $m_{A}(y)=0$ for all $y \notin A$. As such, the Strong Markov Property tells us that
$U \mu_{A}(x)=\mathbb{E}_{x}\left[\mathbf{1}_{\left\{\tau_{A}<\infty\right\}} \int_{\tau_{A}}^{\infty} m_{A}\left(X_{t}\right) \mathrm{d} t\right]=\mathbb{E}_{x}\left[U \mu_{A}\left(X_{\tau_{A}}\right) \mathbf{1}_{\left\{\tau_{A}<\infty\right\}}\right], \quad x \notin A$.

Note, the above equality is also true when $x \in A$ as, in that case, $\tau_{A}=0$.
Replacing $\tau_{A}$ by a general stopping time $\tau$ in the above calculation changes the first equality in (14) to an inequality, thus giving the excessive property

$$
\begin{equation*}
U \mu_{A}(x) \geq \mathbb{E}_{x}\left[U \mu_{A}\left(X_{\tau}\right) \mathbf{1}_{\{\tau<\infty\}}\right], \quad x \in \mathbb{R}^{d} \tag{15}
\end{equation*}
$$

This family of inequalities together with the Strong Markov Property easily gives us the classical result that $\left(U \mu_{A}\left(X_{t}\right), t \geq 0\right)$ is a supermartingale.

Let us now suppose that $\mu$ can be constructed in such a way that it is supported on $A$ such that, for all $x \in A, U \mu(x)=1$. We then recover from identity (14) the corollary to Theorem 1 in Chapter 5 of [13], see also equation (21) in the same chapter, which states that

$$
\mathbb{P}_{x}\left(\tau_{A}<\infty\right)=U \mu(x), \quad x \notin A
$$

Returning to the problem at hand, we can use the principals above to develop a 'guess and verify' approach to the proof, in particular, since we are not chasing an exact formula for $\mathbb{P}_{x}\left(\tau_{\mathbf{S}_{\varepsilon}}<\infty\right)$, but rather the asymptotic leading order behaviour.

Indeed, suppose we can 'guess' a measure, say $\mu_{\varepsilon}^{\mathrm{S}}$, supported on $\mathrm{S}_{\varepsilon}$, such that

$$
\begin{equation*}
U \mu_{\varepsilon}^{\mathrm{S}}(x)=1+o(1), \quad x \in \mathrm{~S}_{\varepsilon} \text { as } \varepsilon \rightarrow 0 \tag{16}
\end{equation*}
$$

so that

$$
\begin{equation*}
(1+o(1)) \mathbb{P}_{x}\left(\tau_{\mathrm{S}_{\varepsilon}}<\infty\right)=U \mu_{\varepsilon}^{\mathrm{S}}(x), \quad x \notin \mathrm{~S}_{\varepsilon} \tag{17}
\end{equation*}
$$

Then, this would be a good basis from which to draw out the leading order decay in $\varepsilon$, especially if our guess of $\mu_{\varepsilon}$ is such that $U \mu_{\varepsilon}$ is tractable.

In one dimension, we know from Lemma 1 of [26], that for a one-dimensional symmetric stable process, the unique measure that satisfies (16) has density ( $1-$ $y)^{-\alpha / 2}(1+y)^{-\alpha / 2}$, i.e.

$$
\begin{equation*}
\int_{-1}^{1}|x-y|^{\alpha-1}(1-y)^{-\alpha / 2}(1+y)^{-\alpha / 2} \mathrm{~d} y=1, \quad x \in[-1,1] . \tag{18}
\end{equation*}
$$

We can use this to build a reasonable choice of $\mu_{\varepsilon}^{S}$. Indeed, writing $X=$ $|X| \arg (X)$, when $X$ begins in the neighbourhood of S , then $|X|$ begins in the neighbourhood of 1 and $\arg (X)$, essentially, from within S . On short-time scales and short-range, the time change $|X|$ behaves similarly to a one-dimensional stable process. Moreover, $\arg (X)$ is an isotropic process. A reasonable guess for $\mu_{\varepsilon}^{S}$ would be to base it on the measure

$$
\begin{equation*}
\mu_{\varepsilon}(\mathrm{d} y)=c_{\alpha, d}(|y|-(1-\varepsilon))^{-\alpha / 2}(1+\varepsilon-|y|)^{-\alpha / 2} \ell_{d}(\mathrm{~d} y), \tag{19}
\end{equation*}
$$

restricted to $\mathrm{S}_{\varepsilon}$, where we recall $c_{\alpha, d}$ is a constant to be determined so that (16) holds. As we will shortly see, when $\alpha \in(0,1)$, the constant $c_{\alpha, d}$ does not depend on $\varepsilon$, however, when $\alpha=1$, in order to respect (16) we need to make it depend on $\varepsilon$.

## 6 Proof of Theorem 2 (i)

As alluded to in the previous section, we will work with the guess $\mu_{\varepsilon}^{\mathrm{S}}$ given by the measure $\mu_{\varepsilon}$ defined in (19) restricted to $\mathrm{S}_{\varepsilon}$. In order to show (16), we will take advantage of some of the symmetric features of $\mu_{\varepsilon}$, when seen as a measure over $\mathbb{S}_{\varepsilon}^{d-1}=\left\{x \in \mathbb{R}^{d}: 1-\varepsilon \leq|x| \leq 1+\varepsilon\right\}$. For a subset $A \subset \mathbb{S}_{\varepsilon}^{d-1}$ we define $\mu_{\varepsilon}^{A}$ the restriction of $\mu_{\varepsilon}$ to $A$. In particular, writing $\mu_{\varepsilon}^{(1)}$ as $\mu_{\varepsilon}$ restricted to $\mathbb{S}_{\varepsilon}^{d-1}$ and $\mu_{\varepsilon}^{(2)}$ as $\mu_{\varepsilon}$ restricted to $\hat{\mathrm{S}}_{\varepsilon}:=\mathbb{S}_{\varepsilon}^{d-1} \backslash \mathbf{S}_{\varepsilon}$, we have the obvious difference

$$
\begin{equation*}
U \mu_{\varepsilon}^{\mathrm{S}}(x)=U \mu_{\varepsilon}^{(1)}(x)-U \mu_{\varepsilon}^{(2)}(x), \quad x \in \mathrm{~S}_{\varepsilon} \tag{20}
\end{equation*}
$$

Moreover, we would like to introduce

$$
\mu_{\varepsilon, \delta}^{(2)}:=\left.\mu_{\varepsilon}\right|_{\hat{\mathbf{S}}_{\varepsilon}^{\delta}}
$$

where $\hat{S}_{\varepsilon}^{\delta}=\mathbb{S}_{\varepsilon}^{d-1} \backslash \mathrm{~S}_{\varepsilon}^{\delta}$ and
$\mathbf{S}_{\varepsilon}^{\delta}:=\left\{x \in \mathbb{R}^{d}: 1-\varepsilon<|x|<1+\varepsilon\right.$ and $\left.\arg (x) \in \mathbf{S}^{\delta}\right\}$, where $\mathrm{S}^{\delta}=\left\{x \in \mathbb{S}^{d-1}: \operatorname{dist}(\arg (x), \mathrm{S})<\delta\right\}$,
for some small $\delta>0$, which, in due course, will depend on $\varepsilon$. Note that, since $S$ is closed, $\mathbf{S}^{\delta}$ (resp. $\mathbf{S}_{\varepsilon}^{\delta}$ ) shrinks to S (resp. $\mathrm{S}_{\varepsilon}$ ) when $\delta \rightarrow 0$. Then, we also have that

$$
\begin{equation*}
U \mu_{\varepsilon}^{\mathrm{S}^{\delta}}(x)=U \mu_{\varepsilon}^{(1)}(x)-U \mu_{\varepsilon, \delta}^{(2)}(x), \quad x \in \mathrm{~S}_{\varepsilon} . \tag{21}
\end{equation*}
$$

The estimate (21) will be useful for a certain lower bound that will give us what we need to prove Theorem 2. We need to prove two technical lemmas first. The first one deals with the term $U \mu_{\varepsilon}^{(1)}$.

Lemma 1 Suppose that we choose

$$
c_{\alpha, d}=\frac{\Gamma((d+\alpha-2) / 2)}{2^{\alpha} \pi^{d / 2} \Gamma(1-\alpha) \Gamma((2-\alpha) / 2)} .
$$

Then,

$$
\lim _{\varepsilon \rightarrow 0} \sup _{x \in \mathbb{S}_{\varepsilon}^{d-1}}\left|U \mu_{\varepsilon}^{(1)}(x)-1\right|=0
$$

Proof Appealing to (77), we have, for $x \in \mathbb{S}_{\varepsilon}^{d-1}$,

$$
\begin{align*}
& U \mu_{\varepsilon}^{(1)}(x) \\
& =c_{\alpha, d} \int_{\mathbb{S}_{\varepsilon}^{d-1}}|x-y|^{\alpha-d}(|y|-(1-\varepsilon))^{-\alpha / 2}(1+\varepsilon-|y|)^{-\alpha / 2} \ell_{d}(\mathrm{~d} y) \\
& =\frac{2 c_{\alpha, d} \pi^{(d-1) / 2}}{\Gamma((d-1) / 2)} \int_{1-\varepsilon}^{1+\varepsilon} \frac{r^{d-1}}{(r-(1-\varepsilon))^{\alpha / 2}(1+\varepsilon-r)^{\alpha / 2}} \mathrm{~d} r \int_{0}^{\pi} \frac{\sin ^{d-2} \theta \mathrm{~d} \theta}{\left(|x|^{2}-2|x| r \cos \theta+r^{2}\right)^{(d-\alpha) / 2}} \\
& =\frac{2 c_{\alpha, d} \pi^{d / 2}}{\Gamma(d / 2)}|x|^{\alpha-d} \int_{1-\varepsilon}^{|x|} \frac{2 F_{1}\left(\frac{d-\alpha}{2}, 1-\frac{\alpha}{2} ; \frac{d}{2} ;(r /|x|)^{2}\right) r^{d-1}}{(r-(1-\varepsilon))^{\alpha / 2}(1+\varepsilon-r)^{\alpha / 2}} \mathrm{~d} r \\
& \quad+\frac{2 c_{\alpha, d} \pi^{d / 2}}{\Gamma(d / 2)} \int_{|x|}^{1+\varepsilon} \frac{{ }_{2} F_{1}\left(\frac{d-\alpha}{2}, 1-\frac{\alpha}{2} ; \frac{d}{2} ;(|x| / r)^{2}\right) r^{\alpha-1}}{(r-(1-\varepsilon))^{\alpha / 2}(1+\varepsilon-r)^{\alpha / 2}} \mathrm{~d} r . \tag{22}
\end{align*}
$$

With a simple change of variables we can reduce this more simply to

$$
\begin{align*}
& U \mu_{\varepsilon}^{(1)}(x)=\frac{2 c_{\alpha, d} \pi^{d / 2}}{\Gamma(d / 2)} \int_{\frac{1-\varepsilon}{|x|}}^{1} \frac{{ }_{2} F_{1}\left(\frac{d-\alpha}{2}, 1-\frac{\alpha}{2} ; \frac{d}{2} ; r^{2}\right) r^{d-1}}{\left(r-\frac{1-\varepsilon}{|x|}\right)^{\alpha / 2}\left(\frac{1+\varepsilon}{|x|}-r\right)^{\alpha / 2}} \mathrm{~d} r \\
&+\frac{2 c_{\alpha, d} \pi^{d / 2}}{\Gamma(d / 2)} \int_{1}^{\frac{1+\varepsilon}{|x|}} \frac{2 F_{1}\left(\frac{d-\alpha}{2}, 1-\frac{\alpha}{2} ; \frac{d}{2} ; r^{-2}\right) r^{\alpha-1}}{\left(r-\frac{1-\varepsilon}{|x|}\right)^{\alpha / 2}\left(\frac{1+\varepsilon}{|x|}-r\right)^{\alpha / 2}} \mathrm{~d} r . \tag{23}
\end{align*}
$$

