

Exchange-Correlation Functional with Uncertainty Quantification Capabilities for Density Functional Theory

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 - Data setup
 - Numerical results
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 - Atomisation Energies
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Introduction

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 - numerical approximations, e.g., pseudopotentials, PAW, ...

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 - numerical approximations, e.g., pseudopotentials, PAW, ...
- One approximation is necessary:
 - Exchange-Correlation Functional (unknown in general)

Density Functional Theory

- Approximates the ground state energy of a material system with charge density n .
- Minimisation of the energy functional $E^{DFT}[n]$ for a given system

$$\begin{aligned} E^{DFT}[n] &= \int n(\mathbf{r})v(\mathbf{r}) d\mathbf{r} + T_0[n] + U[n] + E^{xc}[n] \\ &= E^b[n] + E^{xc}[n] = E^b[n] + E^x[n] + E^c[n] \end{aligned}$$

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- $E^{xc}[n]$ not known... Need to specify an approximation

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Density Functional Theory: Kohn-Sham method

- Kohn and Sham (1965) introduced a methodology to solve the equations: most common formulation nowadays
- Solve a self-consistent problem using independent electrons in an effective potential:
 - $[-\frac{1}{2}\nabla^2 + v_{\text{eff}}(\mathbf{r})] \psi_i(\mathbf{r}) = \epsilon_i \psi_i(\mathbf{r})$
 - $v_{\text{eff}}(\mathbf{r}) = v(\mathbf{r}) + \int \frac{n(\mathbf{r}')}{|\mathbf{r}-\mathbf{r}'|} d\mathbf{r}' + \frac{\delta E^{\text{xc}}[n]}{\delta n(\mathbf{r})}$
 - $n(\mathbf{r}) = \sum_i |\psi_i(\mathbf{r})|^2$

DFT: Exchange-correlation energies

- How to approximate $E^{\text{xc}}[n]$?
- We can write $E^{\text{xc}}[n] = \int n\epsilon^{\text{xc}}(n; \mathbf{r}) d\mathbf{r}$
 - $n\epsilon^{\text{xc}}(n; \mathbf{r})$: XC energy density

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 - $\varepsilon^{\text{xc}}(n; \mathbf{r}) = \varepsilon^{\text{xc}}[n(\mathbf{r}), \nabla n(\mathbf{r}), \tau(\mathbf{r})]$: meta-GGA
($\tau(\mathbf{r}) = 2 \sum_i \frac{1}{2} |\nabla \psi_i(\mathbf{r})|^2$)

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($\tau(\mathbf{r}) = 2 \sum_i' \frac{1}{2} |\nabla \psi_i(\mathbf{r})|^2$)
- We can add exact exchange: $E^{\text{x}}[n] = -\frac{1}{2} \sum_i \int \frac{\psi_i^*(\mathbf{r}) \psi_i(\mathbf{r}')}{|\mathbf{r} - \mathbf{r}'|} d\mathbf{r} d\mathbf{r}'$
 - But still need to approximate the correlation energy...

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- The double integral in the exact exchange makes it much more costly
- We chose the meta-GGA framework to build our approximation,

$$E^{\text{xc}}[n] = \int n \varepsilon^{\text{xc}}(n(\mathbf{r}), \nabla n(\mathbf{r}), \tau(\mathbf{r})) d\mathbf{r}$$

- We will focus on exchange energy only,

$$E^{\text{xc}}[n] = E^{\text{c}}[n] + \int n \varepsilon^{\text{x}}(n(\mathbf{r}), \nabla n(\mathbf{r}), \tau(\mathbf{r})) d\mathbf{r}$$

- We can transform the dependence on $\nabla n(\mathbf{r})$ and $\tau(\mathbf{r})$ into two dimensionless parameters s and α :

$$s = \frac{|\nabla n|}{2(3\pi^2)^{1/3}n^{4/3}}; \quad \alpha = \frac{\tau - \tau^W}{\tau^{UEG}},$$

- $\tau^W = |\nabla n|^2/8n$: Weizsäcker kinetic energy density
- $\tau^{UEG} = \frac{3}{10}(3\pi^2)^{2/3}n^{5/3}$: UEG kinetic energy density
- Also, we can group all non-local contributions in the *exchange enhancement factor* $F^x(s, \alpha)$,

$$E^x[n] = \int n \varepsilon^x(n, \nabla n, \tau) d\mathbf{r} = \int n \varepsilon_{UEG}^x(n) F^x(s, \alpha) d\mathbf{r}$$

Linear model for the enhancement factor

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Linear model for the enhancement factor

- To specify a model for the exchange energy we just need to specify a model for the enhancement factor
- We specify a linear model

$$F^x(\mathbf{s}, \alpha) = \sum_i \xi_i^x \phi_i(\mathbf{s}, \alpha) = (\boldsymbol{\xi}^x)^T \boldsymbol{\phi}(\mathbf{s}, \alpha)$$

- $\phi_i(\mathbf{s}, \alpha)$: basis functions
- ξ_i^x : linear model coefficients

- Given this model, the exchange energy becomes a linear model

$$\begin{aligned} E^x[n; \xi^x] &= \sum_{i=1}^M \xi_i^x \int n \varepsilon_{UEG}^x(n) \phi_i(s, \alpha) d\mathbf{r} \\ &= \sum_{i=0}^{M-1} \xi_i^x E^x[n; \hat{\mathbf{e}}_i] = (\xi^x)^T \mathbf{E}^x[n; \hat{\mathbf{e}}] \end{aligned}$$

- $E^x[n; \hat{\mathbf{e}}_i] = \int n \varepsilon_{UEG}^x(n) \phi_i(s, \alpha) d\mathbf{r}$ is the “basis exchange energy”, which is obtained if we use $\xi^x = \hat{\mathbf{e}}_i$, i.e., only basis ϕ_i with $\xi_i = 1$,

Linear model for the exchange energy

- How do we choose the basis functions $\phi_i(s, \alpha)$?
 - Follow selection in Wellendorff *et al.*
 - Physically based: inspired in previous non-empirical functionals (PBEsol, MS)
 - Complete basis: 2D Legendre Polynomials

¹J. Wellendorff *et al.* (2014)

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 - Complete basis: 2D Legendre Polynomials
- 2D Legendre polynomials: argument in $[-1, 1]$
- Rational transformations from s, α to the interval $[-1, 1]$

$$PBEsol \rightarrow t_s(s) = \frac{2s^2}{q + s^2} - 1$$

$$MS \rightarrow t_\alpha(\alpha) = \frac{(1 - \alpha^2)^3}{1 + \alpha^3 + \alpha^6}$$

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- The final exchange enhancement model is

$$F^x(s, \alpha) = \sum_i^{M_s} \sum_j^{M_\alpha} \xi_{ij}^x P_i(t_s(s)) P_j(t_\alpha(\alpha))$$

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- How to obtain the coefficients?

