

# Explicit Convergence Rates of Underdamped Langevin Dynamics under weighted and Weak Poincaré–Lions Inequalities

**Andi Q. Wang (University of Warwick)**

Joint work with Giovanni Brigati (IST Austria), Gabriel Stoltz (Ecole des Ponts ParisTech) and Lihan Wang (Carnegie Mellon University)

Recent advances in Stochastic Algorithms

Clermont-Ferrand

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**Sampling problem:** Given a target probability measure  $\mu$  on  $\mathbb{R}^d$  defined by

$$\mu(dx) = \frac{1}{Z} e^{-\phi(x)} dx$$

with some unknown normalization constant  $Z$ , generate a sequence of random samples  $\{x_i\} \sim \mu$ .

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Applications:

- Bayesian inference
- Molecular dynamics
- Machine learning
- Economics (Pareto distribution,  $\psi(v)$  represents the money of agent)

- Overdamped Langevin Dynamics

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$$\partial_t h_t = \Delta_x h_t - \nabla_x \phi \cdot \nabla_x h_t \quad =: -\nabla_x^* \nabla_x h_t.$$

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- **Question: How fast does it converge?** Will study the continuous-time process using **functional inequalities**.

# Poincaré Inequality

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**Question: What if (PI) fails?**

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Assume moment bound

$$\sup_{t>0} \int W^\sigma h_t^2 d\mu < \infty,$$

use Hölder inequality

$$\begin{aligned} \text{Var}_\mu(h_t) &\leq \left( \int W^\sigma h_t^2 d\mu \right)^{\frac{2}{2+\sigma}} \left( \int W^{-2} h_t^2 d\mu \right)^{\frac{\sigma}{2+\sigma}} \\ &\lesssim \mathbb{E}_\mu^{\frac{\sigma}{2+\sigma}} |\nabla_x f|^2 \lesssim \left( -\frac{d}{dt} \text{Var}_\mu(h_t) \right)^{\frac{\sigma}{2+\sigma}} \end{aligned}$$

Solving ODE, we get **algebraic** convergence

$$\text{Var}_\mu(h_t) \lesssim t^{-\frac{\sigma}{2}}.$$

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$$\mathrm{Var}_\mu(f) \leq s \mathbb{E}_\mu |\nabla_x f|^2 + \beta(s) \|f\|_{\mathrm{osc}}^2, \quad \forall s > 0. \quad (\text{WPI})$$

$\|f\|_{\mathrm{osc}} = \mathrm{ess\,sup}_\mu f - \mathrm{ess\,inf}_\mu f$ ,  $\beta(s) \downarrow 0$  as  $s \rightarrow \infty$ .

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By maximum principle  $\|h_t\|_{\mathrm{osc}} \leq \|h_0\|_{\mathrm{osc}}$ . Thus (WPI) implies

$$\frac{\mathbb{E}_\mu |\nabla_x h_t|^2}{\|h_0\|_{\mathrm{osc}}^2} \geq \frac{\mathrm{Var}_\mu(h_t)}{s \|h_0\|_{\mathrm{osc}}^2} - \frac{\beta(s)}{s}.$$

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Optimizing over  $s$ , we find an increasing function  $K^*$ , depending on  $\beta$ , such that

$$-\frac{1}{2} \frac{d}{dt} \frac{\mathrm{Var}_\mu(h_t)}{\|h_0\|_{\mathrm{osc}}^2} = \frac{\mathbb{E}_\mu |\nabla_x h_t|^2}{\|h_0\|_{\mathrm{osc}}^2} \geq K^* \left( \frac{\mathrm{Var}_\mu(h_t)}{\|h_0\|_{\mathrm{osc}}^2} \right).$$

Explicit decay rates can be obtained through Bihari–Lasalle argument.

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Now take any  $s > 0$ ,

$$\int (f - \mu(f))^2 d\mu \leq s \int_{\{s > W^2\}} W^{-2}(f - \mu(f))^2 d\mu + \int_{\{s \leq W^2\}} (f - \mu(f))^2 d\mu$$

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⇒ **Weighted Poincaré implies weak Poincaré!** [Cattiaux/Goźlan/Guillin/Roberto 2010].

# Examples (convergence of overdamped Langevin)

Example 1:  $\phi(x) = \langle x \rangle^\alpha := (1 + |x|^2)^{\frac{\alpha}{2}}$  with  $\alpha \in (0, 1)$ .

- weighted Poincaré  $\Rightarrow W(x) = \langle x \rangle^{1-\alpha}$ ;
- weak Poincaré  $\Rightarrow \beta(s) \sim \mu(s \leq W^2) \sim \exp(-s^{\frac{\alpha}{2-2\alpha}})$   
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These tail bounds match the optimal convergence rates in total variation [Brešar/Mijatovic '24].