For the first term on the right-hand side of (23), we can appeal to (71) and (72) to deduce that

$$
\begin{align*}
& \lim _{\varepsilon \rightarrow 0} \sup _{x \in S_{\varepsilon}^{d-1}} \left\lvert\, \frac{2 c_{\alpha, d} \pi^{d / 2}}{\Gamma(d / 2)} \int_{\frac{1-\varepsilon}{|x|}}^{1} \frac{2 F_{1}\left(\frac{d-\alpha}{2}, 1-\frac{\alpha}{2} ; \frac{d}{2} ; r^{2}\right) r^{d-1}}{\left(r-\frac{1-\varepsilon}{|x|}\right)^{\alpha / 2}\left(\frac{1+\varepsilon}{|x|}-r\right)^{\alpha / 2}} \mathrm{~d} r\right. \\
& \quad-\frac{2 c_{\alpha, d} \pi^{d / 2} \Gamma(1-\alpha)}{\Gamma((d-\alpha) / 2) \Gamma((2-\alpha) / 2)} \int_{\frac{1-\varepsilon}{|x|}}^{1} \frac{\left(1-r^{2}\right)^{\alpha-1} r^{d-1}}{\left(r-\frac{1-\varepsilon}{|x|}\right)^{\alpha / 2}\left(\frac{1+\varepsilon}{|x|}-r\right)^{\alpha / 2}} \mathrm{~d} r \\
&  \tag{24}\\
& \left.\quad \quad-\frac{2 c_{\alpha, d} \pi^{d / 2} \Gamma(1-\alpha)}{\Gamma(\alpha / 2) \Gamma((d+\alpha-2) / 2)} \int_{\frac{1-\varepsilon}{|x|}}^{1} \frac{r^{d-1}}{\left(r-\frac{1-\varepsilon}{|x|}\right)^{\alpha / 2}\left(\frac{1+\varepsilon}{|x|}-r\right)^{\alpha / 2}} \mathrm{~d} r \right\rvert\,=0 .
\end{align*}
$$

Note that, by using the transformation $r=(1-\varepsilon+2 \varepsilon u) /|x|$,

$$
\begin{align*}
\int_{\frac{1-\varepsilon}{|x|}}^{1} r^{d-1}(r- & \left.\frac{1-\varepsilon}{|x|}\right)^{-\alpha / 2}\left(\frac{1+\varepsilon}{|x|}-r\right)^{-\alpha / 2} \mathrm{~d} r \\
& =|x|^{\alpha-d}(2 \varepsilon)^{1-\alpha} \int_{0}^{(|x|-1+\varepsilon) / 2 \varepsilon}(2 \varepsilon u+1-\varepsilon)^{d-1} u^{-\alpha / 2}(1-u)^{-\alpha / 2} \mathrm{~d} u \\
& \leq|x|^{\alpha-d}(2 \varepsilon)^{1-\alpha} \frac{\Gamma((2-\alpha) / 2)^{2}}{\Gamma(2-\alpha)} \tag{25}
\end{align*}
$$

which tends to zero uniformly in $x \in \mathbb{S}_{\varepsilon}^{d-1}$ as $\varepsilon \rightarrow 0$.
The asymptotic (25) also tells us that the approximating term of interest in (24) is the middle term. For that, we can use (78) to observe

$$
\begin{align*}
\lim _{\varepsilon \rightarrow 0} \sup _{x \in S_{\varepsilon}^{d-1}} \left\lvert\, \int_{\frac{1-\varepsilon}{|x|}}^{1}\left(1-r^{2}\right)^{\alpha-1} r^{d-1}(r\right. & \left.-\frac{1-\varepsilon}{|x|}\right)^{-\alpha / 2}\left(\frac{1+\varepsilon}{|x|}-r\right)^{-\alpha / 2} \mathrm{~d} r \\
& \left.-2^{\alpha-1} \int_{\frac{1-\varepsilon}{|x|}}^{1}(1-r)^{\alpha-1}\left(r-\frac{1-\varepsilon}{|x|}\right)^{-\alpha / 2}\left(\frac{1+\varepsilon}{|x|}-r\right)^{-\alpha / 2} \mathrm{~d} r \right\rvert\,=0 \tag{26}
\end{align*}
$$

and

$$
\begin{align*}
& \frac{2^{\alpha} c_{\alpha, d} \pi^{d / 2} \Gamma(1-\alpha)}{\Gamma(\alpha / 2) \Gamma((d+\alpha-2) / 2)} \int_{\frac{1-\varepsilon}{|x|}}^{1}(1-r)^{\alpha-1}\left(r-\frac{1-\varepsilon}{|x|}\right)^{-\alpha / 2}\left(\frac{1+\varepsilon}{|x|}-r\right)^{-\alpha / 2} \mathrm{~d} r \\
& =\frac{2^{\alpha} c_{\alpha, d} \pi^{d / 2} \Gamma(1-\alpha)}{\Gamma(\alpha / 2) \Gamma((d+\alpha-2) / 2)} \int_{0}^{1-\frac{1-\varepsilon}{|x|}} u^{\alpha-1}\left(1-\frac{1-\varepsilon}{|x|}-u\right)^{-\alpha / 2}\left(\frac{1+\varepsilon}{|x|}-1+u\right)^{-\alpha / 2} \mathrm{~d} r \\
& =\frac{2^{\alpha} c_{\alpha, d} \pi^{d / 2} \Gamma(1-\alpha) \Gamma((2-\alpha) / 2) \Gamma(\alpha)}{\Gamma(\alpha / 2) \Gamma((d+\alpha-2) / 2) \Gamma((2+\alpha) / 2)}\left(\frac{|x|-1+\varepsilon}{1+\varepsilon-|x|}\right)^{\alpha / 2} \\
& \quad{ }_{2} F_{1}\left(\alpha / 2, \alpha ; 1+\alpha / 2 ;-\frac{|x|-1+\varepsilon}{1+\varepsilon-|x|}\right) . \tag{27}
\end{align*}
$$

The second term on the right-hand side of (23) can be dealt with similarly. Indeed, using (72) we can produce an analogous statement to (24), from which, the leading order approximating term is the integral

$$
\begin{align*}
& \frac{2 c_{\alpha, d} \pi^{d / 2} \Gamma(1-\alpha)}{\Gamma(\alpha / 2) \Gamma((d+\alpha-2) / 2)} \int_{1}^{\frac{1+\varepsilon}{|x|}}\left(1-r^{-2}\right)^{\alpha-1} r^{d-1}\left(r-\frac{1-\varepsilon}{|x|}\right)^{-\alpha / 2}\left(\frac{1+\varepsilon}{|x|}-r\right)^{-\alpha / 2} \mathrm{~d} r \\
& \sim \frac{2^{\alpha} c_{\alpha, d} \pi^{d / 2} \Gamma(1-\alpha)}{\Gamma(\alpha / 2) \Gamma((d+\alpha-2) / 2)} \int_{1}^{\frac{1+\varepsilon}{|x|}}(r-1)^{\alpha-1}\left(r-\frac{1-\varepsilon}{|x|}\right)^{-\alpha / 2}\left(\frac{1+\varepsilon}{|x|}-r\right)^{-\alpha / 2} \mathrm{~d} r \\
& =\frac{2^{\alpha} c_{\alpha, d} \pi^{d / 2} \Gamma(1-\alpha)}{\Gamma(\alpha / 2) \Gamma((d+\alpha-2) / 2)} \int_{0}^{\frac{1+\varepsilon}{|x|}-1} u^{\alpha-1}\left(u+1-\frac{1-\varepsilon}{|x|}\right)^{-\alpha / 2}\left(\frac{1+\varepsilon}{|x|}-1-u\right)^{-\alpha / 2} \mathrm{~d} u \\
& =\frac{2^{\alpha} c_{\alpha, d} \pi^{d / 2} \Gamma(1-\alpha) \Gamma((2-\alpha) / 2) \Gamma(\alpha)}{\Gamma(\alpha / 2) \Gamma((d+\alpha-2) / 2) \Gamma((2+\alpha) / 2)}\left(\frac{1+\varepsilon-|x|}{|x|-1+\varepsilon}\right)^{\alpha / 2} \\
& { }_{2} F_{1}\left(\alpha / 2, \alpha ; 1+\alpha / 2 ;-\frac{1+\varepsilon-|x|}{|x|-1+\varepsilon}\right), \tag{28}
\end{align*}
$$

uniformly for $x \in \mathbb{S}_{\varepsilon}^{d-1}$ as $\varepsilon \rightarrow 0$, where we have again used (78) to develop the right-hand side.

Somewhat remarkably, if we add together the right-hand side of (27) and (28), using the identity in (76), we see that the sum is equal to

$$
\begin{equation*}
\frac{2^{\alpha} c_{\alpha, d} \pi^{d / 2} \Gamma(1-\alpha) \Gamma((2-\alpha) / 2)}{\Gamma((d+\alpha-2) / 2)}=1 \tag{29}
\end{equation*}
$$

where the equality with unity follows from the choice of $c_{\alpha, d}$ in the statement of the lemma.

Piecing together then uniform estimates above as well as the simplification of the two integrals (27) and (28) as well as the decay of the term (25) in (24) and the
analogous term when dealing with the second term on the right-hand side of (23), the statement of the lemma follows.

