Differential Geometric Simulation Methods for Uncertainty Quantification in Complex Model Systems

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- This talk will concentrate on the issues associated with sampling from the induced measures from such models and associated data

System Identification: Nonlinear ODE Oscillator Model



$$\frac{dx_1}{dt} = \frac{k_1}{1+x_n^{\rho}} - m_1 x_1$$
$$\frac{dx_2}{dt} = k_2 x_1 - m_2 x_2$$
$$\vdots$$
$$\frac{dx_n}{dt} = k_n x_{n-1} - m_n x_n$$

Systems Identification - Posterior Inference



Mixing of Markov Chains



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Forward state is *w* and *u* the logarithm of distributed thermal conductivity on Ω , n the unit outward normal on Ω , and *Bi* the Biot number.

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- Consider induced bivariate posterior for varying forms of prior Gaussian measure



MCMC from Diffusions and Geodesics



Riemann manifold Langevin and Hamiltonian Monte Carlo Methods Girolami, M. & Calderhead, B. *J.R.Statist. Soc.* B, with discussion, (2011), 73, 2, 123 - 214.

http://www2.warwick.ac.uk/mgirolami

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- Further Work and Conclusions

• Monte Carlo method employs samples from $\pi(\theta)$ to obtain estimate

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- Success of MCMC reliant upon appropriate proposal design

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Jeffreys, 1948 to first order

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- Can geometric structure be employed in Monte Carlo methodology?

Manifolds

A manifold is a smooth, curved surface: A set *embedded* in \mathbb{R}^d , that locally looks like \mathbb{R}^n (n < d).

Example: the unit sphere (2-sphere): d = 3, n = 2

$$\mathcal{S}_2 = \{x \in \mathbb{R}^3 : \sum_i x_i^2 = 1\}$$



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 $x \in \mathbb{R}^d$ are the embedded coordinates

Coordinate systems and Riemannian metrics

Can parameterise the manifold with a coordinate system in $q \in \mathbb{R}^n$





 $(\sin q_1 \sin q_2, \cos q_1 \sin q_2, \cos q_2), \qquad q_1 \in [0, 2\pi], q_2 \in [0, \pi]$

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The Euclidean metric $\|\cdot\|$ induces a Riemannian metric *G* in the coordinate system:

$$\left\|\mathsf{d}x\right\|^2 = \sum_{i,j} G(q) \mathsf{d}q_i \mathsf{d}q_j$$

M.C. Escher, Heaven and Hell, 1960



Geodesics

Geodesics are the paths of shortest distance.

• On a sphere, these are the great circles





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Given an initial velocity $v(0) \perp x(0)$, we have a nice explicit form

$$\begin{bmatrix} x(t) & v(t) \end{bmatrix} = \begin{bmatrix} x(0) & v(0) \end{bmatrix} \begin{bmatrix} 1 & \\ & \alpha^{-1} \end{bmatrix} \begin{bmatrix} \cos(\alpha t) & -\sin(\alpha t) \\ \sin(\alpha t) & \cos(\alpha t) \end{bmatrix} \begin{bmatrix} 1 & \\ & \alpha \end{bmatrix}$$

here $\alpha = \|v(0)\|$.





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$$\frac{d^2\theta^i}{dt^2} + \sum_{k,l} \Gamma^i_{kl} \frac{d\theta^k}{dt} \frac{d\theta^l}{dt} = 0$$

Fisher-Rao metric

A family of probability densities $\{p(\cdot | \theta) : \theta \in \Theta\}$, the Fisher information forms a natural Riemannian metric, known as Fisher–Rao metric:

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Examples:





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- Consider densities $\mathcal{N}(0,1)$ & $\mathcal{N}(1,1)$ and $\mathcal{N}(0,2)$ & $\mathcal{N}(1,2)$

Normal Density - Euclidean Parameter space



Normal Density - Riemannian Functional space



▶ Continuous Langevin diffusion with invariant measure $\pi(\theta) \equiv \exp(\mathcal{L}(\theta))$

$$d heta = rac{1}{2}oldsymbol{G}^{-1}(oldsymbol{ heta})
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Discretised Langevin diffusion on manifold defines proposal mechanism

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Manifold with constant curvature then proposal mechanism reduces to

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Proposal mechanism diffuses approximately along the manifold





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So for $\phi(\mathbf{x}) = -\log \pi(\mathbf{x}) + \frac{1}{2}\log \det(\mathbf{G}(\mathbf{x}))$ then it follows that marginally

$$p(\mathbf{x}) = \exp\left(-\phi(\mathbf{x})\right) = \pi(\mathbf{x})$$

Numerical integration forms basis of MCMC scheme..... however

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 Stochastic Hamiltonian on Manifold - Lie-Trotter Splitting of Hamiltonian (deterministic and stochastic OU Process) and using symplectic integrator - samples drawn from invariant measure via RMHMC

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Drift function can be seen to describe Lagrangian dynamics

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- Markov chain Monte Carlo from Lagrangian Dynamics Journal of Comp.Graph.Stats, 2014

Univariate finite mixture model

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Figure : Arrows correspond to the gradients and ellipses to the inverse metric tensor. Dashed lines are isocontours of the joint log density

