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$$\begin{cases} dX_t = \nabla_v \psi(V_t) dt; \\ dV_t = -\nabla_x \phi(X_t) - \gamma \nabla_v \psi(V_t) + \sqrt{2\gamma} dW_t. \end{cases}$$

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- Natural choice  $\psi(v) = \frac{|v|^2}{2}$  standard Gaussian.
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- Strong second-order lift of overdamped Langevin dynamics [Eberle/Lörler '24] if  $\mathbb{E}_\nu \nabla_v^2 \psi = \text{Id}$
- $h_t = \mathbb{E}_{x,v} h_0(X_t, V_t)$  satisfies

$$\begin{aligned} \partial_t h_t &= \nabla_v \psi \cdot \nabla_x h_t - \nabla_x \phi \cdot \nabla_v h_t + \gamma \Delta_v h_t - \gamma \nabla_v \psi \cdot \nabla_v h_t \\ &=: \mathcal{T} h_t - \gamma \nabla_v^* \nabla_v h_t. \end{aligned}$$

Energy estimate

$$\frac{d}{dt} \mathbb{E}_{\mu \otimes \nu} \text{Var}(h_t) = -2\gamma \mathbb{E}_{\mu \otimes \nu} |\nabla_{\nu} h_t|^2.$$

Lacking diffusion in  $x$  variable!

# Nonreversible dynamics

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Lacking diffusion in  $x$  variable! This will be 0 if  $h_t$  only depends on  $x$ .

So there is no hope to use standard Poincaré inequality, where we had

$$\frac{d}{dt} \text{Var}_{\mu}(h_t) = -2\mathbb{E}_{\mu} |\nabla_x h_t|^2 \leq -2C_P \text{Var}_{\mu}(h_t).$$

The dynamics are **hypocoercive**.

# Previous works on Hypocoercivity: Strong Confinement

- Early works of [Kolmogorov 1934], [Hörmander 1967];
- Convergence using Lyapunov function approach [Wu '01; Mattingly/Stuart/Higham '02; Rey-Bellet '06];
- Convergence in  $H^1$  norm [Talay 2002; Villani 2009];
- Convergence in a modified  $L^2$  norm [Dolbeault/Mouhot/Schmeiser '09; '15] (also earlier idea from [Herau '06]);
- Convergence in Wasserstein distance: using Bakry-Émery framework [Baudoin '16]; by coupling approaches [Eberle/Guillin/Zimmer '19];
- Resolvent analysis using Schur Complement [Bernard/Fathi/Levitt/Stoltz '20];
- **Space-time Poincaré inequality [Albritton/Armstrong/Mourrat/Novack '19; Cao/Lu/Wang '20; Brigati/Stoltz '23; Eberle/Lörler '24]**  
Optimal scaling of friction and convergence rate for convex  $\phi$ .
- See also recent [Eberle/Guillin/Hahn/Lörler/Michael '25].

# Previous works on Hypocoercivity: Weak Confinement

- Combination of weak Poincaré with hypocoercivity tools: [Hu/Wang '19] using Villani's techniques; [Grothaus/Wang '19] using the DMS framework; [Andrieu/Dobson/W. '21] for PDMPs;
- Weighted Poincaré with DMS framework: [Cao '19] stretched exponential spatial potential and interpolation of spaces; [Bouin/Dolbeault/Ziviani '23] algebraic decays;
- Space-time Poincaré framework of Armstrong/Mourrat: [Brigati '23; Brigati/Stoltz '23], weak confinement in velocity, but requires strong confinement in space; [Dietert '23] allows non-gradient drift, although explicit estimates are absent.

These works require  $\nu$  to be Gaussian.

# The $L^2$ hypocoercivity approach: DMS

Assume (PI) in both  $\mu$  and  $\nu$ . Perturb the  $L^2$  energy:

$$\Phi(f) := \frac{1}{2} \|f\|^2 + \varepsilon \langle (1 + (\mathcal{T}\Pi_\nu)^*(\mathcal{T}\Pi_\nu))^{-1} (\mathcal{T}\Pi_\nu)^* f, f \rangle.$$

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Differentiate along the dynamics, we get for small  $\varepsilon$  and some  $\lambda = \lambda_\varepsilon$

$$\frac{d}{dt} \Phi(h_t) \leq -\lambda \Phi(h_t).$$

Convergence rate is difficult to optimize in  $\varepsilon$ .

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Fix time  $\tau > 0$ , and consider the **time-averaged** energy  $\mathcal{H}_\tau(t) := \int_t^{t+\tau} \|h_t\|_{L^2}^2 dt$ .