Next we deal with the term $U \mu_{\varepsilon, \delta}^{(2)}$.
Lemma 2 Recalling that $c_{\alpha, d}$ is the constant given in Lemma 1, take $\delta(\varepsilon)=$ $\varepsilon^{(1-\alpha) / 2(d-\alpha)}$, then

$$
\limsup _{\varepsilon \rightarrow 0} \sup _{x \in S_{\varepsilon}} \varepsilon^{(\alpha-1) / 2} U \mu_{\varepsilon, \delta(\varepsilon)}^{(2)}(x) \leq C_{\alpha, d},
$$

where

$$
C_{\alpha, d}=c_{\alpha, d} \frac{2^{2-\alpha} \pi^{(d-1) / 2} \Gamma((2-\alpha) / 2)^{2}}{\Gamma(2-\alpha) \Gamma((d-1) / 2)}
$$

In particular,

$$
\lim _{\varepsilon \rightarrow 0} \sup _{x \in \mathcal{S}_{\varepsilon}} U \mu_{\varepsilon, \delta(\varepsilon)}^{(2)}(x)=0
$$

Proof Since $x \in \mathrm{~S}_{\varepsilon}$ and $y \in \hat{\mathrm{~S}}_{\varepsilon}^{\delta}$, i.e. $|x-y|>\delta$, we have,

$$
\begin{align*}
\sup _{x \in \mathrm{~S}_{\varepsilon}} U \mu_{\varepsilon, \delta}^{(2)}(x) & =\int_{\hat{\mathbf{S}}_{\varepsilon}^{\delta}} \frac{1}{|x-y|^{d-\alpha}} \mu_{\varepsilon}(\mathrm{d} y) \\
& \leq \frac{1}{\delta^{d-\alpha}} \int_{\hat{\mathbf{S}}_{\varepsilon}^{\delta}} \mu_{\varepsilon}(\mathrm{d} y) \\
& \leq \frac{1}{\delta^{d-\alpha}} \frac{2 \pi^{(d-1) / 2}}{\Gamma((d-1) / 2)} \int_{1-\varepsilon}^{1+\varepsilon} r^{d-1} m_{\varepsilon}(r) \mathrm{d} r \tag{30}
\end{align*}
$$

where $m_{\varepsilon}(r)=c_{\alpha, d}(r-(1-\varepsilon))^{-\alpha / 2}(1+\varepsilon-r)^{-\alpha / 2}$. It is easy to see that

$$
\begin{align*}
\int_{1-\varepsilon}^{1+\varepsilon} m_{\varepsilon}(r) \mathrm{d} r & =c_{\alpha, d} \int_{1-\varepsilon}^{1+\varepsilon}(r-(1-\varepsilon))^{-\alpha / 2}(1+\varepsilon-r)^{-\alpha / 2} \mathrm{~d} r \\
& =c_{\alpha, d} \varepsilon^{1-\alpha} 2^{1-\alpha} \frac{\Gamma((2-\alpha) / 2)^{2}}{\Gamma(2-\alpha)} . \tag{31}
\end{align*}
$$

Putting (30) and (31) together we have

$$
\begin{equation*}
\sup _{x \in \mathbf{S}_{\varepsilon}} U \mu_{\varepsilon, \delta}^{(2)}(x) \leq c_{\alpha, d} \frac{2^{2-\alpha} \pi^{(d-1) / 2} \Gamma((2-\alpha) / 2)^{2}}{\Gamma(2-\alpha) \Gamma((d-1) / 2)} \times \frac{\varepsilon^{1-\alpha}}{\delta^{d-\alpha}}(1+\varepsilon)^{d-1} . \tag{32}
\end{equation*}
$$

By choosing $\delta=\delta(\varepsilon)$, the result follows.

Let us now return to the proof of Theorem 2. We show that we can make careful sense of (16) and (17). Using (20) in (14) we see that for $x \notin \mathrm{~S}$,

$$
\begin{align*}
U \mu_{\varepsilon}^{\mathrm{S}}(x) & =\mathbb{E}_{x}\left[\left(U \mu_{\varepsilon}^{(1)}\left(X_{\tau_{\mathrm{S}_{\varepsilon}}}\right)-1\right) ; \tau_{\mathrm{S}_{\varepsilon}}<\infty\right]+\mathbb{P}_{x}\left(\tau_{\mathrm{S}_{\varepsilon}}<\infty\right)-\mathbb{E}_{x}\left[U \mu_{\varepsilon}^{(2)}\left(X_{\tau_{\mathrm{S}_{\varepsilon}}}\right) ; \tau_{\mathbf{S}_{\varepsilon}}<\infty\right] \\
& \leq \mathbb{E}_{x}\left[\left(U \mu_{\varepsilon}^{(1)}\left(X_{\tau_{\mathbf{S}_{\varepsilon}}}\right)-1\right) ; \tau_{\mathrm{S}_{\varepsilon}}<\infty\right]+\mathbb{P}_{x}\left(\tau_{\mathrm{S}_{\varepsilon}}<\infty\right) \tag{33}
\end{align*}
$$

Then, due to Lemma 1 , for each $x \notin \mathrm{~S}$ and $v>0$, we can choose $\varepsilon$ sufficiently small such that

$$
\begin{equation*}
U \mu_{\varepsilon}^{\mathrm{S}}(x) \leq(1+v) \mathbb{P}_{x}\left(\tau_{\mathrm{S}_{\varepsilon}}<\infty\right) \tag{34}
\end{equation*}
$$

Since we can take $v$ arbitrarily small, we have the lower bound on a liminf version of the statement of Theorem 2 given by

$$
\begin{equation*}
\liminf _{\varepsilon \rightarrow 0} \varepsilon^{\alpha-1} U \mu_{\varepsilon}^{\mathrm{S}}(x) \leq \liminf _{\varepsilon \rightarrow 0} \varepsilon^{\alpha-1} \mathbb{P}_{x}\left(\tau_{\mathbf{S}_{\varepsilon}}<\infty\right), \quad x \notin \mathbf{S} \tag{35}
\end{equation*}
$$

On the other hand, suppose instead of S , we replace its role by $\mathrm{S}^{\delta(\varepsilon)}$, where $\delta(\varepsilon)$ was given in the statement of Lemma 2, we have from the excessive property (15) associated to $U \mu_{\varepsilon}^{\mathrm{S}^{\delta(\varepsilon)}}$ that

$$
\begin{equation*}
U \mu_{\varepsilon}^{\mathrm{S}^{\delta(\varepsilon)}}(x) \geq \mathbb{E}_{x}\left[U \mu_{\varepsilon}^{\mathrm{S}^{\delta(\varepsilon)}}\left(X_{\tau_{\mathrm{S}_{\varepsilon}}}\right) ; \tau_{\mathrm{S}_{\varepsilon}}<\infty\right], \quad x \notin \mathrm{~S} . \tag{36}
\end{equation*}
$$

Now appealing to (21), we get
$U \mu_{\varepsilon}^{\delta^{\delta(\varepsilon)}}(x) \geq \mathbb{E}_{x}\left[U \mu_{\varepsilon}^{(1)}\left(X_{\tau_{S_{\varepsilon}}}\right)-1 ; \tau_{\mathrm{S}_{\varepsilon}}<\infty\right]+\mathbb{P}_{x}\left(\tau_{\mathrm{S}_{\varepsilon}}<\infty\right)-\mathbb{E}_{x}\left[U \mu_{\varepsilon, \delta(\varepsilon)}^{(2)}\left(X_{\boldsymbol{S}_{\varepsilon}}\right) ; \tau_{\mathrm{S}_{\varepsilon}}<\infty\right]$.
Appealing to Lemmas 1 and 2, for each $v>0$, we can choose $\varepsilon$ small enough such that, for each $x \notin \mathrm{~S}$,

$$
\begin{equation*}
U \mu_{\varepsilon}^{\mathrm{S}^{\delta(\varepsilon)}}(x) \geq(1-v) \mathbb{P}_{x}\left(\tau_{\mathbf{S}_{\varepsilon}}<\infty\right) \tag{37}
\end{equation*}
$$

Hence, since we can choose $v$ as small as we like, we have

$$
\begin{equation*}
\limsup _{\varepsilon \rightarrow 0} \varepsilon^{\alpha-1} U \mu_{\varepsilon}^{\mathrm{S}^{\delta(\varepsilon)}}(x) \geq \limsup _{\varepsilon \rightarrow 0} \varepsilon^{\alpha-1} \mathbb{P}_{x}\left(\tau_{\mathbf{S}_{\varepsilon}}<\infty\right), \quad x \notin \mathrm{~S} . \tag{38}
\end{equation*}
$$

It follows from (35) and (38) that, as soon as

$$
\begin{equation*}
\limsup _{\varepsilon \rightarrow 0} \varepsilon^{\alpha-1} U \mu_{\varepsilon}^{\delta^{\delta(\varepsilon)}}(x)=\liminf _{\varepsilon \rightarrow 0} \varepsilon^{\alpha-1} U \mu_{\varepsilon}^{\mathrm{S}}(x), \quad x \notin \mathrm{~S}, \tag{39}
\end{equation*}
$$

and noting that $U \mu_{\varepsilon}^{\mathrm{S}} \leq U \mu_{\varepsilon}^{\mathrm{S}^{\delta(\varepsilon)}}$, we have

$$
\lim _{\varepsilon \rightarrow 0} \varepsilon^{\alpha-1} U \mu_{\varepsilon}^{\mathrm{S}}(x)=\lim _{\varepsilon \rightarrow 0} \varepsilon^{\alpha-1} \mathbb{P}_{x}\left(\tau_{\mathbf{S}_{\varepsilon}}<\infty\right), \quad x \notin \mathrm{~S}
$$

Let us thus complete the proof by verifying the limit on the equality (39) holds and by finding the left-hand side limit in the previous equation.

To this end, using that $|x-y|^{\alpha-d}$ is continuous on $\mathrm{S}_{\varepsilon}$ and, when $x \notin \mathrm{~S}$, without loss of generality, we can take $\varepsilon$ small enough so that $x \notin \mathrm{~S}_{\varepsilon}$. For each $x \notin \mathrm{~S}$, using the Mean Valued Theorem, there exists a $r_{\varepsilon}^{*} \in(1-\varepsilon, 1+\varepsilon)$ such that

$$
\begin{align*}
U \mu_{\varepsilon}^{\mathrm{S}}(x) & =\int_{\mathrm{S}_{\varepsilon}}|x-y|^{\alpha-d} m_{\varepsilon}(|y|) \ell_{d}(\mathrm{~d} y) \\
& =\left(r_{\varepsilon}^{*}\right)^{d-1} \int_{\mathrm{S}}\left|x-r_{\varepsilon}^{*} \theta\right|^{\alpha-d} \sigma_{1}(\mathrm{~d} \theta) \int_{1-\varepsilon}^{1+\varepsilon} m_{\varepsilon}(r) \mathrm{d} r \tag{40}
\end{align*}
$$

where we recall that $m_{\varepsilon}(r)=c_{\alpha, d}(r-(1-\varepsilon))^{-\alpha / 2}(1+\varepsilon-r)^{-\alpha / 2}$. By using (31) we get

$$
\begin{equation*}
\varepsilon^{\alpha-1} U \mu_{\varepsilon}^{\mathrm{S}}(x)=\left(r_{\varepsilon}^{*}\right)^{d-1} 2^{1-\alpha} c_{\alpha, d} \frac{\Gamma((2-\alpha) / 2)^{2}}{\Gamma(2-\alpha)} \int_{\mathrm{S}}\left|x-r_{\varepsilon}^{*} \theta\right|^{\alpha-d} \sigma_{1}(\mathrm{~d} \theta), \quad x \notin \mathrm{~S} . \tag{41}
\end{equation*}
$$

Taking limits in (41) as $\varepsilon \rightarrow 0$ and recalling the value of $c_{\alpha, d}$ from the statement of Lemma 1, we have, for $x \notin \mathrm{~S}$

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0} \varepsilon^{\alpha-1} U \mu_{\varepsilon}^{\mathrm{S}}(x)=2^{1-2 \alpha} \frac{\Gamma((d+\alpha-2) / 2)}{\pi^{d / 2} \Gamma(1-\alpha)} \frac{\Gamma((2-\alpha) / 2)}{\Gamma(2-\alpha)} \int_{\mathrm{S}}|x-\theta|^{\alpha-d} \sigma_{1}(\mathrm{~d} \theta) . \tag{42}
\end{equation*}
$$

An application of the recursion formula for gamma functions allows us to identify the right-hand side as equal to that of the right-hand side of (5). Very little changes in the above calculation if we replace $S$ by $\mathrm{S}^{\delta(\varepsilon)}$. As such, (42) allows us to conclude (39), and thus gives the statement of the Theorem 2.