The joint density for Poisson counts and latent Gaussian field

 $p(\mathbf{y}, \mathbf{x}|\mu, \sigma, \beta) \propto \prod_{i,j}^{64} \exp\{y_{i,j} x_{i,j} - m \exp(x_{i,j})\} \exp(-(\mathbf{x} - \mu \mathbf{1})^{\mathsf{T}} \boldsymbol{\Sigma}_{\theta}^{-1} (\mathbf{x} - \mu \mathbf{1})/2)$

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Metric tensors

$$\begin{aligned} \mathbf{G}(\boldsymbol{\theta})_{i,j} &= \frac{1}{2} trace \left(\boldsymbol{\Sigma}_{\boldsymbol{\theta}}^{-1} \frac{\partial \boldsymbol{\Sigma}_{\boldsymbol{\theta}}}{\partial \theta_i} \boldsymbol{\Sigma}_{\boldsymbol{\theta}}^{-1} \frac{\partial \boldsymbol{\Sigma}_{\boldsymbol{\theta}}}{\partial \theta_j} \right) \\ \mathbf{G}(\mathbf{x}) &= \Lambda + \boldsymbol{\Sigma}_{\boldsymbol{\theta}}^{-1} \end{aligned}$$

where Λ is diagonal with elements $m \exp(\mu + (\Sigma_{\theta})_{i,i})$

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- Latent field metric tensor defining flat manifold is 4096 × 4096, O(N³) obtained from parameter conditional
- MALA requires transformation of latent field to sample with separate tuning in transient and stationary phases of Markov chain

RMHMC for Log-Gaussian Cox Point Processes



Figure : Data, Latent Field, Poisson Mean Field

Table : Sampling the latent variables of a Log-Gaussian Cox Process - Comparison of sampling methods

Method	Time	ESS (Min, Med, Max)	s/Min ESS	Rel. Speed
MALA (Transient)	31,577	(3, 8, 50)	10,605	×1
MALA (Stationary)	31,118	(4, 16, 80)	7836	$\times 1.35$
mMALA	634	(26, 84, 174)	24.1	×440
RMHMC	2936	(1951, 4545, 5000)	1.5	×7070

Heat conduction problem governed by an elliptic partial differential equation in the open and bounded domain Ω:

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Forward state is *w* and *u* the logarithm of distributed thermal conductivity on Ω , n the unit outward normal on Ω , and *Bi* the Biot number.

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 Take one finite element discretisation of domain and one observation at leftmost boundary

- Take one finite element discretisation of domain and one observation at leftmost boundary
- Consider induced bivariate posterior for varying forms of prior Gaussian measure



The dynamics for k-th component of **u** is given by Hamiltons equations

$$\begin{aligned} \frac{d\mathbf{u}_{k}}{dt} &= \frac{\partial \mathcal{H}}{\partial \mathbf{p}_{k}} = \left(\mathbf{G}(\mathbf{u})^{-1}\mathbf{p}\right)_{k} \\ \frac{d\mathbf{p}_{k}}{dt} &= -\frac{\partial \mathcal{H}}{\partial \mathbf{u}_{k}} = -\nabla_{k}J\left(\mathbf{u}\right) - \frac{1}{2}Tr\left[\mathbf{G}(\mathbf{u})^{-1}\frac{\partial \mathbf{G}(\mathbf{u})}{\partial \mathbf{u}_{k}}\right] \\ &+ \frac{1}{2}\mathbf{p}^{T}\mathbf{G}(\mathbf{u})^{-1}\frac{\partial \mathbf{G}(\mathbf{u})}{\partial \mathbf{u}_{k}}\mathbf{G}(\mathbf{u})^{-1}\mathbf{p} \end{aligned}$$

Gradient

Gradient

$$\langle \nabla \mathcal{J}(\boldsymbol{u}), \tilde{\boldsymbol{u}} \rangle = \int_{\Omega} \tilde{\boldsymbol{u}} \boldsymbol{e}^{\boldsymbol{u}} \nabla \boldsymbol{w} \cdot \nabla \lambda \, \boldsymbol{d}\Omega,$$

First Order Forward

$$\int_{\Omega} e^{u} \nabla w \cdot \nabla \hat{\lambda} \, d\Omega + \int_{\partial \Omega \setminus \Gamma_{R}} Bi \, w \hat{\lambda} \, ds = \int_{\Gamma_{R}} \hat{\lambda} \, ds,$$

First Order Adjoint

$$\int_{\Omega} e^{u} \nabla \lambda \cdot \nabla \hat{w} \, d\Omega + \int_{\partial \Omega \setminus \Gamma_{R}} Bi \, \lambda \hat{w} \, ds = -\frac{1}{\sigma^{2}} \sum_{j=1}^{K} \left(w \left(\mathbf{x}_{j} \right) - d_{j} \right) \hat{w} \left(\mathbf{x}_{j} \right),$$