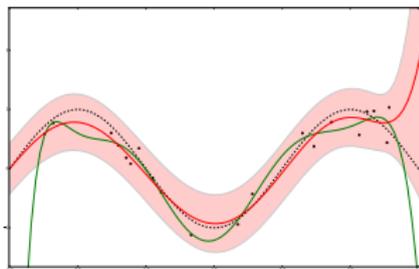
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Bayesian linear regression

- We want to account for uncertainty in the model
- Classical least squares fitting gives a point estimate, not appropriate
- A Bayesian model will give us probability distributions for the coefficients
 - Uncertainty from a limited data set for the regression
 - Uncertainty from an inadequate model

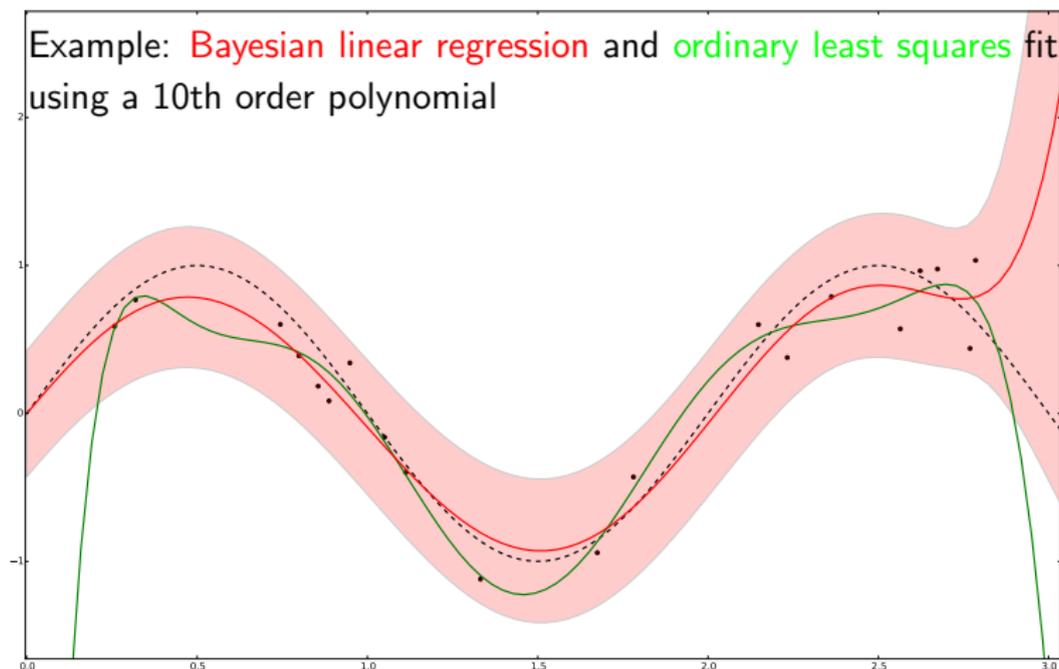
$$p(\boldsymbol{\xi}, \boldsymbol{\beta} \mid \mathbf{t}) \propto \mathcal{L}(\mathbf{t} \mid \mathbf{x}, \boldsymbol{\xi}, \boldsymbol{\beta})p(\boldsymbol{\xi}, \boldsymbol{\beta})$$

Example: Bayesian linear regression and ordinary least squares fit using a 10th order polynomial



¹C. Bishop (2006)

Bayesian linear regression



Some definitions

- $\mathbf{t} = (t_1, t_2, \dots, t_N)^T$: given data (experimental, simulation, mix)
- $\mathbf{n} = (n_1, n_2, \dots, n_N)^T$: input points (densities for DFT)
- $\mathcal{L}(\mathbf{t} \mid \mathbf{n}, \text{model})$: likelihood function
- $\mathcal{N}(x \mid \mu, \nu)$: normal distribution on x with mean μ and variance ν
- $\mathcal{G}(x \mid \alpha, \beta)$: gamma distribution on x with parameters α and β
- $\mathcal{St}(x \mid \mu, \lambda, \nu)$: Student t-distribution on x with parameters μ , λ and ν

Assumptions

- The observed data \mathbf{t} follow on average our model and have a noise ε (includes model inaccuracy),

$$t_i = (\boldsymbol{\xi}^x)^T \mathbf{E}^x[n; \hat{\mathbf{e}}] + \varepsilon_i$$

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$$\mathcal{L}(\mathbf{t} \mid \mathbf{n}, \boldsymbol{\xi}, \beta) = \prod_{i=1}^N \mathcal{N}(t_i \mid \boldsymbol{\xi}^T \mathbf{E}^x[n_i; \hat{\mathbf{e}}], \beta^{-1})$$

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- We choose conjugate priors for $\boldsymbol{\xi}$ and β

- Incorporate prior beliefs into the model
- Depend on extra parameters: *hyperparameters*
- Conjugate priors keep the posterior probability distribution in the same family as the prior probability distribution
- Prior on ξ : $p(\xi \mid \beta, \mathbf{m}_0, \mathbf{S}_0) = \mathcal{N}(\xi \mid \mathbf{m}_0, \beta^{-1}\mathbf{S}_0)$
- Prior on β : $p(\beta \mid a_0, b_0) = \mathcal{G}(\beta \mid a_0, b_0)$
- Joint prior: $p(\xi, \beta) = p(\xi \mid \beta)p(\beta) = \mathcal{N}(\xi \mid \mathbf{m}_0, \beta^{-1}\mathbf{S}_0)\mathcal{G}(\beta \mid a_0, b_0)$

- Probability of a set of parameters given the data
- Proportional to the prior distribution of parameters
- Proportional to the likelihood of the data

$$\begin{aligned} p(\boldsymbol{\xi}, \beta \mid \mathbf{t}) &= \frac{\mathcal{L}(\mathbf{t} \mid \mathbf{x}, \boldsymbol{\xi}, \beta) p(\boldsymbol{\xi}, \beta)}{\int \mathcal{L}(\mathbf{t} \mid \mathbf{x}, \boldsymbol{\xi}, \beta) p(\boldsymbol{\xi}, \beta) d\boldsymbol{\xi} d\beta} \\ &= \mathcal{N}(\boldsymbol{\xi} \mid \mathbf{m}_N, \beta^{-1} \mathbf{S}_N) \mathcal{G}(\beta \mid a_N, b_N) \end{aligned}$$