## 7 Proof of Theorem 2 (ii)

The proof needs some adaptation when we deal with the case $\alpha=1$. Principally, we need to focus on Lemmas 1 and 2. What is different in these two lemmas is that the normalisation constant $c_{\alpha, d}$ must now depend on $\varepsilon$. The replacement for Lemma 1 and Lemma 2 (combined into one result) now takes the following form.

Lemma 3 Suppose that we define, for $0<\varepsilon<1$,

$$
\begin{equation*}
\mu_{\varepsilon}(\mathrm{d} y)=\frac{c_{1, d}}{|\log \varepsilon|}(|y|-(1-\varepsilon))^{-\alpha / 2}(1+\varepsilon-|y|)^{-\alpha / 2} \ell_{d}(\mathrm{~d} y), \tag{43}
\end{equation*}
$$

and

$$
c_{1, d}=\frac{\Gamma((d-1) / 2)}{\pi^{(d+1) / 2}} .
$$

(i) We have

$$
\lim _{\varepsilon \rightarrow 0} \sup _{x \in \mathbb{S}_{\varepsilon}^{d-1}}\left|U \mu_{\varepsilon}^{(1)}(x)-1\right|=0
$$

(ii) take $\delta(\varepsilon)=|\log \varepsilon|^{-1 / 2(d-1)}$, then

$$
\limsup _{\varepsilon \rightarrow 0} \sup _{x \in \mathcal{S}_{\varepsilon}} \sqrt{|\log \varepsilon|} U \mu_{\varepsilon, \delta(\varepsilon)}^{(2)}(x)<\infty,
$$

so that

$$
\lim _{\varepsilon \rightarrow 0} \sup _{x \in \mathcal{S}_{\varepsilon}} U \mu_{\varepsilon, \delta(\varepsilon)}^{(2)}(x)=0
$$

Proof We give only a sketch proof of both parts for the interested reader to use as a guide to reproduce the finer details.
(i) The essence of the proof is an adaptation of the proof of Lemma 1. We pick up the proof of the latter at the analogue of (23), albeit $\alpha=1$ and $c_{\alpha, d}$ is replaced by $c_{1, d} /|\log \varepsilon|$, i.e.

$$
\begin{align*}
U \mu_{\varepsilon}^{(1)}(x)=\frac{2 c_{1, d} \pi^{d / 2}}{|\log \varepsilon| \Gamma(d / 2)} \int_{\frac{1-\varepsilon}{|x|}}^{1} & \frac{2 F_{1}\left(\frac{d-1}{2}, \frac{1}{2} ; \frac{d}{2} ; r^{2}\right) r^{d-1}}{\left(r-\frac{1-\varepsilon}{|x|}\right)^{1 / 2}\left(\frac{1+\varepsilon}{|x|}-r\right)^{1 / 2}} \mathrm{~d} r \\
& \quad+\frac{2 c_{1, d} \pi^{d / 2}}{|\log \varepsilon| \Gamma(d / 2)} \int_{1}^{\frac{1+\varepsilon}{|x|}} \frac{2 F_{1}\left(\frac{d-1}{2}, \frac{1}{2} ; \frac{d}{2} ; r^{-2}\right) r^{1-1}}{\left(r-\frac{1-\varepsilon}{|x|}\right)^{1 / 2}\left(\frac{1+\varepsilon}{|x|}-r\right)^{1 / 2}} \mathrm{~d} r . \tag{44}
\end{align*}
$$

Appealing to (74), noting that $\log \left(1-r^{2}\right) \sim \log (1-r)+\log 2$, as $r \rightarrow 1$, we can deduce that there is an unimportant constant, say $\chi$, such that

$$
\begin{align*}
\lim _{\varepsilon \rightarrow 0} \sup _{x \in \mathbb{S}_{\varepsilon}^{d-1}} \mid & \frac{2 c_{1, d} \pi^{d / 2}}{|\log \varepsilon| \Gamma(d / 2)} \int_{\frac{1-\varepsilon}{|x|}}^{1} \frac{{ }_{2} F_{1}\left(\frac{d-1}{2}, \frac{1}{2} ; \frac{d}{2} ; r^{2}\right) r^{d-1}}{\left(r-\frac{1-\varepsilon}{|x|}\right)^{1 / 2}} \frac{\left(\frac{1+\varepsilon}{|x|}-r\right)^{1 / 2}}{} \mathrm{~d} r \\
& +\frac{2 c_{1, d} \pi^{d / 2}}{|\log \varepsilon| \Gamma((d-1) / 2) \Gamma(1 / 2)} \int_{\frac{1-\varepsilon}{|x|}}^{1} \frac{r^{d-1} \log (1-r)}{\left(r-\frac{1-\varepsilon}{|x|}\right)^{1 / 2}\left(\frac{1+\varepsilon}{|x|}-r\right)^{1 / 2}} \mathrm{~d} r \\
& \left.\quad-\frac{c_{1, d} \chi}{|\log \varepsilon|} \int_{\frac{1-\varepsilon}{|x|}}^{1} \frac{r^{d-1}}{\left(r-\frac{1-\varepsilon}{|x|}\right)^{1 / 2}\left(\frac{1+\varepsilon}{|x|}-r\right)^{1 / 2}} \mathrm{~d} r \right\rvert\,=0 . \tag{45}
\end{align*}
$$

A similar uniform limiting control can be undertaken by subtracting off analogous terms from the second integral in (44), i.e. the integral

$$
\frac{2 c_{1, d} \pi^{d / 2}}{|\log \varepsilon| \Gamma(d / 2)} \int_{1}^{\frac{1+\varepsilon}{|x|}} \frac{{ }_{2} F_{1}\left(\frac{d-1}{2}, \frac{1}{2} ; \frac{d}{2} ; r^{-2}\right)}{\left(r-\frac{1-\varepsilon}{|x|}\right)^{1 / 2}\left(\frac{1+\varepsilon}{|x|}-r\right)^{1 / 2}} \mathrm{~d} r
$$

Using (25), again noting $\alpha=1$, we can uniformly control the last term in (45) and note that it is $O(1 /|\log \varepsilon|)$. Similarly to (26), the second term in (45) has the same behaviour as

$$
\begin{equation*}
-\frac{2 c_{1, d} \pi^{d / 2}}{|\log \varepsilon| \Gamma((d-1) / 2) \Gamma(1 / 2)} \int_{0}^{1-\frac{1-\varepsilon}{|x|}} \frac{\log u}{\left(\frac{1+\varepsilon-|x|}{|x|}+u\right)^{1 / 2}\left(\frac{|x|-(1-\varepsilon)}{|x|}-u\right)^{1 / 2}} \mathrm{~d} u . \tag{46}
\end{equation*}
$$

To evaluate (46), using the change of variable $u=a-(a+b) /\left(t^{2}+1\right)$

$$
\begin{align*}
\int_{0}^{a} \frac{\log u}{\sqrt{(b+u)(a-u)}} \mathrm{d} u & =2 \int_{\sqrt{\frac{b}{a}}}^{\infty} \log \left(a-\frac{a+b}{t^{2}+1}\right) \frac{\mathrm{d} t}{t^{2}+1} \\
& =\int_{0}^{\arctan \sqrt{\frac{a}{b}}} \log \left(a-(a+b) \sin ^{2} w\right) \mathrm{d} w \\
& =\int_{0}^{\arctan \sqrt{\frac{a}{b}}} \log a+\log \left(1-\frac{\sin ^{2} w}{\frac{a}{a+b}}\right) \mathrm{d} w \\
& =\arctan \sqrt{\frac{a}{b}} \log (a+b)-L\left(\frac{\pi}{2}-2 \arctan \sqrt{\frac{a}{b}}\right)-\frac{\pi}{2} \log 2 \tag{47}
\end{align*}
$$

where we have used formula $4.226(5)$ of [15], which tells us that

$$
\begin{equation*}
\int_{0}^{u} \log \left(1-\frac{\sin ^{2} w}{\sin ^{2} v}\right) \mathrm{d} w=-u \log \sin ^{2} v-L\left(\frac{\pi}{2}-v+u\right)-L\left(\frac{\pi}{2}-v-u\right) \tag{48}
\end{equation*}
$$

for any $-\frac{\pi}{2} \leq u \leq \frac{\pi}{2}$ and $|\sin u| \leq|\sin v|$ where $L(x)$ is the Lobachevsky function. Note that, Lobachevsky's function is defined and represented as

$$
\begin{equation*}
L(x)=-\int_{0}^{x} \log \cos \theta \mathrm{~d} \theta=x \log 2-\frac{1}{2} \sum_{k=1}^{\infty}(-1)^{k-1} \frac{\sin 2 k x}{k^{2}} . \tag{49}
\end{equation*}
$$

Using (47) to evaluate (46) as well to evaluate the partner integral to (46), which comes from the analogous control of the second integral in (44), we get a nice cancellation of terms (as happened at this stage of the argument for $\alpha \in(0,1)$ ), to give us the controlled feature that

$$
\lim _{\varepsilon \rightarrow 0} \sup _{x \in \mathbb{S}_{\varepsilon}^{d-1}}\left|U \mu_{\varepsilon}^{(1)}(x)+\frac{2 c_{1, d} \pi^{d / 2}}{|\log \varepsilon| \Gamma((d-1) / 2) \Gamma(1 / 2)} \frac{\pi}{2} \log \varepsilon\right|=0 .
$$

Noting that with the indicated choice of $c_{1, d}$, we have

$$
\frac{2 c_{1, d} \pi^{d / 2}}{\Gamma((d-1) / 2) \Gamma(1 / 2)} \frac{\pi}{2}=1,
$$

which concludes the proof of part (i).
(ii) For the second part, the proof is almost identical to the proof of Lemma 2. Indeed, following the calculations through to (32), recalling that we have replaced $c_{\alpha, d}$ by $c_{1, d} /|\log \varepsilon|$, we get, up to an unimportant constant $\chi^{\prime}$,

$$
\begin{equation*}
\sup _{x \in \mathbf{S}_{\varepsilon}} U \mu_{\varepsilon, \delta}^{(2)}(x) \leq \chi^{\prime} \frac{1}{|\log \varepsilon| \delta^{d-1}} \tag{50}
\end{equation*}
$$

Hence, by taking $\delta=\delta(\varepsilon)=|\log \varepsilon|^{-1 / 2(d-1)}$ the statement of part (ii) follows.

With Lemma 3 in hand, we can now complete the proof of Theorem 2 (ii). Inequalities (34) and (37) are still at our disposal for the same reasons as before. The proof thus boils down to the asymptotic treatment of the term $U \mu_{\varepsilon}^{\mathrm{S}}(x)$ as in (40) for $x \notin \mathrm{~S}$. Recalling that we have replaced $c_{\alpha, d}$ by $c_{1, d} /|\log \varepsilon|$ we get from (31) and the constant $c_{1, d}$ given in the statement of Lemma 3,
$\lim _{\varepsilon \rightarrow 0}|\log \varepsilon| \mathbb{P}_{x}\left(\tau_{\mathrm{S}_{\varepsilon}}<\infty\right)=\frac{\Gamma((d-1) / 2)}{\pi^{(d+1) / 2}} \Gamma(1 / 2)^{2} H_{\mathrm{S}}(x)=\frac{\Gamma((d-1) / 2)}{\pi^{(d-1) / 2}} H_{\mathrm{S}}(x)$,
where we have used that $\Gamma(1 / 2)=\sqrt{\pi}$.