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- Poincaré–Lions inequality in **space and time**

$$\|f - (f)\|_{L^2([0,\tau] \times \mu)} \lesssim \|(\nabla_x, \partial_t)f\|_{H^{-1}([0,\tau] \times \mu)}.$$

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- Averaging functional inequality in **velocity**

$$\begin{aligned} \|(\nabla_x \Pi_\nu h, \partial_t \Pi_\nu h)\|_{H^{-1}([0,\tau] \times \mu)} &\lesssim \|h - \Pi_\nu h\|_{L^2([0,\tau] \times \mu \times \nu)} \\ &\quad + \|(\partial_t + \mathcal{T})h\|_{L^2([0,\tau] \times \mu, H^{-1}(\nu))}. \end{aligned}$$

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- Use the equation, dissipation of **time-averaged**  $L^2$  energy controls the  $L^2$ -energy itself. Convergence rate  $O(\sqrt{C_P})$  for convex  $\phi$  when taking  $\gamma = O(\sqrt{C_P})$  [Cao/Lu/Wang '20], optimal second-order lift [Eberle/Lörler '24].

# Poincaré–Lions Inequality

Assume that  $\mu$  satisfies (PI), then the space-time strip  $([0, \tau] \times \mathbb{R}^d, dt \otimes d\mu)$  satisfies the Poincaré–Lions inequality: for any  $f(t, x)$ ,

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This is seemingly stronger than tensorized Poincaré inequality ( $H^{-1}$  replaced by  $L^2$ ), but they are equivalent.

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This is seemingly stronger than tensorized Poincaré inequality ( $H^{-1}$  replaced by  $L^2$ ), but they are equivalent.

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Integrate by parts:

$$\langle f, f \rangle = \langle f, -\partial_t F_0 + \nabla_x^* \nabla_x F_1 \rangle = \langle \partial_t f, F_0 \rangle + \langle \nabla_x f, \nabla_x F_1 \rangle \leq \|\partial_t f\| \cdot \|F_0\| + \|\nabla_x f\| \cdot \|\nabla_x F_1\|.$$

# Poincaré–Lions Inequality

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Proof techniques:

- Traditionally: Bogovskii's operator [Bogovskii 1979]. Estimates are suboptimal.
- Explicit construction using  $H^1$ - $L^2$ - $H^{-1}$  duality and solving divergence equation: [Cao/Lu/Wang '20] which requires  $L^2(\mu)$  to have discrete spectrum. This can be removed following [Brigati/Stoltz '23].
- Recent work of [Eberle/Lörler '24] with finer estimates.

# Weighted Poincaré-Lions Inequality

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Theorem (Brigati/Stoltz/W./Wang '24)

Define tilted measure  $d\mu_W \propto W^{-2}d\mu$ , then

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The operator  $W^2 \nabla_x^* \nabla_x$  is self-adjoint and has spectral gap in  $L^2(\mu_W)$ .

# The Averaging Inequality in Velocity

Add the velocity variable

$$\begin{aligned} \|(\nabla_x \Pi_\nu h, \partial_t \Pi_\nu h)\|_{H^{-1}([0,\tau] \times \mu)} &\lesssim \|h - \Pi_\nu h\|_{L^2([0,\tau] \times \mu \times \nu)} \\ &+ \|(\partial_t + \mathcal{T})h\|_{L^2([0,\tau] \times \mu, H^{-1}(\nu))}. \end{aligned}$$

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Idea: take test functions  $Y, Z \in H^1([0, \tau] \times \mu)$  independent of  $\nu$ ,

$$\begin{aligned} \int (\partial_t \Pi_\nu h Y + \nabla_x \Pi_\nu h \cdot Z) &= \int (Y + \nu \cdot Z)(\partial_t + \mathcal{T})\Pi_\nu h \\ &= \int (Y + \nu \cdot Z)(\partial_t + \mathcal{T})h \\ &\quad + \int (h - \Pi_\nu h)(\partial_t + \mathcal{T})(Y + \nu \cdot Z). \end{aligned}$$

# The Weighted Averaging Inequality in Velocity

Add the velocity variable

$$\begin{aligned} \|(\nabla_x \Pi_\nu h, W^{-1} \partial_t \Pi_\nu h)\|_{H^{-1}([0, \tau] \times \mu)} &\lesssim \|h - \Pi_\nu h\|_{L^2([0, \tau] \times \mu \times \nu)} \\ &\quad + \|(\partial_t + \mathcal{T})h\|_{L^2([0, \tau] \times \mu, H^{-1}(\nu))}. \end{aligned}$$

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$$\begin{aligned} \int (\partial_t W^{-1} \Pi_\nu h Y + \nabla_x \Pi_\nu h \cdot Z) &= \int (W^{-1} Y + v \cdot Z) (\partial_t + \mathcal{T}) \Pi_\nu h \\ &= \int (W^{-1} Y + v \cdot Z) (\partial_t + \mathcal{T}) h \\ &\quad + \int (h - \Pi_\nu h) (\partial_t + \mathcal{T}) (W^{-1} Y + v \cdot Z). \end{aligned}$$