- The parameters of the posterior depend on those of the prior and the data,

$$\mathbf{S}_N^{-1} = \mathbf{S}_0^{-1} + \Phi^T \Phi; \quad \mathbf{m}_N = \mathbf{S}_N \left[\mathbf{S}_0^{-1} \mathbf{m}_0 + \Phi^T \mathbf{t} \right]$$

$$a_N = a_0 + N/2$$

$$b_N = b_0 + \frac{1}{2} \left(\mathbf{m}_0^T \mathbf{S}_0^{-1} \mathbf{m}_0 - \mathbf{m}_N^T \mathbf{S}_N^{-1} \mathbf{m}_N + \mathbf{t}^T \mathbf{t} \right)$$

- Φ is the *design matrix*

$$\Phi = \begin{pmatrix} E^x[n_1^*; \hat{\mathbf{e}}_0] & \cdots & E^x[n_1^*; \hat{\mathbf{e}}_{M-1}] \\ \vdots & \ddots & \vdots \\ E^x[n_N^*; \hat{\mathbf{e}}_0] & \cdots & E^x[n_N^*; \hat{\mathbf{e}}_{M-1}] \end{pmatrix} = \begin{pmatrix} \mathbf{E}^x[n_1^*; \hat{\mathbf{e}}]^T \\ \vdots \\ \mathbf{E}^x[n_N^*; \hat{\mathbf{e}}]^T \end{pmatrix}$$

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- Rows: For a given data point, a vector with the “basis exchange energies”
- Columns: For a given basis function, its value for every data point

- Once we have our probability distributions for model parameters, what can we say about new data points?

Predictive distribution

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- Prediction averaged over all possible parameters

- Once we have our probability distributions for model parameters, what can we say about new data points?
- Prediction averaged over all possible parameters
- Predictions given with probability distributions

$$\begin{aligned} p(\tilde{\mathbf{t}} \mid \tilde{\mathbf{n}}, \mathbf{t}) &= \int p(\tilde{\mathbf{t}} \mid \tilde{\mathbf{n}}, \boldsymbol{\xi}, \beta) p(\boldsymbol{\xi}, \beta \mid \mathbf{t}) d\boldsymbol{\xi} d\beta \\ &= \int \mathcal{N}(\tilde{\mathbf{t}} \mid \boldsymbol{\xi}^T \mathbf{E}^x[\tilde{\mathbf{n}}; \hat{\mathbf{e}}], \beta^{-1}) \mathcal{N}(\boldsymbol{\xi} \mid \mathbf{m}_N, \beta^{-1} \mathbf{S}_N) \mathcal{G}(\beta \mid a_N, b_N) d\boldsymbol{\xi} d\beta \\ &= St(\tilde{\mathbf{t}} \mid \mu, \lambda, \nu). \end{aligned}$$

- The Student t-distribution $\mathcal{St}(\tilde{t} \mid \mu, \lambda, \nu)$ has parameters

$$\mu = \mathbf{E}^x[\tilde{n}; \hat{\mathbf{e}}]^T \mathbf{m}_N$$

$$\lambda = \frac{a_N}{b_N} \left(1 + \mathbf{E}^x[\tilde{n}; \hat{\mathbf{e}}]^T \mathbf{S}_N \mathbf{E}^x[\tilde{n}; \hat{\mathbf{e}}] \right)^{-1}$$

$$\nu = 2a_N.$$

- Its mean, variance and mode are

$$\mathbf{E}[\tilde{t}] = \mu; \quad \nu > 1$$

$$\text{cov}[\tilde{t}] = \frac{\nu}{\nu - 2} \lambda^{-1} = \frac{1 + \mathbf{E}^x[\tilde{n}; \hat{\mathbf{e}}]^T \mathbf{S}_N \mathbf{E}^x[\tilde{n}; \hat{\mathbf{e}}]}{\text{mode}[\beta]}; \quad \nu > 2$$

$$\text{mode}[\tilde{t}] = \mu$$

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$$\text{cov}[\tilde{t}] = \frac{\nu}{\nu - 2} \lambda^{-1} = \frac{1 + \mathbf{E}^x[\tilde{n}; \hat{\mathbf{e}}]^T \mathbf{S}_N \mathbf{E}^x[\tilde{n}; \hat{\mathbf{e}}]}{\text{mode}[\beta]}; \quad \nu > 2$$

$$\text{mode}[\tilde{t}] = \mu$$

- The variance of the prediction depends on the data point

- If we make several predictions, they are correlated,

$$p(\tilde{\mathbf{t}} \mid \tilde{\mathbf{n}}, \mathbf{t}) = \int p(\tilde{\mathbf{t}} \mid \tilde{\mathbf{n}}, \boldsymbol{\xi}, \beta) p(\boldsymbol{\xi}, \beta \mid \mathbf{t}) d\boldsymbol{\xi} d\beta = \mathcal{S}t(\tilde{\mathbf{t}} \mid \mu, \Lambda, \nu)$$

- The mean and covariance are

$$E[\tilde{\mathbf{t}}] = \tilde{\boldsymbol{\Phi}} \mathbf{m}_N; \quad \text{cov}[\tilde{\mathbf{t}}] = \frac{\mathbf{I} + \tilde{\boldsymbol{\Phi}} \mathbf{S}_N \tilde{\boldsymbol{\Phi}}^T}{\text{mode}[\beta]}$$

- $\tilde{\boldsymbol{\Phi}}$ is analogous to the design matrix where the rows are the points of the predictions

Hyperparameters: Evidence approximation

- One last point to be solved
- How do we obtain the hyperparameters?