## 8 Proof of Theorem 1

Recall the definition $\tau_{\beta}:=\inf \left\{t>0: 1 / \beta<\left|X_{t}\right|<\beta\right\}$ for $\beta>1$ and fix $\varepsilon_{0}>0$ such that, for all $0<\varepsilon<\varepsilon_{0},(1-\varepsilon, 1+\varepsilon) \subset(1 / \beta, \beta)$. Then, by applying the Markov property at time $t$, we have, for any $\Lambda \in \mathcal{F}_{t}$,

$$
\begin{equation*}
\mathbb{P}_{x}\left(\Lambda, t<\tau_{\beta} \mid \tau_{\mathbf{S}_{\varepsilon}}<\infty\right)=\mathbb{E}_{x}\left[\mathbf{1}_{\left\{\Lambda, t<\tau_{\beta}\right\}} \frac{\mathbb{P}_{X_{t}}\left(\tau_{\mathrm{S}_{\varepsilon}}<\infty\right)}{\mathbb{P}_{x}\left(\tau_{\mathrm{S}_{\varepsilon}}<\infty\right)}\right] \tag{51}
\end{equation*}
$$

The event $\left\{t<\tau_{\beta}\right\}$ implies that either $\left|X_{t}\right|>\beta>1$ or $\left|X_{t}\right|<1 / \beta<1$. Hence, for all $0<\varepsilon<\varepsilon_{0}$ and $y \in \mathbb{S}_{\varepsilon}^{d-1}$, on $\left\{t<\tau_{\beta}\right\}$,

$$
\left|X_{t}-y\right|^{\alpha-d}<\max \left\{\left(\left(1-\varepsilon_{0}\right)-1 / \beta\right)^{\alpha-d},\left(\beta-\left(1+\varepsilon_{0}\right)\right)^{\alpha-d}\right\} .
$$

Hence, on $\left\{t<\tau_{\beta}\right\}$, we have from (37) and (41) that we can choose $\varepsilon$ sufficiently small such that

$$
\varepsilon^{\alpha-1} \mathbb{P}_{X_{t}}\left(\tau_{\mathbf{S}_{\varepsilon}}<\infty\right)<K_{1}
$$

for some constant $K_{1} \in(0, \infty)$. In a similar spirit, using (34) and (41), since $x \notin \mathrm{~S}$ and S is closed, it follows similarly that there is another constant $K_{2} \in(0, \infty)$ such that, for $x$ given in (51), we can choose $\varepsilon$ sufficiently small such that

$$
\varepsilon^{\alpha-1} \mathbb{P}_{x}\left(\tau_{\Omega_{\varepsilon}}<\infty\right)>K_{2}
$$

Theorem 2, dominated convergence and monotone convergence gives us, for all $\Lambda \in \mathcal{F}_{t}, t \geq 0$,

$$
\lim _{\beta \rightarrow 1} \lim _{\varepsilon \rightarrow 0} \mathbb{P}_{x}\left(\Lambda, t<\tau_{\beta} \mid \tau_{S_{\varepsilon}}<\infty\right)=\lim _{\beta \rightarrow 1} \mathbb{E}_{x}\left[\mathbf{1}_{\left\{\Lambda, t<\tau_{\beta}\right\}} \lim _{\varepsilon \rightarrow 0} \frac{\varepsilon^{\alpha-1} \mathbb{P}_{X_{t}}\left(\tau_{S_{\varepsilon}}<\infty\right)}{\varepsilon^{\alpha-1} \mathbb{P}_{x}\left(\tau_{S_{\varepsilon}}<\infty\right)}\right]=\mathbb{E}_{x}\left[\mathbf{1}_{\Lambda} \frac{H_{\mathrm{S}}\left(X_{t}\right)}{H_{\mathrm{S}}(x)}\right],
$$

as required.

## 9 Proof of Theorem 3

Recall the notation for a general Markov process $(Y, \mathrm{P})$ on $E$ preceding the statement of Theorem 3. We will additionally write $\mathcal{P}:=\left(\mathcal{P}_{t}, t \geq 0\right)$ for the semigroup associated to ( $Y, \mathrm{P}$ ).

Theorem 3.5 of Nagasawa [23], shows that, under suitable assumptions on the Markov process, $L$-times form a natural family of random times at which the pathwise time-reversal

$$
\overleftarrow{Y}_{t}:=Y_{(\mathrm{k}-t)-}, \quad t \in[0, \mathrm{k}]
$$

is again a Markov process. Let us state Nagasawa's principle assumptions.
(A) The potential measure $U_{Y}(a, \cdot)$ associated to $\mathcal{P}$, defined by the relation

$$
\begin{equation*}
\int_{E} f(x) U_{Y}(a, \mathrm{~d} x)=\int_{0}^{\infty} \mathcal{P}_{t}[f](a) \mathrm{d} t=\mathrm{E}_{a}\left[\int_{0}^{\infty} f\left(X_{t}\right) \mathrm{d} t\right], \quad a \in E, \tag{52}
\end{equation*}
$$

for bounded and measurable $f$ on $E$, is $\sigma$-finite. Assume that there exists a probability measure, $v$, such that, if we put

$$
\begin{equation*}
\mu(A)=\int U_{Y}(a, A) v(\mathrm{~d} a) \quad \text { for } A \in \mathcal{B}(E) \tag{53}
\end{equation*}
$$

then, there exists a Markov transition semigroup, say $\hat{\mathcal{P}}:=\left(\hat{\mathcal{P}}_{t}, t \geq 0\right)$ such that

$$
\begin{equation*}
\int_{E} \mathcal{P}_{t}[f](x) g(x) \mu(\mathrm{d} x)=\int_{E} f(x) \hat{\mathcal{P}}_{t}[g](x) \mu(\mathrm{d} x), \quad t \geq 0 \tag{54}
\end{equation*}
$$

for bounded, measurable and compactly supported test-functions $f, g$ on $E$.
(B) For any continuous test-function $f \in C_{0}(E)$, the space of continuous and compactly supported functions, and $a \in E$, assume that $\mathcal{P}_{t}[f](a)$ is rightcontinuous in $t$ for all $a \in E$ and, for $q>0, U_{\hat{Y}}^{(q)}[f]\left(\overleftarrow{Y}_{t}\right)$ is right-continuous in $t$, where, for bounded and measurable $f$ on $E$,

$$
U_{\hat{Y}}^{(q)}[f](a)=\int_{0}^{\infty} \mathrm{e}^{-q t} \hat{\mathcal{P}}_{t}[f](a) \mathrm{d} t, \quad a \in E,
$$

is the $q$-potential associated to $\hat{\mathcal{P}}$.
Nagasawa's duality theorem, Theorem 3.5. of [23], now reads as follows.
Theorem 5 (Nagasawa's duality theorem) Suppose that assumptions (A) and (B) hold. For the given starting probability distribution $v$ in $(A)$ and any $L$-time k , the time-reversed process $\overleftarrow{Y}$ under $\mathrm{P}_{v}$ is a time-homogeneous Markov process with transition probabilities

$$
\begin{equation*}
\mathrm{P}_{v}\left(\overleftarrow{Y}_{t} \in A \mid \overleftarrow{Y}_{r}, 0<r<s\right)=\mathrm{P}_{v}\left(\overleftarrow{Y}_{t} \in A \mid \overleftarrow{Y}_{s}\right)=p_{\hat{Y}}\left(t-s, \overleftarrow{Y}_{s}, A\right), \quad \mathrm{P}_{v} \text {-almost surely, } \tag{55}
\end{equation*}
$$

for all $0<s<t$ and closed $A$ in $E$, where $p_{\hat{Y}}(u, x, A), u \geq 0, x \in E$, is the transition measure associated to the semigroup $\hat{\mathcal{P}}$.

### 9.1 Completing the Proof of Theorem 3

We will make a direct application of Theorem 5, with $Y$ taken to be the process $\left(X, \mathbb{P}_{v}\right)$ where $v$ satisfies (8). Recall that its potential is written $U$ and we will denote its transition semigroup by ( $\mathcal{P}_{t}, t \geq 0$ ). Moreover, the dual process, formerly $\hat{Y}$, is taken to be $\left(X, \mathbb{P}^{S}\right)$ and we will, in the obvious way, work with the notation $U^{\mathrm{S}}$ in place of $U_{\hat{Y}}, \mathcal{P}$ in place of $\hat{\mathcal{P}}$ and so on. We need only to verify the two assumptions (A) and (B).

In order to verify (A), writing

$$
U(x, \mathrm{~d} y)=\int_{0}^{\infty} \mathbb{P}_{x}\left(X_{t} \in \mathrm{~d} y\right) \mathrm{d} t=\frac{\Gamma((d-\alpha) / 2)}{2^{\alpha} \pi^{d / 2} \Gamma(\alpha / 2)}|x-y|^{\alpha-d} \ell_{d}(\mathrm{~d} y), \quad x, y \in \mathbb{R}^{d}
$$

we have, up to a multiplicative constant,

$$
\begin{equation*}
\eta(\mathrm{d} x)=\int_{\mathbb{R}^{d}} U(a, \mathrm{~d} x) v(\mathrm{~d} a)=\frac{1}{\sigma_{1}(\mathrm{~S})} \int_{\mathrm{S}}|x-a|^{\alpha-d} \sigma_{1}(\mathrm{~d} a) \propto H_{\mathrm{S}}(x) \mathrm{d} x \tag{56}
\end{equation*}
$$

Now, we need to verify that (54) holds. Hunt's switching identity (cf. Chapter II. 1 of [4]) for $(X, \mathbb{P})$, states that

$$
\mathcal{P}_{t}(y, \mathrm{~d} x) \mathrm{d} y=\mathcal{P}_{t}(x, \mathrm{~d} y) \mathrm{d} x, \quad x, y \in \mathbb{R}^{d} .
$$

Using Hunt's switching identity together with (56), we have for $x, y \in \mathbb{R}^{d} \backslash \mathrm{~S}$
$\mathcal{P}_{t}(y, \mathrm{~d} x) \eta(\mathrm{d} y)=\mathcal{P}_{t}(y, \mathrm{~d} x) H_{\mathrm{S}}(y) \mathrm{d} y=\mathcal{P}_{t}(x, \mathrm{~d} y) \frac{H_{\mathrm{S}}(y)}{H_{\mathrm{S}}(x)} H_{\mathrm{S}}(x) \mathrm{d} x=\mathcal{P}_{t}^{\mathrm{S}}(x, \mathrm{~d} y) \eta(\mathrm{d} x)$.
Let us now turn to the verification of assumption (B). This assumption is immediately satisfied on account of the fact that $\mathcal{P}^{\mathcal{S}}$ is a right-continuous semigroup by virtue of its definition as a Doob $h$-transform with respect to the Feller semigroup $\mathcal{P}$ of the stable process.

With both (A) and (B) in hand, we are ready to apply Theorem 5 and the desired result thus follows.

## 10 Proof of Theorem 4

For the proof of Theorem 4, we focus on just part (i) and (ii) as the proof of parts (iii)-(v) are essentially verbatim the same as for the case of $S \in \mathbb{S}^{d-1}$. Moreover, for both parts (i) and (ii) we will provide only a sketch proof as the reader will quickly see that the proof is not hugely different form that of Theorem 2, albeit for a few technical details.
(i) The substance of the proof of part (i) is thus to follow a similar strategy as with Theorem 2 and build a measure $\rho_{\varepsilon}^{\mathrm{D}}$ such that the analogue of (16) holds, i.e. $U \rho_{\varepsilon}^{\mathrm{D}}(x)=1+o(1)$, for $x \in \mathrm{D}$ so that $(1+o(1)) \mathbb{P}_{x}\left(\tau_{\mathrm{D} \varepsilon}<\infty\right)=U \rho_{\varepsilon}^{\mathrm{D}}(x), x \notin \mathrm{D} \varepsilon$. More precisely, we develop analogues of Lemmas 1 and 2 to help make this precise.