## Theorem (Brigati/Stoltz/W./Wang '24)

If  $\mu$  satisfies (WTPI) with weight  $W$  and  $\sup_{t>0} \int W^\sigma(x) h_t^2 d\mu d\nu < \infty$ , then

① If  $\nu$  satisfies (PI), then

$$\text{Var}_{\mu \otimes \nu}(h_t) \lesssim t^{-\frac{\sigma}{2}}.$$

② If  $\nu$  satisfies (WTPI) with weight  $\mathcal{G}$  and  $\sup_{t>0} \int \mathcal{G}^\delta(v) h_t^2 d\mu d\nu < \infty$ , then

$$\text{Var}_{\mu \otimes \nu}(h_t) \lesssim t^{-\frac{\sigma\delta}{2\sigma+2\delta+4}}.$$

Note that (2) is consistent with (1) when taking  $\delta \rightarrow \infty$ .

# From Weighted to Weak Poincaré–Lions Inequality

Recall the argument

$$\begin{aligned} \int (f - \mu(f))^2 d\mu &\leq s \int_{s > W^2} W^{-2} (f - \mu(f))^2 d\mu + \int_{s \leq W^2} (f - \mu(f))^2 d\mu \\ &\lesssim s \mathbb{E}_\mu |\nabla_x f|^2 + \mu(s \leq W^2) \|f\|_{\text{osc}}^2. \end{aligned}$$

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Repeat this in the space-time domain, and we get the **weak Poincaré–Lions inequality**

$$\|f - (f)\|_{L^2([0, \tau] \times \mu)}^2 \lesssim s \|(\nabla_x f, W^{-1} \partial_t f)\|_{H^{-1}([0, \tau] \times \mu)}^2 + \mu(s \leq W^2) \|f\|_{\text{osc}}^2.$$

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This can be combined with the averaging inequality in velocity and we obtain the next convergence result.

# Second Convergence Result

## Theorem (Brigati/Stoltz/W./Wang '24)

Assume  $h_0 \in L^\infty(\mathbb{R}^{2d})$ , then we have the following convergence rates for explicit examples:

Potential	$\psi(v) = \langle v \rangle^\delta$ $\delta \geq 1$	$\psi(v) = \langle v \rangle^\delta$ $\delta \in (0, 1)$	$\psi(v) =$ $(d + q) \log \langle v \rangle$
$\phi(x) = \langle x \rangle^\alpha$ $\alpha \geq 1$	$\exp(-\lambda t)$	$\exp(-ct^{\frac{\delta}{2-\delta}})$	$t^{-\frac{q}{2}}$
$\phi(x) = \langle x \rangle^\alpha$ $\alpha \in (0, 1)$	$\exp(-ct^{\frac{\alpha}{2-\alpha}})$	$\exp(-ct^{\frac{\alpha\delta}{2\alpha+2\delta-3\alpha\delta}})$	$t^{-\frac{q}{2}-}$
$\phi(x) =$ $(d + p) \log \langle x \rangle$	$t^{-\frac{p}{2}}$	$t^{-\frac{p}{2}-}$	$t^{-\frac{pq}{4+2p+2q}}$

# Comparison with Grothaus/Wang

Potential	$\psi(v) = \langle v \rangle^\delta$ $\delta \geq 1$	$\psi(v) = \langle v \rangle^\delta$ $\delta \in (0, 1)$	$\psi(v) =$ $(d + q) \log \langle v \rangle$
$\phi(x) = \langle x \rangle^\alpha$ $\alpha \geq 1$	$\exp(-\lambda t)$	$\exp(-ct^{\frac{\delta}{4-3\delta}})$	$t^{-\frac{1}{\theta(q)}}$
$\phi(x) = \langle x \rangle^\alpha$ $\alpha \in (0, 1)$	$\exp(-ct^{\frac{\alpha}{8-7\alpha}})$	$\exp(-ct^{\frac{\alpha\delta}{4\alpha+8\delta-11\alpha\delta}})$	$t^{-\frac{1}{\theta(q)}}$
$\phi(x) =$ $(d + p) \log \langle x \rangle$	$t^{-\frac{1}{2\theta(p)}}$	$t^{-\frac{1}{2\theta(p)}}$	$t^{-\frac{1}{2\theta(q)+\theta(p)+2\theta(p)\theta(q)}}$

Here  $\theta(p) = \min\left\{\frac{d+p+2}{p}, \frac{4p+4+2d}{(p^2-2p-2d-4)^+}\right\} > \frac{2}{p}$

# Thanks for listening! I



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# Thanks for listening! II



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