Hyperparameters: Evidence approximation

- One last point to be solved
- How do we obtain the hyperparameters?
- We chose the evidence approximation: maximise the log of the marginal likelihood (*evidence function*)

$$\begin{aligned}\log p(\mathbf{t} \mid \mathbf{m}_0, \mathbf{S}_0, a_0, b_0) &= \\ &= \log \int p(\mathbf{t} \mid \boldsymbol{\xi}, \beta, \mathbf{m}_0, \mathbf{S}_0, a_0, b_0) p(\boldsymbol{\xi}, \beta \mid \mathbf{m}_0, \mathbf{S}_0, a_0, b_0) d\boldsymbol{\xi} d\beta \\ &= \mathcal{E}(\mathbf{m}_0, \mathbf{S}_0, a_0, b_0) = \frac{1}{2} \log \frac{|\mathbf{S}_N|}{|\mathbf{S}_0|} - \frac{N}{2} \log(2\pi) + \log \frac{\Gamma(a_N)}{\Gamma(a_0)} + \\ &\quad + a_0 \log(b_0) - a_N \log(b_N)\end{aligned}$$

Hyperparameters: Evidence approximation

- For a model with M parameters, \mathbf{m}_0 , \mathbf{S}_0^{-1} , a_0 and b_0 have $\sim M^2$ parameters

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- Using $\mathbf{S}_0^{-1} = \text{diag}(\alpha_0, \dots, \alpha_{M-1})$, we have $M + 2$ hyperparameters (*Relevance Vector Machine*)
 - Induces sparsity (some of the α_i go to infinity)
 - Only keeps relevant terms: automatic model selection

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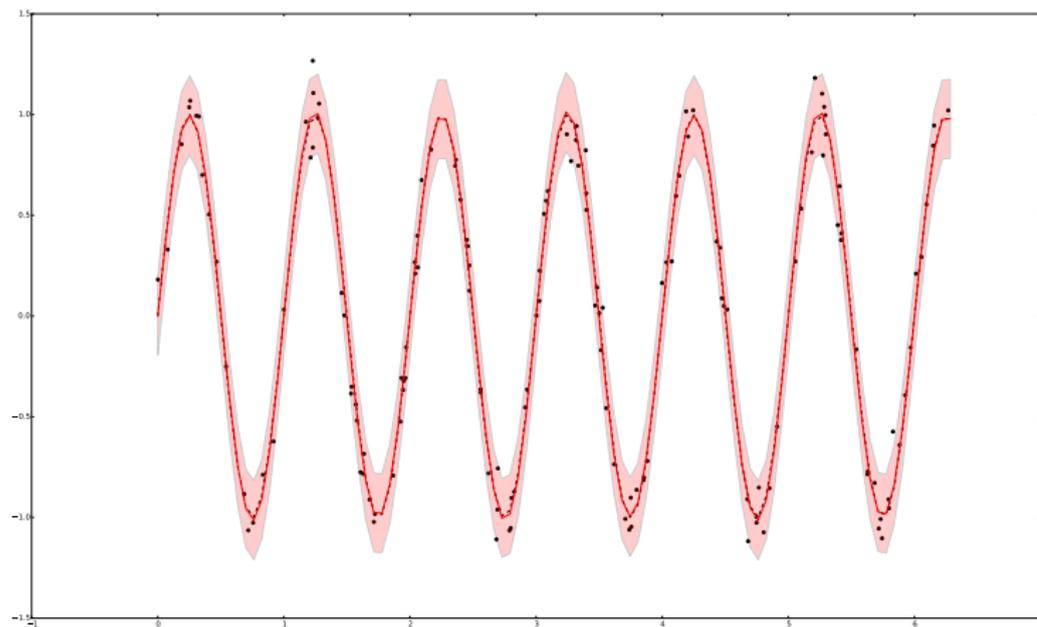
Algorithm 1 Hyperparameter optimisation for the RVM.

- 1: $\mathbf{S}_0^{-1} = \text{diag}(\alpha_0, \dots, \alpha_{M-1})$, $\mathbf{m}_0 = \mathbf{0}$.
 - 2: Initialize α_i from random numbers $r \in (0, 10^{10})$.
 - 3: **repeat**
 - 4: **repeat**
 - 5: **for** all $i = 0, 1, \dots, M - 1$ **do**
 - 6: Update α_i as $\alpha_i^{\text{new}} = \frac{1}{[\mathbf{S}_N]_{ii} + \frac{a_N}{b_N} [\mathbf{m}_N]_i^2}$.
 - 7: Update \mathbf{S}_N , \mathbf{m}_N .
 - 8: **end for**
 - 9: **until** $\Delta\alpha_i < 10^{-5}\%$ or $\alpha_i > 10^{10}$
 - 10: Update a_0, b_0 with a Newton iteration.
 - 11: Update \mathbf{S}_N , \mathbf{m}_N .
 - 12: **until** $\Delta\alpha, \Delta a_0, \Delta b_0 < 10^{-4}\%$
-

Relevance Vector Machine: Example

- Generate data from $f(x) = \sin(2\pi x) + \varepsilon$
 - $\varepsilon \sim \mathcal{N}(\varepsilon | 0, 0.1^2)$
- Use 10 sine basis functions, $\sin(k\pi x)$; $k = 0, 1, \dots, 9$
- Fitted $\mathbf{m}_N = [0, 0.003, 1.006, -0.016, 0, 0, 0, 0, 0, 0]$,
 $mode[\sigma_N] = 0.097$
 - Only three coefficients remain

Relevance Vector Machine: Example



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Training data

- The exchange energy cannot be measured
- Absolute energies cannot be measured
- How to train the model?
 - Easy to get from DFT simulations
 - Experimentally available

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- Absolute energies cannot be measured
- How to train the model?
 - Easy to get from DFT simulations
 - Experimentally available
- Atomisation/cohesive energies
- Given a system $M = A_{n_A} B_{n_B} \dots$,

$$E_{at} = \frac{1}{N} \left(\sum_I n_I E_I - E_M \right); \quad I = A, B, \dots$$

- Decomposing energies into components...

$$\begin{aligned} E_{at} &= \frac{1}{N} \left(\sum_I n_I (E_I^b + E_I^x + E_I^c) - (E_M^b + E_M^x + E_M^c) \right) \\ &= E_{at}^b + E_{at}^x + E_{at}^c, \end{aligned}$$

- Decomposing energies into components...
- Using our exchange energy model...

$$E_{at}^x = \xi^T \frac{1}{N} \left[\sum_I n_I \mathbf{E}^x[n_i; \hat{\mathbf{e}}] - \mathbf{E}^x[n_M; \hat{\mathbf{e}}] \right]$$

- Decomposing energies into components...
- Using our exchange energy model...
- The design matrix becomes...

$$\Phi = \begin{pmatrix} \frac{1}{N} (\sum_{I \in S_1} n_I \mathbf{E}^x[n_i; \hat{\mathbf{e}}] - \mathbf{E}^x[n_{S_1}; \hat{\mathbf{e}}])^T \\ \frac{1}{N} (\sum_{I \in S_2} n_I \mathbf{E}^x[n_i; \hat{\mathbf{e}}] - \mathbf{E}^x[n_{S_2}; \hat{\mathbf{e}}])^T \\ \vdots \\ \frac{1}{N} (\sum_{I \in S_N} n_I \mathbf{E}^x[n_i; \hat{\mathbf{e}}] - \mathbf{E}^x[n_{S_N}; \hat{\mathbf{e}}])^T \end{pmatrix}$$

Atomisation energies

- What is my data vector \mathbf{t} ?
- No access to exchange energy directly

Atomisation energies

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$$\mathbf{t} = \begin{pmatrix} E_{at}^{exp}(s_1) - E_{at}^b[n_1] - E_{at}^c[n_1] \\ E_{at}^{exp}(s_2) - E_{at}^b[n_2] - E_{at}^c[n_2] \\ \vdots \\ E_{at}^{exp}(s_N) - E_{at}^b[n_N] - E_{at}^c[n_N] \end{pmatrix}$$

Atomisation energies

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- We have all we need for the regression

- What if we want to add other data?