Following what we have learned for $\mu_{\varepsilon}^{\mathrm{S}}$, our choice of $\rho_{\varepsilon}^{\mathrm{D}}$ is built from the base measure

$$
\begin{equation*}
\rho_{\varepsilon}(\mathrm{d} y)=k_{\alpha, d}((v, y)+\varepsilon)^{-\alpha / 2}(\varepsilon-(v, y))^{-\alpha / 2} \ell_{d}(\mathrm{~d} y) . \tag{57}
\end{equation*}
$$

for an appropriate choice of $k_{\alpha, d}$. As in (20) we can work with the decomposition,

$$
\begin{equation*}
U \rho_{\varepsilon}^{\mathrm{D}}(x)=U \rho_{\varepsilon}^{(1)}(x)-U \rho_{\varepsilon}^{(2)}(x), \quad x \in \mathrm{D}_{\varepsilon} \tag{58}
\end{equation*}
$$

where $\rho_{\varepsilon}^{(1)}$ (resp. $\rho_{\varepsilon}^{(2)}$ ) is the restriction of $\rho_{\varepsilon}$ to $\mathbb{H}_{\varepsilon}^{d-1}:=\left\{x \in \mathbb{R}^{d}:-\varepsilon<\right.$ $(v, x)<\varepsilon\}$ (resp. to $\hat{\mathrm{D}}_{\varepsilon}:=\mathbb{H}_{\varepsilon}^{d-1} \backslash \mathrm{D}_{\varepsilon}$ ). This helps with lower bounding $\liminf \varepsilon_{\varepsilon \rightarrow 0} \varepsilon^{\alpha-1} \mathbb{P}_{x}\left(\tau_{\mathrm{D}_{\varepsilon}}<\infty\right)$ by following steps of (33)-(35) together with the last paragraph of the Proof of Theorem 2, for which an analogue of Lemma 1 is needed.

For each $|u|<\varepsilon$, define the following sets: $\mathrm{D}^{\delta}=\left\{x \in \mathbb{H}^{d-1}: \operatorname{dist}(x, \mathrm{D})<\delta\right\}$, $\mathrm{D}_{\varepsilon}^{\delta}=\left\{y \in \mathbb{H}_{\varepsilon}^{d-1}: \hat{y} \in \mathrm{D}^{\delta}\right\}$ (recalling $\hat{y}$ is the orthogonal projection of $y$ on to $\mathbb{H}^{d-1}$ ) and $\hat{\mathrm{D}}_{\varepsilon}^{\delta}=\mathbb{H}_{\varepsilon}^{d-1} \backslash \mathrm{D}_{\varepsilon}^{\delta}$. Moreover, for any $u \in \mathbb{R}$, we define $\mathbb{H}^{d-1}(u)=\{x \in$ $\left.\mathbb{R}^{d}:(v, x)=u\right\}, \mathrm{D}^{\varepsilon}(u):=\left\{y \in \mathbb{H}^{d-1}(u): \hat{y} \in \mathrm{D}\right\}, \mathrm{D}^{\delta}(u):=\left\{y \in \mathbb{H}^{d-1}(u): \hat{y} \in\right.$ $\left.\mathrm{D}^{\delta}\right\}$, and $\hat{\mathrm{D}}^{\delta}(u)=\mathbb{H}^{d-1}(u) \backslash \mathrm{D}^{\delta}(u)$. Similarly, in the spirit of (21) we can use the decomposition

$$
\begin{equation*}
U \rho_{\varepsilon}^{\mathrm{D}^{\delta}}(x)=U \rho_{\varepsilon}^{(1)}(x)-U \rho_{\varepsilon, \delta}^{(2)}(x), \quad x \in \mathrm{D}_{\varepsilon}, \tag{59}
\end{equation*}
$$

where $\rho_{\varepsilon, \delta}^{(2)}$ is the restriction of $\rho_{\varepsilon}$ to $\hat{\mathrm{D}}_{\varepsilon}^{\delta}$. which helps with $\lim \sup _{\varepsilon \rightarrow 0} \varepsilon^{\alpha-1} \mathbb{P}_{x}\left(\tau_{\mathrm{D}_{\varepsilon}}<\right.$ $\infty$ ) by following steps (36)-(39) together with the last paragraph of the Proof of Theorem 2, for which an analogue of Lemma 2 is needed.

Let us address the technical detail that differs from the proof of Theorem 2 that we alluded to above. For $x \in \mathrm{D}_{\varepsilon}$,

$$
\begin{aligned}
& U \rho_{\varepsilon}^{(1)}(x) \\
& =k_{\alpha, d} \int_{\mathbb{H}_{\varepsilon}^{d-1}}|x-y|^{\alpha-d}((v, y)+\varepsilon)^{-\alpha / 2}(\varepsilon-(v, y))^{-\alpha / 2} \ell_{d}(\mathrm{~d} y) \\
& =k_{\alpha, d} \int_{-\varepsilon}^{\varepsilon}(u+\varepsilon)^{-\alpha / 2}(\varepsilon-u)^{-\alpha / 2} \mathrm{~d} u \int_{\mathbb{H}^{d-1}(u)}|x-y|^{\alpha-d} \ell_{d-1}(\mathrm{~d} y) \\
& =k_{\alpha, d} \int_{-\varepsilon}^{\varepsilon}(u+\varepsilon)^{-\alpha / 2}(\varepsilon-u)^{-\alpha / 2} \mathrm{~d} u \int_{\mathbb{H}^{d-1}((v, x))}\left(|x-\hat{z}|^{2}+|(v, x)-u|^{2}\right)^{\frac{\alpha-d}{2}} \ell_{d-1}(\mathrm{~d} \hat{z}),
\end{aligned}
$$

where $\hat{z}$ is the orthogonal projection of $y \in \mathbb{H}^{d-1}(u)$ onto $\mathbb{H}^{d-1}((v, x))$, which satisfies $|\hat{z}-y|=|(v, x)-u|$ and $\ell_{d-1}(\mathrm{~d} \hat{z})=\ell_{d-1}(\mathrm{~d} y)$. Note also that $(v, x-\hat{z})=$ 0 , for $\hat{z} \in \mathbb{H}^{d-1}((v, x))$, and hence $x-\mathbb{H}^{d-1}((v, x))$ is equal to $\mathbb{H}^{d-1}(0)$, which, in turn, can otherwise be identified as $\mathbb{R}^{d-1}$. Therefore, if we used generalised polar coordinates to integrate over $\mathbb{H}^{d-1}((v, x))$ identified as $x-\mathbb{R}^{d-1}$, we have

$$
\begin{align*}
& U \rho_{\varepsilon}^{(1)}(x) \\
& =k_{\alpha, d} \int_{-\varepsilon}^{\varepsilon}(u+\varepsilon)^{-\alpha / 2}(\varepsilon-u)^{-\alpha / 2} \mathrm{~d} u \int_{\mathbb{H}^{d-1}((v, x))}\left(|x-\hat{z}|^{2}+|(v, x)-u|^{2}\right)^{\frac{\alpha-d}{2}} \ell_{d-1}(\mathrm{~d} \hat{z}) \\
& =\frac{2 k_{\alpha, d} \pi^{(d-2) / 2}}{\Gamma((d-2) / 2)} \int_{-\varepsilon}^{\varepsilon}(u+\varepsilon)^{-\alpha / 2}(\varepsilon-u)^{-\alpha / 2} \mathrm{~d} u \int_{0}^{\infty} \int_{\mathbb{S}^{d-2}}\left(r^{2}+|(v, x)-u|^{2}\right)^{\frac{\alpha-d}{2}} r^{d-2} \mathrm{~d} r \sigma_{1}(d \theta) \\
& =\frac{2 k_{\alpha, d} \pi^{(d-2) / 2}}{\Gamma((d-2) / 2)} \int_{-\varepsilon}^{\varepsilon}(u+\varepsilon)^{-\alpha / 2}(\varepsilon-u)^{-\alpha / 2} \mathrm{~d} u \int_{0}^{\infty}\left(r^{2}+|(v, x)-u|^{2}\right)^{\frac{\alpha-d}{2}} r^{d-2} \mathrm{~d} r \\
& =\frac{k_{\alpha, d} \pi^{(d-2) / 2}}{\Gamma((d-2) / 2)} \int_{-\varepsilon}^{\varepsilon}(u+\varepsilon)^{-\alpha / 2}(\varepsilon-u)^{-\alpha / 2} \mathrm{~d} u \int_{0}^{\infty}\left(w+|(v, x)-u|^{2}\right)^{\frac{\alpha-d}{2}} w^{\frac{d-3}{2}} \mathrm{~d} w  \tag{60}\\
& =\frac{k_{\alpha, d} \pi^{(d-2) / 2} \Gamma\left(\frac{1-\alpha}{2}\right) \Gamma\left(\frac{d-1}{2}\right)}{\Gamma\left(\frac{d-2}{2}\right) \Gamma\left(\frac{d-\alpha}{2}\right)} \int_{-\varepsilon}^{\varepsilon}(u+\varepsilon)^{-\alpha / 2}(\varepsilon-u)^{-\alpha / 2}|(v, x)-u|^{\alpha-1} \mathrm{~d} u \\
& =\frac{k_{\alpha, d} \pi^{(d-2) / 2} \Gamma\left(\frac{1-\alpha}{2}\right) \Gamma\left(\frac{d-1}{2}\right)}{\Gamma\left(\frac{d-2}{2}\right) \Gamma\left(\frac{d-\alpha}{2}\right)} \int_{-1}^{1}(1+w)^{-\alpha / 2}(1-w)^{-\alpha / 2}\left|\varepsilon^{-1}(v, x)-w\right|^{\alpha-1} \mathrm{~d} w, \tag{61}
\end{align*}
$$

where, in the penultimate equality, we used a classical representation of the Beta function (see formula 3.191.2 in [15]), which tells us that, for any $\operatorname{Re}(v)>\operatorname{Re}(\gamma)>$ 0 and $z>0$,

$$
\int_{0}^{\infty}(y+z)^{-v} y^{\gamma-1} \mathrm{~d} y=z^{\gamma-v} \frac{\Gamma(v-\gamma) \Gamma(\gamma)}{\Gamma(v)}
$$

and in the final equality, we have changed variables using $w=\varepsilon u$. Next, we observe that $\left|\varepsilon^{-1}(v, x)\right| \leq 1$ on account of the fact that $x \in \mathrm{D}_{\varepsilon} \subseteq \mathbb{H}_{\varepsilon}^{d-1}$. Now choose $k_{\alpha, d}$, so that

$$
\begin{equation*}
\frac{k_{\alpha, d} \pi^{(d-2) / 2} \Gamma\left(\frac{1-\alpha}{2}\right) \Gamma\left(\frac{d-1}{2}\right)}{\Gamma\left(\frac{d-2}{2}\right) \Gamma\left(\frac{d-\alpha}{2}\right)}=1 \tag{62}
\end{equation*}
$$