The dynamics for k-th component of u is given by Hamiltons equations

$$\frac{du_{k}}{dt} = \frac{\partial \mathcal{H}}{\partial \mathbf{p}_{k}} = \left(\mathbf{G}(\boldsymbol{u})^{-1}\mathbf{p}\right)_{k}$$
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$$+ \frac{1}{2}\mathbf{p}^{T}\mathbf{G}(\boldsymbol{u})^{-1}\frac{\partial \mathbf{G}(\boldsymbol{u})}{\partial u_{k}}\mathbf{G}(\boldsymbol{u})^{-1}\mathbf{p}$$

Hessian/Fisher Information Matrix

Fisher Matrix-Vector Product

$$\left\langle \left\langle {{\boldsymbol{G}}\left({\boldsymbol{u}} \right),{{\boldsymbol{\tilde u}}} \right
angle ,{{\boldsymbol{u}}^2} }
ight
angle = \int_\Omega {{\tilde {\boldsymbol{u}}}{{\boldsymbol{e}}^{\boldsymbol{u}}}
abla {\boldsymbol{w}} \cdot
abla {{\tilde {\lambda }}^2}\,d\Omega },$$

Second Order Forward

$$\int_{\Omega} e^{u} \nabla w^{2} \cdot \nabla \hat{\lambda} \, d\Omega + \int_{\partial \Omega \setminus \Gamma_{R}} Bi \, w^{2} \hat{\lambda} \, ds = - \int_{\Omega} u^{2} e^{u} \nabla w \cdot \nabla \hat{\lambda} \, d\Omega,$$

Second Order Adjoint

$$\int_{\Omega} e^{u} \nabla \tilde{\lambda^{2}} \cdot \nabla \hat{w} \, d\Omega + \int_{\partial \Omega \setminus \Gamma_{R}} Bi \, \tilde{\lambda^{2}} \, \hat{w} \, ds = -\frac{1}{\sigma^{2}} \sum_{j=1}^{K} w^{2} \left(\mathbf{x}_{j}\right) \hat{w} \left(\mathbf{x}_{j}\right).$$

The dynamics for k-th component of u is given by Hamiltons equations

$$\begin{aligned} \frac{du_k}{dt} &= \frac{\partial \mathcal{H}}{\partial \mathbf{p}_k} = \left(\mathbf{G}(u)^{-1} \mathbf{p} \right)_k \\ \frac{d\mathbf{p}_k}{dt} &= -\frac{\partial \mathcal{H}}{\partial u_k} = -\nabla_k J(u) - \frac{1}{2} \operatorname{Tr} \left[\mathbf{G}(u)^{-1} \frac{\partial \mathbf{G}(u)}{\partial u_k} \right] \\ &+ \frac{1}{2} \mathbf{p}^T \mathbf{G}(u)^{-1} \frac{\partial \mathbf{G}(u)}{\partial u_k} \mathbf{G}(u)^{-1} \mathbf{p} \end{aligned}$$
Derivative of Fisher Information Matrix

Derivative of Fisher Matrix-Matrix product

$$\left\langle \left\langle \left\langle T\left(u\right),\tilde{u}\right\rangle,u^{2}\right\rangle,u^{3}\right\rangle :=\left\langle \nabla\left\langle \left\langle G\left(u\right),\tilde{u}\right\rangle,u^{2}\right\rangle,u^{3}\right\rangle \\ =\int_{\Omega}\tilde{u}u^{3}e^{u}\nabla w\cdot\nabla\tilde{\lambda}^{2}\,d\Omega+\int_{\Omega}\tilde{u}e^{u}\nabla w^{3}\cdot\nabla\tilde{\lambda}^{2}\,d\Omega+\int_{\Omega}\tilde{u}e^{u}\nabla w\cdot\nabla\lambda^{2,3}\,d\Omega \right.$$

Third Order Forward

$$\int_{\Omega} e^{u} \nabla w^{3} \cdot \nabla \hat{\lambda} \, d\Omega + \int_{\partial \Omega \setminus \Gamma_{R}} B_{i} \, w^{3} \hat{\lambda} \, ds = - \int_{\Omega} u^{3} e^{u} \nabla w \cdot \nabla \hat{\lambda} \, d\Omega$$

Third Order Adjoint

$$\begin{split} \int_{\Omega} e^{u} \nabla \lambda^{2,3} \cdot \nabla \hat{w} \, d\Omega + \int_{\partial \Omega \setminus \Gamma_{R}} B_{i} \, \lambda^{2,3} \hat{w} \, ds &= -\frac{1}{\sigma^{2}} \sum_{j=1}^{K} w^{2,3} \left(\mathbf{x}_{j} \right) \hat{w} \left(\mathbf{x}_{j} \right) \\ &- \int_{\Omega} u^{3} e^{u} \nabla \tilde{\lambda^{2}} \cdot \nabla \hat{w} \, d\Omega, \end{split}$$

Two-parameter Case



- $\varepsilon = 0.7$ for simRMMALA and RMMALA
- $\varepsilon = 0.02$, L = 100 for simRMHMC and RMHMC

1025-parameter Case



1025-parameter Case



Bui-Thanh, T., and Girolami, M., *Solving Large-scale PDE using Riemann Manifold Hamiltonian MCMC*, Inverse Problems, 30, 92014) 114014, IoP Publishing.

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