¹A. B. Alchagirov *et al.* (2001)

Indirect measurements

- What if we want to add other data?
- Linear on the energy?

¹A. B. Alchagirov *et al.* (2001)

Indirect measurements

- What if we want to add other data?
- Linear on the energy?
 - We can use it directly within our model

¹A. B. Alchagirov *et al.* (2001)

Indirect measurements

- What if we want to add other data?
- Linear on the energy?
- Not linear?

¹A. B. Alchagirov *et al.* (2001)

Indirect measurements

- What if we want to add other data?
- Linear on the energy?
- Not linear?
 - Transform into impact on energy if possible
 - Non-analytically tractable posterior otherwise

¹A. B. Alchagirov *et al.* (2001)

- What if we want to add other data?
- Linear on the energy?
- Not linear?
- Example: equilibrium volume (V_0), bulk modulus and pressure derivative (B_0, B_1) $\rightarrow E(V)$ through equation of state

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- Example: equilibrium volume (V_0), bulk modulus and pressure derivative (B_0, B_1) $\rightarrow E(V)$ through equation of state

$$E(V) = a + b \frac{V_0^{1/3}}{V^{1/3}} + c \frac{V_0^{2/3}}{V^{2/3}} + d \frac{V_0}{V} = \gamma^T \phi(V)$$

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- where

$$\begin{pmatrix} 1 & 1 & 1 & 1 \\ 3 & 2 & 1 & 0 \\ 18 & 10 & 4 & 0 \\ 108 & 50 & 16 & 0 \end{pmatrix} \gamma = \begin{pmatrix} -E_0 \\ 0 \\ 9V_0 B_0 \\ 27V_0 B_0 B_1 \end{pmatrix}$$

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- Example: equilibrium volume (V_0), bulk modulus and pressure derivative (B_0, B_1) $\rightarrow E(V)$ through equation of state

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- From experimental V_0, B_0, B_1 we can also obtain cohesive energies of the strained material

¹A. B. Alchagirov *et al.* (2001)

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- Elementary solids
- Molecules

- Elementary solids
 - 20 cubic solids: 13 training + 7 testing
 - Extended using bulk properties (4 strains each)
- Molecules

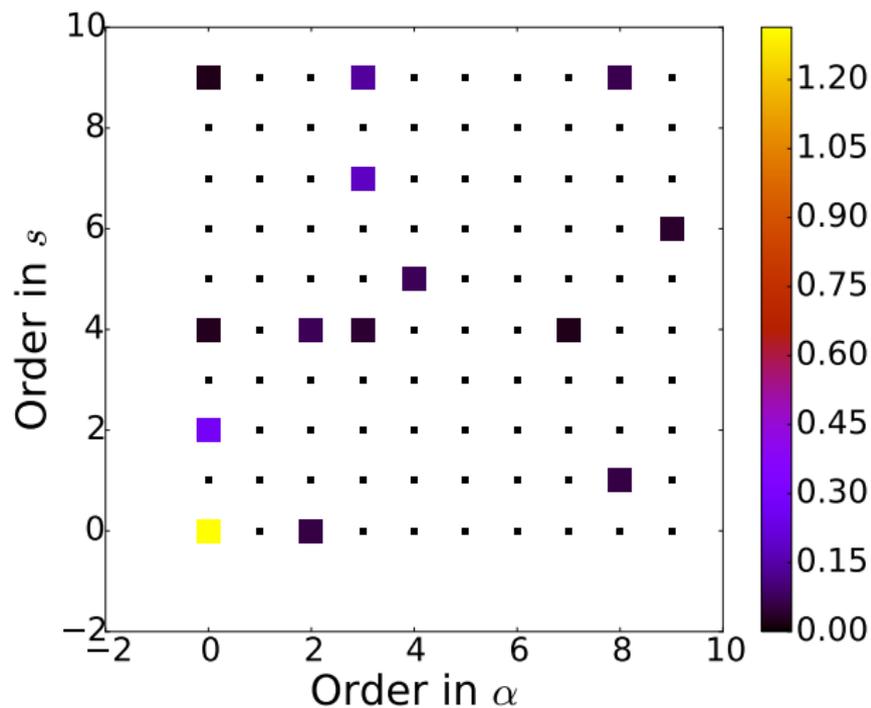
- Elementary solids
 - 20 cubic solids: 13 training + 7 testing
 - Extended using bulk properties (4 strains each)
- Molecules
 - G2/97 data set (small molecules): 120 training + 28 testing
 - Larger molecules from G3/99 only for testing