We can now appeal directly to (18) to deduce that, for $x \in \mathrm{D}_{\varepsilon}$

$$
\begin{equation*}
U \rho_{\varepsilon}^{(1)}(x)=1 \tag{63}
\end{equation*}
$$

In the spirit of (33)-(35), it now follows that, for $x \notin \mathrm{D}$ and $\varepsilon$ sufficiently small,

$$
U \rho_{\varepsilon}^{\mathrm{D}}(x) \leq \mathbb{P}_{x}\left(\tau_{\mathrm{D}_{\varepsilon}}<\infty\right)
$$

So that

$$
\begin{equation*}
\liminf _{\varepsilon \rightarrow 0} U \rho_{\varepsilon}^{\mathrm{D}}(x) \leq \liminf _{\varepsilon \rightarrow 0} \mathbb{P}_{x}\left(\tau_{\mathrm{D}_{\varepsilon}}<\infty\right), \quad x \notin \mathrm{D} \tag{64}
\end{equation*}
$$

Now we turn our attention to (59). Noting that when $x \in \mathrm{D}_{\varepsilon},|x-y|>\delta$ for $y \in \hat{\mathrm{D}}_{\varepsilon}^{\delta}$, we have, for all $x \in \mathrm{D}_{\varepsilon}$,

$$
\begin{aligned}
U \rho_{\varepsilon, \delta}^{(2)}(x) & =k_{\alpha, d} \int_{\hat{\mathrm{D}}_{\varepsilon}^{\delta}}|x-y|^{\alpha-d}((v, y)+\varepsilon)^{-\alpha / 2}(\varepsilon-(v, y))^{-\alpha / 2} \ell_{d}(\mathrm{~d} y) \\
& \leq k_{\alpha, d} \delta^{\alpha-d} \int_{-\varepsilon}^{\varepsilon}(u+\varepsilon)^{-\alpha / 2}(\varepsilon-u)^{-\alpha / 2} \mathrm{~d} u \int_{\mathrm{D}^{\delta}((v, x))} \ell_{d-1}(\mathrm{~d} \hat{y}) \\
& \leq \delta^{\alpha-d} k_{\alpha, d} \ell_{d-1}\left(\mathrm{D}^{\delta}\right) \int_{-\varepsilon}^{\varepsilon}(u+\varepsilon)^{-\alpha / 2}(\varepsilon-u)^{-\alpha / 2} \mathrm{~d} u \\
& =\delta^{\alpha-d} \varepsilon^{1-\alpha} k_{\alpha, d} \ell_{d-1}\left(\mathrm{D}^{\delta}\right) 2^{1-\alpha} \frac{\Gamma((2-\alpha) / 2)^{2}}{\Gamma(2-\alpha)},
\end{aligned}
$$

where we have used the calculation in (31) in the final equality. Choosing $\delta=$ $\delta(\varepsilon)=\varepsilon^{(1-\alpha) / 2(d-\alpha)}$, and noting that $\ell_{d-1}\left(\mathrm{D}^{\delta}\right)$ is uniformly bounded from above by an unimportant constant for e.g. all $\delta<1$ (thanks to the assumption that $\ell_{d-1}$ (D) $<$ $\infty$ ), we see that

$$
\lim _{\varepsilon \rightarrow 0} \sup _{x \in \mathrm{D}_{\varepsilon}} U \rho_{\varepsilon, \delta(\varepsilon)}^{(2)}(x)=0 .
$$

In a similar spirit to (36)-(38), we now have that

$$
\begin{equation*}
\limsup _{\varepsilon \rightarrow 0} \varepsilon^{\alpha-1} U \rho_{\varepsilon}^{\mathrm{D}^{\delta(\varepsilon)}}(x) \geq \limsup _{\varepsilon \rightarrow 0} \varepsilon^{\alpha-1} \mathbb{P}_{x}\left(\tau_{\mathrm{D}_{\varepsilon}}<\infty\right), \quad x \notin \mathrm{D} . \tag{65}
\end{equation*}
$$

Matching up the left-hand side of (64) with that of (65), we can proceed in a similar fashion to (41)-(42), leading to the statement of Theorem 4(i) as promised. The calculation is based around the fact that

$$
\begin{align*}
\lim _{\varepsilon \rightarrow 0} \varepsilon^{\alpha-1} U \rho_{\varepsilon}^{\mathrm{D}}(x) & =\lim _{\varepsilon \rightarrow 0} k_{\alpha, d} \varepsilon^{\alpha-1} \int_{\mathrm{D}_{\varepsilon}}|x-y|^{\alpha-d}((v, y)+\varepsilon)^{-\alpha / 2}(\varepsilon-(v, y))^{-\alpha / 2} \ell_{d}(\mathrm{~d} y) \\
& =\lim _{\varepsilon \rightarrow 0} k_{\alpha, d} \varepsilon^{\alpha-1} \int_{-\varepsilon}^{\varepsilon}(u+\varepsilon)^{-\alpha / 2}(\varepsilon-u)^{-\alpha / 2} \mathrm{~d} u \int_{\mathrm{D}(u)}|x-\hat{y}|^{\alpha-d} \ell_{d-1}(\mathrm{~d} \hat{y}) \\
& =k_{\alpha, 2^{2}} 1^{1-\alpha} \frac{\Gamma((2-\alpha) / 2)^{2}}{\Gamma(2-\alpha)} \int_{\mathrm{D}}|x-y|^{\alpha-d} \ell_{d-1}(\mathrm{~d} y) \\
& =2^{1-\alpha} \pi^{-(d-2) / 2} \frac{\Gamma\left(\frac{d-2}{2}\right) \Gamma\left(\frac{d-\alpha}{2}\right) \Gamma\left(\frac{2-\alpha}{2}\right)^{2}}{\Gamma\left(\frac{1-\alpha}{2}\right) \Gamma\left(\frac{d-1}{2}\right) \Gamma(2-\alpha)} \int_{\mathrm{D}}|x-y|^{\alpha-d} \ell_{d-1}(\mathrm{~d} y), \tag{66}
\end{align*}
$$

where we have used the calculation in (31) and (62) in the third equality.
(ii) The setting $\alpha=1$ requires yet another delicate handling of the associated potentials. Given that all the main ideas are now present in the paper, we simply lay out the key points of the proof, leaving the remaining detail for the reader.

Our calculations begin in the same way as in part (i), in particular, we work with the core measure $\rho_{\varepsilon}$ as in (57), albeit (as with Theorem 2(ii)) replacing $k_{1, d}$ by $k_{1, d} /|\log \varepsilon|$, to be used in the constructions (58) and (59). An immediate complication we have is in evaluating $U \rho_{\varepsilon}^{(1)}(x)$, for $x \in \mathrm{D}_{\varepsilon}$, can be seen when we pick up the computations for part (i) at (60). Indeed, at that point, we are confronted with the integral

$$
\int_{0}^{\infty}\left(w+|(v, x)-u|^{2}\right)^{\frac{1-d}{2}} w^{\frac{d-3}{2}} \mathrm{~d} w=\infty
$$

The solution to this is to adjust the core measure $\rho_{\varepsilon}$ as follows. Since D is bounded, we can choose an $R>0$ sufficiently large that, $\mathrm{D} \subset \mathbb{S}^{d-2}(0, R):=\left\{y \in \mathbb{H}^{d-1}\right.$ : $|y| \leq R\}$ strictly contains $D$. Denote $\mathbb{S}_{\varepsilon}^{d-2}(0, R)=\left\{x \in \mathbb{H}_{\varepsilon}^{d-1}: \hat{x} \in \mathbb{S}^{d-2}(0, R)\right\}$, where $\hat{x}$ is the orthogonal projection of $x$ on to $\mathbb{H}^{d-1}$. Suppose we now make a slight adjustment and replace $\rho_{\varepsilon}$ by

$$
\rho_{\varepsilon}(\mathrm{d} y)=\frac{k_{1, d, R}}{|\log \varepsilon|}((v, y)+\varepsilon)^{-\alpha / 2}(\varepsilon-(v, y))^{-\alpha / 2} \mathbf{1}_{\left(y \in \mathbb{S}_{\varepsilon}^{d-2}(0, R)\right)} \ell_{d}(\mathrm{~d} y),
$$

for an appropriate choice of $k_{1, d, R}$. We may now continue the argument from (60) with the calculation

$$
\begin{equation*}
|\log \varepsilon| U \rho_{\varepsilon}^{(1)}(x)=\frac{k_{1, d, R} \pi^{(d-2) / 2}}{\Gamma((d-2) / 2)} \int_{-\varepsilon}^{\varepsilon}(u+\varepsilon)^{-1 / 2}(\varepsilon-u)^{-1 / 2} \mathrm{~d} u \int_{0}^{R}\left(w+|(v, x)-u|^{2}\right)^{\frac{1-d}{2}} w^{\frac{d-3}{2}} \mathrm{~d} w . \tag{67}
\end{equation*}
$$

Denote by $I(R, \varepsilon, x)$ the right-hand side of 67 , ensuring that $\varepsilon$ is small enough that $\varepsilon \ll R$.

Appealing to (78),

$$
\begin{aligned}
& I(R, \varepsilon, x)= \frac{k_{1, d, R} \pi^{(d-2) / 2}}{\Gamma((d-2) / 2)} \int_{-\varepsilon}^{\varepsilon}(u+\varepsilon)^{-1 / 2}(\varepsilon-u)^{-1 / 2} \mathrm{~d} u \int_{0}^{R}\left(w+|(v, x)-u|^{2}\right)^{\frac{1-d}{2}} w^{\frac{d-3}{2}} \mathrm{~d} w \\
&= \frac{k_{1, d, R} \pi^{(d-2) / 2}}{\Gamma((d-2) / 2)} \int_{-\varepsilon}^{\varepsilon}\left(\varepsilon^{2}-u^{2}\right)^{-1 / 2}|(v, x)-u|^{1-d} \mathrm{~d} u \int_{0}^{R}\left(\frac{w}{|(v, x)-u|^{2}}+1\right)^{\frac{1-d}{2}} w^{\frac{d-3}{2}} \mathrm{~d} w \\
&= \frac{k_{1, d, R} \pi^{(d-2) / 2}}{\Gamma((d-2) / 2)} \int_{-\varepsilon}^{\varepsilon}\left(\varepsilon^{2}-u^{2}\right)^{-1 / 2}|(v, x)-u|^{1-d} \\
& \quad \frac{R^{(d-1) / 2}}{(d-1) / 2} 2 F_{1}\left(\frac{d-1}{2}, \frac{d-1}{2} ; \frac{d+1}{2} ;-\frac{R}{|(v, x)-u|^{2}}\right) \mathrm{d} u,
\end{aligned}
$$

where we have used the identity in (79). One of the many identities for hypergeometric functions, see [2], offers us the growth condition, for $c-a \in \mathbb{N}$, as $|z| \rightarrow \infty$,

$$
\begin{equation*}
{ }_{2} F_{1}(a, a ; c ; z) \sim \frac{\Gamma(c)(\log (-z)-\psi(c-a)-\psi(a)-2 \gamma)(-z)^{-a}}{\Gamma(a)(c-a-1)!}+\frac{\Gamma(c) 2(-z)^{-c}}{\Gamma(a)^{2}((c-a)!)^{2}}, \tag{68}
\end{equation*}
$$