- Linear model with 10×10 terms

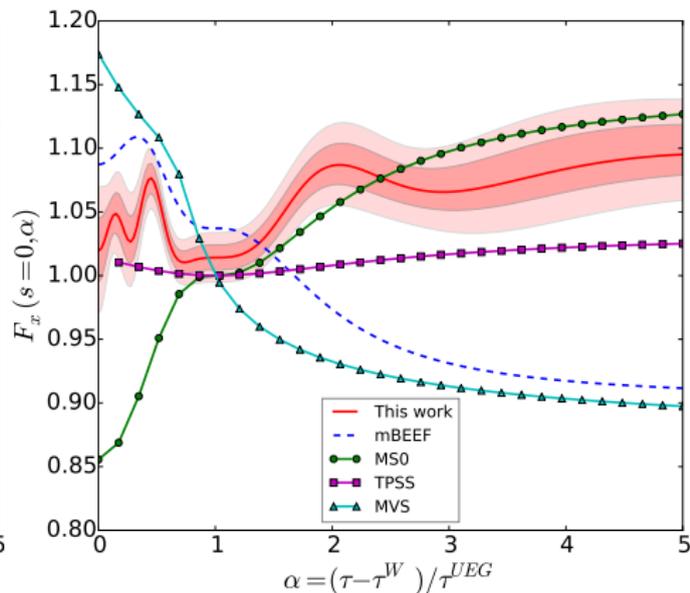
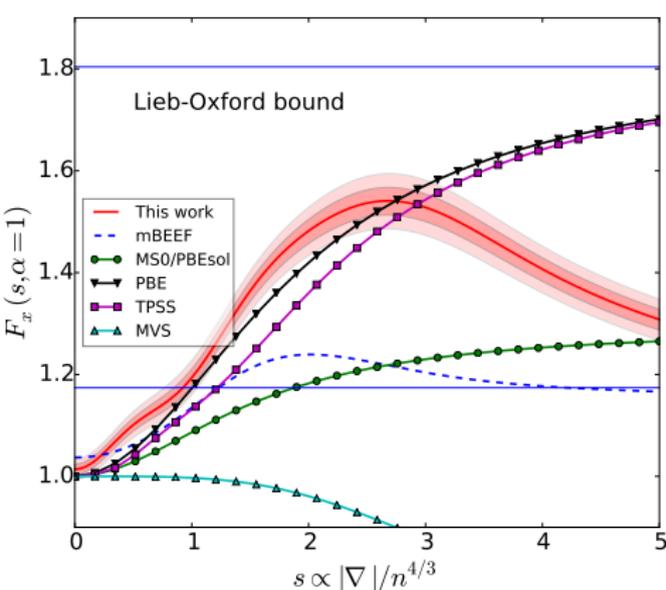
$$F^x(s, \alpha) = \sum_{i=0}^9 \sum_{j=0}^9 \xi_{ij}^x P_i(t_s(s)) P_j(t_\alpha(\alpha))$$

- Linear model with 10×10 terms
- DFT simulations run
 - Using PBE functional
 - Cut-off energy of 800 eV
 - Monkhorst-Pack mesh ($16 \times 16 \times 16$)
 - Relaxing with a maximum force criterion of $0.05 \text{ eV}/\text{\AA}$

Results



- Enhancement factor with 1 and 2 σ intervals

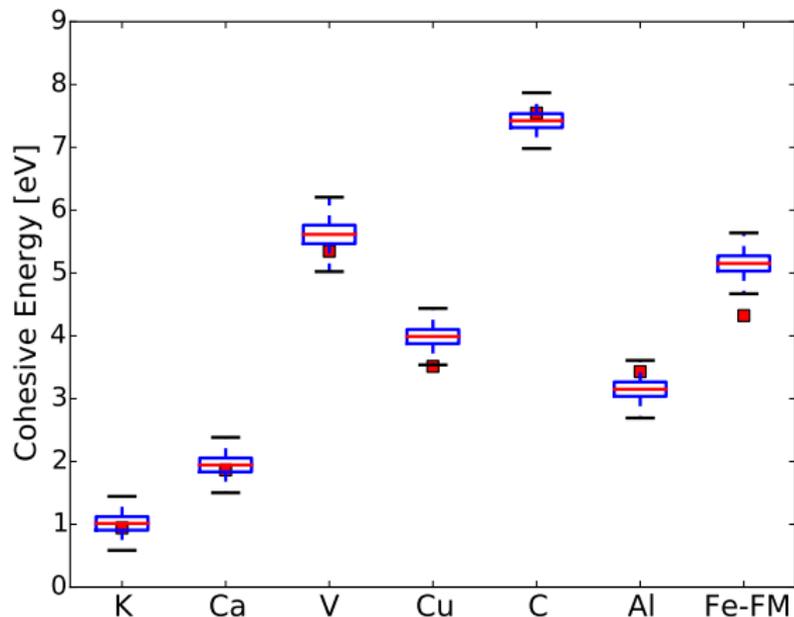


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Atomisation energies

- Prediction of atomisation energies in the 7 test solids



Atomisation energies

- Prediction of atomisation energies in the 7 test solids
- Error in predictions for solids and molecules

| XC functional | Error | G2/97-test | G2/97 | EL20-test | EL20 |
|---------------|-------|------------|-------|-----------|--------|
| This work | MAE | 0.116 | 0.103 | 0.243 | 0.0975 |
| | MARE | 3.27 | 1.46 | 8.56 | 5.62 |
| PBE | MAE | | 0.703 | | 0.238 |
| | MARE | | 5.09 | | 6.88 |

Table: Mean absolute error (in eV) and mean absolute relative error (in %) of the predictions of atomisation energies using the average model for the training sets containing molecules (G2/97) and solids (EL20).

Atomisation energies

- Prediction of atomisation energies in the 7 test solids
- Error in predictions for solids and molecules

| XC functional | Error | G2/97-test | G2/97 | EL20-test | EL20 |
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- G2/97 MAE better than TPSS (0.28 eV), BEEF-vdW (0.16 eV), B3LYP (0.14 eV) or PBE0 (0.21 eV)

Correlation functional

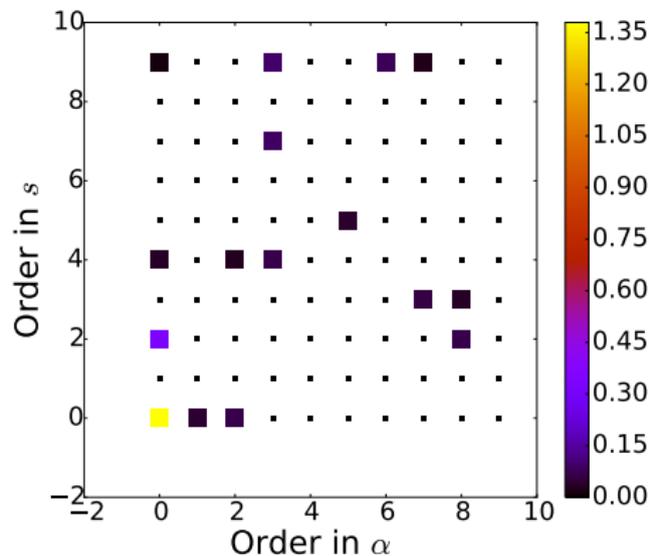
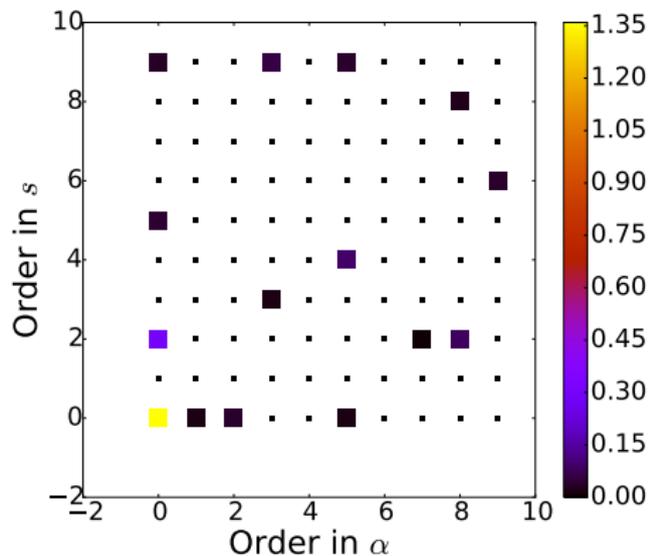
- We assumed a fixed correlation energy functional