where $\gamma$ is an unimportant constant and $\psi(z)=\Gamma^{\prime}(z) / \Gamma(z)$ is the di-gamma function. In the spirit of previous calculations, we can thus find to leading order, uniformly over $x \in \mathrm{D}_{\varepsilon}$,

$$
\begin{equation*}
U \rho_{\varepsilon}^{(1)}(x) \sim 2 \frac{\pi^{d / 2} k_{1, d, R}}{\Gamma((d-2) / 2)} \tag{69}
\end{equation*}
$$

which remarkably does not depend on $R$. This means we should choose the constant

$$
k_{1, d, R}=\frac{\Gamma((d-2) / 2)}{2 \pi^{d / 2}}
$$

for this asymptotic to serve our purpose.
At this point in the proof, recalling the fundamental decomposition (58), it is worth bringing in the term $U \mu_{\varepsilon}^{(2)}$ and noting that one can compute with relatively coarse estimates that

$$
\sup _{x \in \mathrm{D}_{\varepsilon}}\left|U \rho_{\varepsilon}^{(2)}(x)\right| \leq \frac{C}{|\log \varepsilon|},
$$

for some unimportant constant $C>0$. Together with (69), in a calculation similar to (66) we can put the pieces together to get the asymptotic, for $x \notin \mathrm{D}$ and $\varepsilon$ sufficiently small,

$$
\begin{align*}
\lim _{\varepsilon \rightarrow 0}|\log \varepsilon| \mathbb{P}_{x}\left(\tau_{\mathrm{D}_{\varepsilon}}<\infty\right) & =\lim _{\varepsilon \rightarrow 0}|\log \varepsilon| U \rho_{\varepsilon}^{\mathrm{D}}(x) \\
& =\lim _{\varepsilon \rightarrow 0} \frac{\Gamma((d-2) / 2)}{2 \pi^{d / 2}} \int_{\mathrm{D}_{\varepsilon}}|x-y|^{1-d}\left(\varepsilon^{2}-(v, y)^{2}\right)^{-1 / 2} \ell_{d}(\mathrm{~d} y) \\
& =\lim _{\varepsilon \rightarrow 0} \frac{\Gamma((d-2) / 2)}{\pi^{d / 2}} \int_{-\varepsilon}^{\varepsilon}\left(\varepsilon^{2}-u^{2}\right)^{-1 / 2} \mathrm{~d} u \int_{\mathrm{D}(u)}|x-\hat{y}|^{1-d} \ell_{d-1}(\mathrm{~d} \hat{y}) \\
& =\frac{\Gamma((d-2) / 2)}{\pi^{(d-2) / 2}} M_{\mathrm{D}}(x) . \tag{70}
\end{align*}
$$

The proof is complete.

## Appendix: Hypergeometric Identities

We work with the standard definition for the hypergeometric function,

$$
{ }_{2} F_{1}(a, b, c ; z)=\sum_{n=0}^{\infty} \frac{(a)_{n}(b)_{n}}{(c)_{n}} \frac{z^{n}}{n!}, \quad|z|<1
$$

Of the many identities for hypergeometric functions, we need the following:

$$
\begin{gather*}
{ }_{2} F_{1}(a, b, c ; z)=\frac{\Gamma(c) \Gamma(a+b-c)}{\Gamma(a) \Gamma(b)}(1-z)^{c-a-b}{ }_{2} F_{1}(c-a, c-b, 1+c-a-b ; 1-z) \\
+\frac{\Gamma(c) \Gamma(c-a-b)}{\Gamma(c-a) \Gamma(c-b)}{ }_{2} F_{1}(a, b, a+b-c+1 ; 1-z), \tag{71}
\end{gather*}
$$

for $c-a-b \notin \mathbb{Z}$. Hence, thanks to continuity,

$$
\begin{align*}
& \lim _{\varepsilon \rightarrow 0} \sup _{r \in[1-\varepsilon, 1]} \left\lvert\, 2 F_{1}\left(\frac{d-\alpha}{2}, 1-\frac{\alpha}{2} ; \frac{d}{2} ; r^{2}\right)\right. \\
& \left.-\frac{\Gamma(d / 2) \Gamma(1-\alpha)}{\Gamma((d-\alpha) / 2) \Gamma((2-\alpha) / 2)}\left(1-r^{2}\right)^{\alpha-1}-\frac{\Gamma(d / 2) \Gamma(\alpha-1)}{\Gamma(\alpha / 2) \Gamma((d+\alpha-2) / 2)} \right\rvert\,=0 . \tag{72}
\end{align*}
$$

We will need to apply a similar identity to (71) but for the setting that $c-a-b=$ 0 , which violates the assumption behind (71). In that case, we need to appeal to the formula

$$
\begin{align*}
{ }_{2} F_{1}(a, b, a+b, z) & =\frac{\Gamma(a+b)}{\Gamma(a) \Gamma(b)}\left(\sum_{k=0}^{\infty} \frac{(a)_{k}(b)_{k}}{(k!)^{2}}(2 \psi(k+1)-\psi(a+k)-\psi(b+k))(1-z)^{k}\right. \\
& \left.-\log (1-z){ }_{2} F_{1}(a, b, 1,1-z)\right), \tag{73}
\end{align*}
$$

for $|1-z|<1$ where the di-gamma function $\psi(z)=\Gamma^{\prime}(z) / \Gamma(z)$ is defined for all $z \neq-n, n \in \mathbb{N}$.

Again, thanks to continuity, we can write

$$
\begin{align*}
\lim _{\varepsilon \rightarrow 0} \sup _{r \in[1-\varepsilon, 1]} \left\lvert\, 2 F_{1}\left(\frac{d-1}{2}, \frac{1}{2} ; \frac{d}{2} ; r^{2}\right)\right. & +\frac{\Gamma(d / 2)}{\Gamma((d-1) / 2) \Gamma(1 / 2)} \log \left(1-r^{2}\right) \\
& \left.-\frac{2 \Gamma(d / 2)(\psi(1)-\psi((d-1) / 2)-\psi(1 / 2))}{\Gamma((d-1) / 2) \Gamma(1 / 2)} \right\rvert\,=0 \tag{74}
\end{align*}
$$

A second identity that is needed is the following combination formula, which states that for any $|z|<1$, we have

$$
\begin{align*}
&{ }_{2} F_{1}(a, b ; c ; z)=\frac{\Gamma(b-a) \Gamma(c)}{\Gamma(c-a) \Gamma(b)}(-z)^{-a}{ }_{2} F_{1}\left(a, a-c+1 ; a-b+1 ; \frac{1}{z}\right) \\
&+\frac{\Gamma(a-b) \Gamma(c)}{\Gamma(c-b) \Gamma(a)}(-z)^{-b}{ }_{2} F_{1}\left(b-c+1, b ;-a+b+1 ; \frac{1}{z}\right) . \tag{75}
\end{align*}
$$

Its proof can be found, for example at [1]. In the main body of the text, we use this identity for the setting that $a=\alpha / 2, b=\alpha$ and $c=1+\alpha / 2$. This gives us the identity

$$
\begin{aligned}
{ }_{2} F_{1}\left(\frac{\alpha}{2}, \alpha ; 1+\frac{\alpha}{2} ; z\right)= & \frac{\Gamma(\alpha / 2) \Gamma((2+\alpha) / 2)}{\Gamma(\alpha)}(-z)^{-\alpha / 2}{ }_{2} F_{1}\left(\alpha / 2,0 ; 1-\alpha / 2 ; \frac{1}{z}\right) \\
& \quad+\frac{\Gamma(-\alpha / 2) \Gamma((2+\alpha) / 2)}{\Gamma((2-\alpha) / 2) \Gamma(\alpha / 2)}(-z)^{-\alpha}{ }_{2} F_{1}\left(\alpha / 2, \alpha ; 1+\alpha / 2 ; \frac{1}{z}\right) \\
= & \frac{\Gamma(\alpha / 2) \Gamma((2+\alpha) / 2)}{\Gamma(\alpha)}(-z)^{-\alpha / 2} \\
& \quad-(-z)^{-\alpha}{ }_{2} F_{1}\left(\alpha / 2, \alpha ; 1+\alpha / 2 ; \frac{1}{z}\right),
\end{aligned}
$$

where we have used the recursion formula for gamma functions twice in the final equality. This allows us to come to rest at the following useful identity

$$
\begin{equation*}
(-z)^{-\alpha / 2}{ }_{2} F_{1}\left(\alpha / 2, \alpha ; 1+\alpha / 2 ; \frac{1}{z}\right)+(-z)^{\alpha / 2}{ }_{2} F_{1}\left(\frac{\alpha}{2}, \alpha ; 1+\frac{\alpha}{2} ; z\right)=\frac{\Gamma(\alpha / 2) \Gamma((2+\alpha) / 2)}{\Gamma(\alpha)} . \tag{76}
\end{equation*}
$$

We are also interested in integral formulae, for which the hypergeometric function is used to evaluate an integral. The first is aversion of formula 3.665(2) in [15] which states that, for any $0<|a|<r$ and $v>0$, as

$$
\begin{equation*}
\int_{0}^{\pi} \frac{\sin ^{d-2} \phi}{\left(a^{2}+2 a r \cos \phi+r^{2}\right)^{\nu}} \mathrm{d} \phi=\frac{1}{r^{2 v}} B\left(\frac{d-1}{2}, \frac{1}{2}\right)_{2} F_{1}\left(v, v-\frac{d}{2}+1 ; \frac{d}{2} ; \frac{a^{2}}{r^{2}}\right), \tag{77}
\end{equation*}
$$

where $B(a, b)=\Gamma(a) \Gamma(b) / \Gamma(a+b)$ is the Beta function. The second is formula 3.197.8 in [15], which states that, for $\operatorname{Re}(\mu)>0, \operatorname{Re}(\nu)>0$ and $|\arg (u / \beta)|<\pi$,

$$
\begin{equation*}
\int_{0}^{u} x^{\nu-1}(u-x)^{\mu-1}(x+\beta)^{\lambda} \mathrm{d} x=\beta^{\lambda} u^{\mu+\nu-1} B(\mu, \nu)_{2} F_{1}\left(-\lambda, \nu ; \mu+\nu ;-\frac{u}{\beta}\right) . \tag{78}
\end{equation*}
$$

The third is 3.194.1 of [15] and states that, for $|\arg (1+\beta u)|>\pi$ and $\operatorname{Re}(\mu)>0$, $\operatorname{Re}(v)>0$,

$$
\begin{equation*}
\int_{0}^{u} x^{\mu-1}(1+\beta x)^{-v} \mathrm{~d} x=\frac{u^{\mu}}{\mu} 2 F_{1}(v, v-\mu ; 1+\mu ;-\beta u), \tag{79}
\end{equation*}
$$

where ${ }_{2} F_{1}$ in the above identity is understood as its analytic extension in the event that $|\beta u|>1$.

Acknowledgments TS acknowledges support from a Schlumberger Faculty of the Future award. SP acknowledges support from the Royal Society as a Newton International Fellow Alumnus (AL201023) and UNAM-DGAPA-PAPIIT grant no. IA103220.

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[^1]:    ${ }^{1}$ We will distinguish integrals with respect to one-dimensional Lebesgue measure as taking the form $\int \cdot \mathrm{d} x$, where as higher dimensional integrals will always indicate the dimension, for example $\int \cdot \ell_{d}(\mathrm{~d} x)$.