Correlation functional

- We assumed a fixed correlation energy functional
- What's the impact of our choice?

Correlation functional

- We assumed a fixed correlation energy functional
- What's the impact of our choice?
- Coefficients with PBEsol (left) and vPBE (right) correlations:



Correlation functional

- We assumed a fixed correlation energy functional
- What's the impact of our choice?
- Errors:

| C functional | Error | G2/97-test | G2/97 | EL20-test | EL20 |
|--------------|-------|------------|-------|-----------|--------|
| PBE | MAE | 0.116 | 0.103 | 0.243 | 0.0975 |
| | MARE | 3.27 | 1.46 | 8.56 | 5.62 |
| PBEsol | MAE | 0.116 | 0.108 | 0.204 | 0.172 |
| | MARE | 2.91 | 1.55 | 6.12 | 4.98 |
| vPBE | MAE | 0.110 | 0.107 | 0.226 | 0.184 |
| | MARE | 2.72 | 1.41 | 6.45 | 5.17 |
| TPSS | MAE | 0.108 | 0.104 | 0.227 | 0.190 |
| | MARE | 2.68 | 1.42 | 6.85 | 5.53 |

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- SL20 test set including
 - 13 elemental solids
 - I-VII, II-VI, III-V and IV-IV compounds
 - 7 of them in the training set (elemental solids)

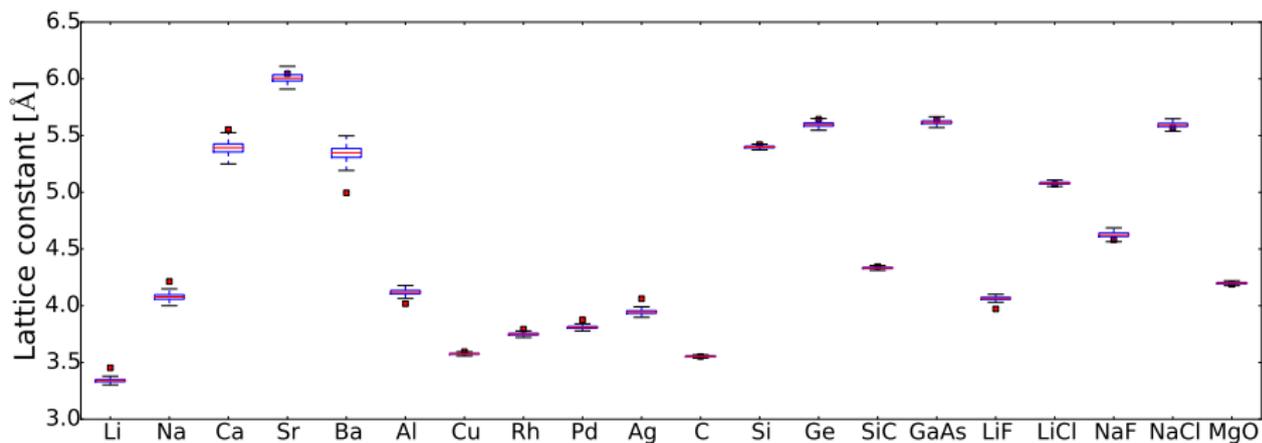
Uncertainty propagation from DFT

- Bulk properties are not obtained directly from DFT simulations
- How to propagate the uncertainty?
- We use a nested Monte Carlo approach
 - Sample model coefficients from the posterior distribution
 - Fit the EOS to the values from this X energy (Bayesian fit)
 - Sample coefficients from the fitting to calculate V_0 , B_0 , B_1

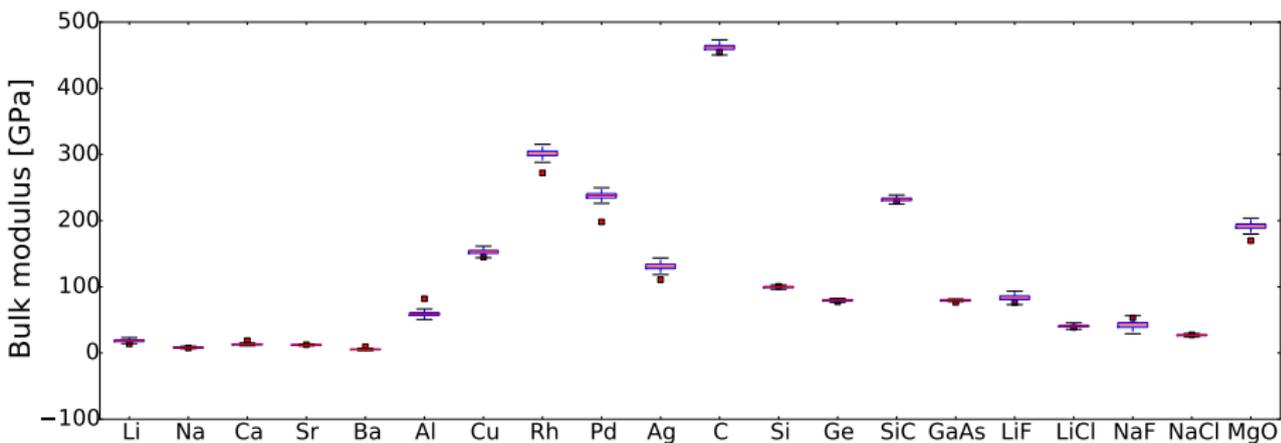
Algorithm 2 Calculation of uncertainty for V_0 and B_0 .

- 1: Input: system s with unit cell $(\mathbf{x}_1, \mathbf{x}_2, \mathbf{x}_3)$.
 - 2: Input: N_1^{max} , N_2^{max} , the maximum iterations.
 - 3: **for** 5 strains $0.95 \leq \sigma_i \leq 1.05$ **do**
 - 4: Strain unit cell by σ_i : $\mathbf{x}_\alpha \rightarrow \sigma_i \mathbf{x}_\alpha$, $\alpha = 1, 2, 3$.
 - 5: Self-consistent simulation of strained system.
 - 6: Keep the self-consistent electron density $n_i^* = n(\sigma_i)$.
 - 7: **end for**
 - 8: $N_1 = 0$
 - 9: **repeat**
 - 10: Sample ξ_{N_1}, β_{N_1} .
 - 11: Non self-consistent simulation with ξ_{N_1}, β_{N_1} and n_i^* .
 - 12: $N_2 = 0$
 - 13: **repeat**
 - 14: Sample γ_{N_2} .
 - 15: Calculate \hat{V}_0, B_0 .
 - 16: **until** $N_2 = N_2^{max}$
 - 17: $N_1 = N_1 + 1$
 - 18: **until** $N_1 = N_1^{max}$
 - 19: Collect statistics on calculated V_0, B_0 .
-

● Equilibrium lattice constants for SL20 materials



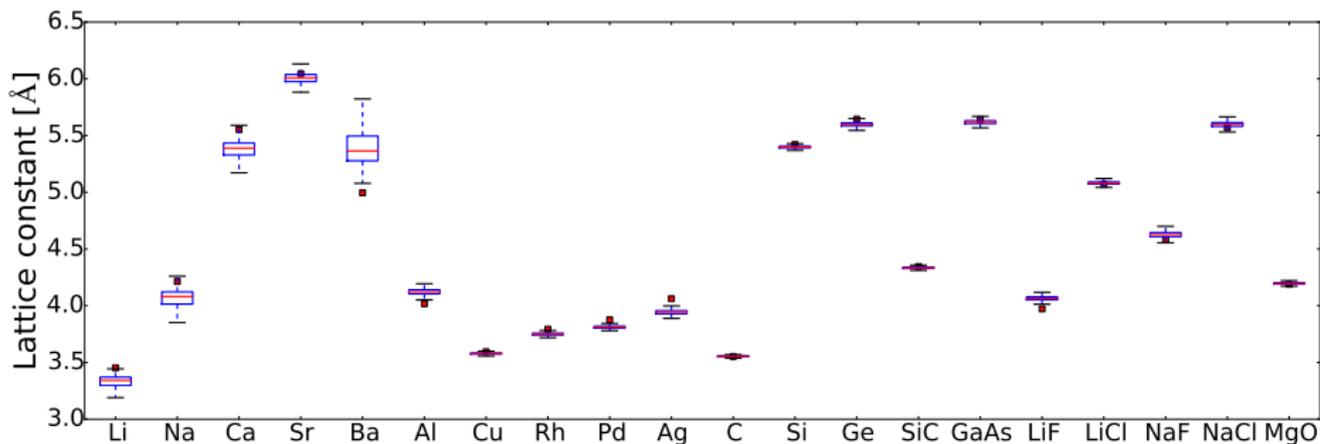
● Equilibrium bulk moduli for SL20 materials



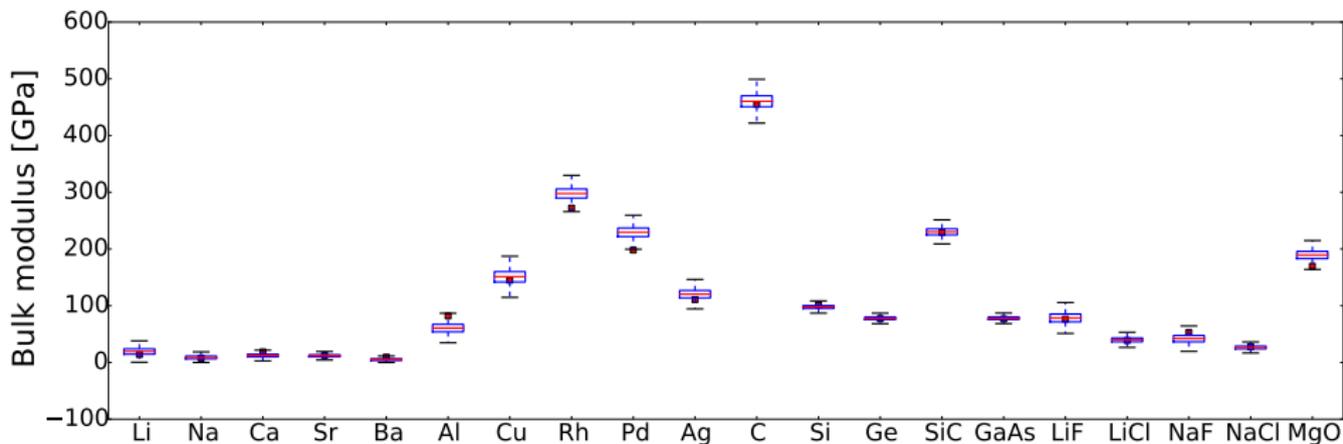
- What if the DFT results have another error sources?

- What if the DFT results have another error sources?
- For example, assume a numerical error with Gaussian distribution and standard deviation 10 mV

● Equilibrium lattice constants for SL20 materials



- Equilibrium bulk moduli for SL20 materials



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- Bayesian framework to obtain an exchange energy functional
 - Use of a linear model
 - Coefficients of the model are random variables
 - Basis functions are fixed

- Bayesian framework to obtain an exchange energy functional
- Use of a relevance vector machine to find hyperparameters
 - Automatic selection of relevant basis functions (model selection)

- Bayesian framework to obtain an exchange energy functional
- Use of a relevance vector machine to find hyperparameters
- Obtained exchange energy from a simulation has an uncertainty
 - Limited data in the training (can be reduced to zero asymptotically with more data)
 - Limited model space, meta-GGA (cannot be reduced to zero asymptotically with more expansion terms)

- Bayesian framework to obtain an exchange energy functional
- Use of a relevance vector machine to find hyperparameters
- Obtained exchange energy from a simulation has an uncertainty
- This uncertainty can be propagated to other derived quantities
 - Bulk properties (shown)
 - Band diagrams, phonon properties, transport coefficients, energy barriers, etc.
 - Further models based on DFT results (e.g., cluster expansion for alloy modelling)

Thank you for your attention!

Acknowledgments

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